

Gaussian Quantum Markov Semigroups: Irreducibility and Normal Invariant States

Franco Fagnola

Politecnico di Milano

(joint work with J. Agredo and D. Poletti)



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1 Gaussian QMS

- Fock space & CCR
- Gaussian - GKLS
- Construction (minimal semigroup)
- Action on Gaussian states

2 Invariant States

- One mode $d = 1$
- Existence: sharp conditions
- Interpretation
- Uniqueness & convergence

3 Irreducibility

- Two noise operators
- One noise operator L and the Hörmander condition
- Irreducibility: multimode $d > 1$

$\mathfrak{h} = \Gamma(\mathbb{C}^d)$ $(e_n)_n$ o.n. basis $e_\alpha = e_{n_1} \otimes \cdots \otimes e_{n_d}$
Annihilation and creation operators a_j, a_j^\dagger

$$a_j e_n = \sqrt{n_j} e_{n_1} \otimes \cdots \otimes e_{n_{j-1}} \otimes \cdots \otimes e_{n_d},$$
$$a_j^\dagger e_n = \sqrt{n_j + 1} e_{n_1} \otimes \cdots \otimes e_{n_{j+1}} \otimes \cdots \otimes e_{n_d}.$$

Domain $\{\xi = \sum_\alpha \xi_n e_n \mid \sum_n |n| |\xi_n|^2 < \infty\}$, CCR $[a_j, a_k^\dagger] = \delta_{jk} \mathbb{1}$
 q_j, p_j the position and momentum operators on \mathfrak{h}

$$q_j = (a_j + a_j^\dagger) / \sqrt{2}, \quad p_j = (a_j - a_j^\dagger) / \sqrt{2} i.$$

q_j, p_k selfadjoint and satisfy $[q_j, p_k] = i \delta_{jk} \mathbb{1}$

$$H = \sum_{j,k=1}^d \left(\Omega_{jk} a_j^\dagger a_k + \frac{\kappa_{jk}}{2} a_j^\dagger a_k^\dagger + \frac{\bar{\kappa}_{jk}}{2} a_j a_k \right) + \frac{1}{2} \sum_{j=1}^d \left(\zeta_j a_j^\dagger + \bar{\zeta}_j a_j \right)$$

$$L_\ell = \sum_{j=1}^d \left(v_{\ell j} a_j + u_{\ell j} a_j^\dagger \right) \quad \ell = 1, \dots, m; \quad m \leq 2d$$

$v_{\ell j}, u_{\ell j}, \zeta_j \in \mathbb{C}, \quad \Omega_{jk}, \kappa_{jk} \in \mathbb{C} \quad m \text{ noise multiplicity}$

$\Omega = (\Omega_{jk})_{1 \leq j, k \leq d}$ Hermitian $K = (\kappa_{jk})_{1 \leq j, k \leq d}$ symmetric

$D := \text{Lin}\{e_n \mid n \in \mathbb{N}^d\}$ (finite linear span) $\xi', \xi \in D, x \in \mathcal{B}(\mathfrak{h})$

$$\begin{aligned} \mathcal{L}(x)[\xi', \xi] &= i \langle H\xi', x\xi \rangle - i \langle \xi', xH\xi \rangle \\ &\quad - \frac{1}{2} \sum_{\ell=1}^m (\langle \xi', xL_\ell^* L_\ell \xi \rangle - 2 \langle L_\ell \xi', xL_\ell \xi \rangle + \langle L_\ell^* L_\ell \xi', x\xi \rangle) \end{aligned}$$

H, L_ℓ unbounded \rightsquigarrow minimal qds $\mathcal{T} \rightsquigarrow$ it is unique and $\mathcal{T}_t(\mathbb{1}) = \mathbb{1}$

Theorem

There exists a unique QMS, $\mathcal{T} = (\mathcal{T}_t)_t$ such that

$$\left. \frac{d}{dt} \langle \xi', \mathcal{T}_t(x)\xi \rangle \right|_{t=0} = \mathcal{L}(x)[\xi', \xi] \quad \forall \xi, \xi' \in D$$

Action on Weyl operators

$$\text{Weyl operator} \quad W(z) = \exp \left(\sum_{j=1}^d (z_j a_j^\dagger - \bar{z}_j a_j) \right)$$

define matrices $U = (u_{j\ell})$, $V = (v_{j\ell})$ & *real* linear operators on \mathbb{C}^d

$$Zz = \left[\frac{\overline{(U^*U - V^*V)}}{2} + i\Omega \right] z + \left[\frac{(U^T V - V^T U)}{2} + iK \right] \bar{z}$$

$$Cz = \left(\overline{U^*U + V^*V} \right) z + (U^T V + V^T U) \bar{z}$$

$$\mathcal{T}_t(W(z)) = \exp \left(-\frac{1}{2} \int_0^t \operatorname{Re} \langle e^{sZ} z, C e^{sZ} z \rangle ds + i \int_0^t \operatorname{Re} \langle \zeta, e^{sZ} z \rangle ds \right) W(e^{tZ} z)$$

Definition

ρ is a gaussian state if $\forall z \in \mathbb{C}^d$

$$\hat{\rho}(z) = \text{tr}(\rho W(z)) = \exp\left(-i \text{Re} \langle \omega, z \rangle - \frac{1}{2} \text{Re} \langle z, Sz \rangle\right)$$

where $\omega \in \mathbb{C}^d$ (mean), $S : \mathbb{C}^d \rightarrow \mathbb{C}^d$ invertible, real linear.

$\rho =: \rho_{\omega, S}$.

If $z = x + iy$ $x, y \in \mathbb{R}^d$

$$\begin{aligned} \hat{\rho}(z) = \text{tr}(\rho W(z)) &\leftrightarrow \mathbb{E}_{\rho} \left[e^{-ix\sqrt{2}p + iy\sqrt{2}q} \right] \\ &= \exp \left\{ -i \left\langle \begin{pmatrix} \text{Re } \omega \\ \text{Im } \omega \end{pmatrix}, \begin{pmatrix} x \\ y \end{pmatrix} \right\rangle - \frac{1}{2} \left\langle \begin{pmatrix} x \\ y \end{pmatrix}, \mathbf{S} \begin{pmatrix} x \\ y \end{pmatrix} \right\rangle \right\} \end{aligned}$$

Theorem

If $\rho = \rho_{(\omega, S)}$ then $\rho_t := \mathcal{T}_{*t}(\rho)$ is still a gaussian state $\rho_{(\omega_t, S_t)}$ with

$$\begin{aligned}\omega_t &= e^{tZ^T} \omega - \int_0^t e^{sZ^T} \zeta ds \\ S_t &= e^{tZ^T} S e^{tZ} + \int_0^t e^{sZ^T} C e^{sZ} ds,\end{aligned}$$

where Z and C are the *real* linear operators

$$\begin{aligned}Zz &= [(\overline{U^*U - V^*V})/2 + i\Omega] z + [(U^T V - V^T U)/2 + iK] \bar{z} \\ Cz &= (\overline{U^*U + V^*V}) z + (U^T V + V^T U) \bar{z}\end{aligned}$$

Conversely, (essentially proved in

B. Demoen, P. Vanheuverzwijn, A. Verbeure, Lett. Math. Phys. 2, 161 (1977)

M. Fannes, Commun. Math. Phys 51, 55 (1976))

if \mathcal{T} is a w^* -continuous semigroup of normal CP maps on $\mathcal{B}(\Gamma(\mathbb{C}^d))$ s.t.,

$\forall \rho = \rho_{(\omega, S)}$ Gaussian state, $\rho_t := \mathcal{T}_{*t}(\rho)$ is still a gaussian state, then one can find a GKLS generator with unbounded $H, L_\ell \dots$

Definition

QMSs, $\mathcal{T} = (\mathcal{T}_t)_t$, constructed from the above choices of H, L_ℓ are called Gaussian QMSs.

One mode Fock space $d = 1$

$$H = \Omega a^\dagger a + \frac{\kappa}{2} a^{\dagger 2} + \frac{\bar{\kappa}}{2} a^2 + \frac{\zeta}{2} a^\dagger + \frac{\bar{\zeta}}{2} a$$
$$L_\ell = \bar{v}_\ell a + u_\ell a^\dagger, \quad \ell = 1, 2$$

with $\Omega \in \mathbb{R}$, $\kappa, \zeta, v_\ell, u_\ell \in \mathbb{C}$ (not all = 0).

Sometimes, when stated, we consider *only one* Kraus operator L .

Definition

Invariant state: ρ positive, $\text{tr } \rho = 1$, $\mathcal{T}_{*t}(\rho) = \rho \quad \forall t \geq 0$.

Theorem

A Gaussian semigroup has a normal *invariant state* if and only if

$$\gamma := \frac{1}{2} \sum_{\ell} (|v_{\ell}|^2 - |u_{\ell}|^2) > 0 \quad \text{and} \quad \gamma^2 + \Omega^2 - |\kappa|^2 > 0$$

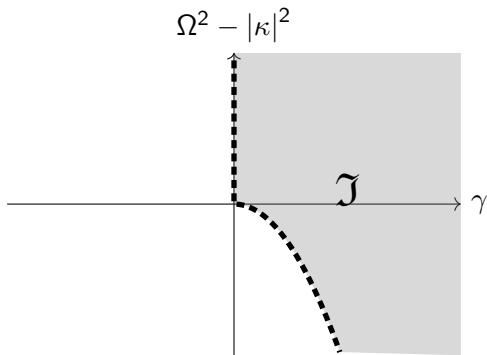
the invariant state is *unique* in the class of normal states and *Gaussian* with

$$\omega = (Z^T)^{-1} \zeta = \frac{(-\gamma + i\Omega)\zeta - i\kappa\bar{\zeta}}{\gamma^2 + \Omega^2 - |\kappa|^2}, \quad S = \int_0^{\infty} e^{sZ^T} C e^{sZ} ds.$$

Note: sum on $\ell = 1, 2$ for two ops L_{ℓ} , sum on $\ell = 1$ for *only one* L_{ℓ}

Recall: $H = \Omega a^{\dagger} a + \frac{\kappa}{2} a^{\dagger 2} + \frac{\bar{\kappa}}{2} a^2 + \frac{\zeta}{2} a^{\dagger} + \frac{\bar{\zeta}}{2} a$, $L_{\ell} = \bar{v}_{\ell} a + u_{\ell} a^{\dagger}$

$$H = \Omega a^\dagger a + \frac{\kappa}{2} a^{\dagger 2} + \frac{\zeta}{2} a^\dagger + \text{h.c.} \quad L_\ell = \bar{v}_\ell a + u_\ell a^\dagger \quad \gamma = \sum_\ell \frac{|v_\ell|^2 - |u_\ell|^2}{2}$$



$$\mathfrak{J} := \left\{ (\gamma, \Omega^2 - |\kappa|^2) \mid \gamma > 0, \gamma^2 + \Omega^2 - |\kappa|^2 > 0 \right\}$$

$$H = \Omega a^\dagger a + \frac{\kappa}{2} a^{\dagger 2} + \frac{\bar{\kappa}}{2} a^2 + \frac{\zeta}{2} a^\dagger + \frac{\bar{\zeta}}{2} a$$
$$\lambda_{\pm} = -\gamma \pm \sqrt{|\kappa|^2 - \Omega^2} \quad \text{eigenvalues of } \mathbf{Z}$$

An invariant state exists if and only if the parameters are in \mathfrak{J}

$$\gamma > 0 \text{ and } \gamma^2 + \Omega^2 - |\kappa|^2 > 0 \quad \Leftrightarrow \quad e^{t\mathbf{Z}} \rightarrow 0 \text{ as } t \rightarrow +\infty$$

- H semi-bounded (below or above) iff $\Omega^2 - |\kappa|^2 > 0$
- H semi-bounded \Rightarrow existence of invariant state determined only by noise,
- H not semi-bounded \Rightarrow a “strong” noise is necessary for existence of an invariant state.

Intuition behind invariant states

For an invariant state

$$\mathbf{S}_t = \mathbf{S}_0 \iff \int_0^t e^{s\mathbf{Z}^T} (\mathbf{Z}^T \mathbf{S}_0 + \mathbf{S}_0 \mathbf{Z} + \mathbf{C}) e^{s\mathbf{Z}} ds = 0$$

The operator

$$\mathbf{S} = \int_0^\infty e^{s\mathbf{Z}^T} \mathbf{C} e^{s\mathbf{Z}} ds$$

is a solution of $\mathbf{Z}^T \mathbf{S} + \mathbf{S} \mathbf{Z} = -\mathbf{C}$ if it exists.

It *does exist* when \mathbf{Z} is stable

$$\mathbf{Z} = \begin{pmatrix} -\gamma - \text{Im}(\kappa) & \text{Re}(\kappa) - \Omega \\ \text{Re}(\kappa) + \Omega & -\gamma + \text{Im}(\kappa) \end{pmatrix}, \quad \lambda_{1,2} = -\gamma \pm \sqrt{|\kappa|^2 - \Omega^2}$$

Recall that an invariant state exists if the parameters are in \mathfrak{I}

Theorem

If

$$\gamma > 0 \text{ and } \gamma^2 + \Omega^2 - |\kappa|^2 > 0$$

and ρ_∞ is the unique invariant state, for any initial state ρ_0

$$\lim_{t \rightarrow +\infty} \mathcal{T}_{*t}(\rho_0) = \rho_\infty$$

Remark. ρ_∞ is either faithful or pure

Multimode $d > 1$

- All eigenvalues of Z with $\Re(\lambda) < 0 \Rightarrow \exists$ invariant state(s),
- If \exists eigenvalues of Z with $\Re(\lambda) > 0 \Rightarrow$ might not \exists invariant states.

If \exists eigenvalues with $\Re(\lambda) = 0$ both situations.

Irreducibility

A semigroup is **irreducible** if and only if for all p projection

$$\mathcal{T}_t(p) \geq p \quad \Rightarrow \quad p = 0, \quad p = \mathbb{1}$$

$\mathcal{T}_t : \mathcal{A} \rightarrow \mathcal{A}$, $\mathcal{T}_t(p) \geq p$ means

- $\mathcal{T}_t(p\mathcal{A}p) \subseteq p\mathcal{A}p$
- ρ density s.t. $\rho = p\rho p \Leftrightarrow \mathcal{T}_{*t}(\rho) = p\mathcal{T}_{*t}(\rho)p$

One mode $d = 1$ case.

Theorem

Gaussian QMS

Number of Kraus' operators (L_ℓ)

two

one

Irreducible

Consider $2\Omega\bar{v}u - \bar{v}^2\kappa - u^2\bar{\kappa}$

$\neq 0$

$= 0$

Irreducible

Not Irreducible

Theorem (FF & R. Rebolledo)

p projection, $\text{Ran}(p)$ its range. $\mathcal{T}_t(p) \geq p$ if and only if

- i $\text{Ran}(p)$ is invariant for the strongly continuous contraction semigroup e^{tG}
- ii $L_\ell u = pL_\ell u$, for $u \in \text{Dom}(G) \cap \text{Ran}(p)$

$$G = -iH - \frac{1}{2} \sum_{\ell} L_{\ell}^* L_{\ell}$$

Two linearly independent $L_{\ell} \Rightarrow$ a candidate invariant subspace must be invariant for $a, a^{\dagger} \dots \Rightarrow$ irreducible!

In the case of a single Kraus operator L a candidate invariant subspace should be invariant also for

$$-2[G, L] = [L^*, L]L + 2i(\bar{v}\Omega - u\bar{\kappa})a + 2i(\bar{v}\kappa - u\Omega)a^\dagger + 2i(\bar{v}\zeta - u\bar{\zeta})$$

We have another linear operator in creation and annihilation operators

$$\tilde{L} = (\bar{v}\Omega - u\bar{\kappa})a + (\bar{v}\kappa - u\Omega)a^\dagger$$

The condition for linear independence of L and \tilde{L} is

$$\det \begin{bmatrix} \bar{v}\Omega - u\bar{\kappa} & \bar{v}\kappa - u\Omega \\ \bar{v} & u \end{bmatrix} \neq 0$$

$$L_\ell = \sum_{j=1}^d \left(\bar{v}_{\ell j} a_j + u_{\ell j} a_j^\dagger \right) \rightsquigarrow \begin{bmatrix} \bar{v}_{\ell \bullet} \\ u_{\ell \bullet} \end{bmatrix}$$

$$-2[G, L_\ell] = 2i[H, L_\ell] + \sum_k [L_k^*, L_\ell] L_\ell + \sum_k L_k^* [L_k, L_\ell]$$

linear in a_j a_j^\dagger \rightsquigarrow vector $\begin{bmatrix} \bar{v}_{\ell \bullet}^{(1)} \\ u_{\ell \bullet}^{(1)} \end{bmatrix}$

Iterate commutators $[G, [G, L_\ell]]$ and get vectors

$$\begin{bmatrix} \bar{v}_{\ell \bullet}^{(n)} \\ u_{\ell \bullet}^{(n)} \end{bmatrix} \quad n \geq 0.$$

Conjecture: the gaussian QMS is irreducible if and only if

$$\text{Lin}_{\mathbb{R}} \left\{ \left[\begin{array}{c} \bar{v}_{\ell \bullet}^{(n)} \\ u_{\ell \bullet}^{(n)} \end{array} \right] \mid n \geq 0 \right\}$$

as a **real** subspace of \mathbb{C}^d is $2d$ -dimensional.

Thank you!

