

In memory of Professor Andrzej Kossakowski

Dariusz Chruściński

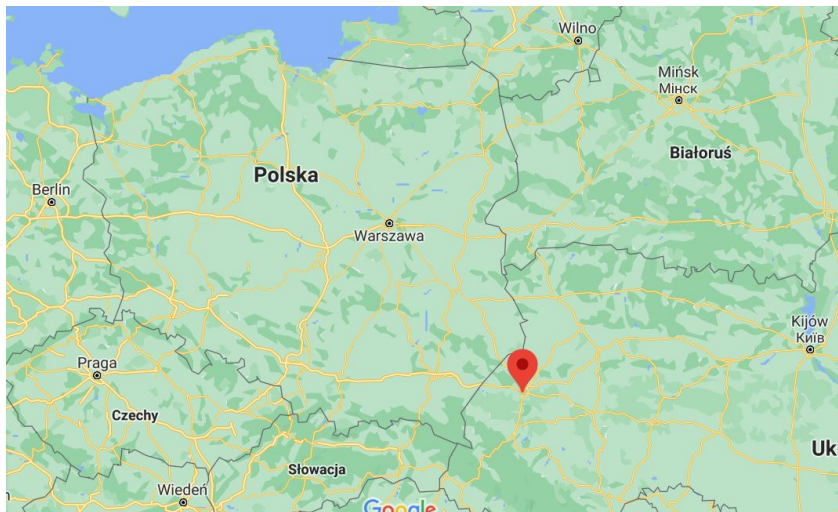
Nicolaus Copernicus University, Toruń, POLAND

SMP 52, Toruń June 14 2021

Andrzej Kossakowski (1938-2021)



Born in Lwów (Lviv)



Nicolaus Copernicus University (PhD 1966)





Scientific activity

- open quantum systems
- quantum statistical mechanics (3D Ising model)
- quantum chaos
- quantum entropies
- quantum information theory
- quantum probability
- ...

Markovian semigroups

Vol. 3 (1972)

REPORTS ON MATHEMATICAL PHYSICS

No. 4

ON QUANTUM STATISTICAL MECHANICS OF NON-HAMILTONIAN SYSTEMS

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*(Received November 9, 1971)**

An axiomatic definition of time evolution (dynamical semi-group) of a physical system has been given. A dynamical semi-group is defined as a one-parameter semi-group of linear endomorphisms of the set of all density operators corresponding to the physical system in question. Some classes of dynamical semi-groups (quantum Poisson and Brownian processes) induced by Markov processes on topological groups are described. Examples of dynamical semi-groups for the harmonic oscillator are given.

Lindblad cited in his 1976 paper

Generators of Quantum Semigroups

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As special cases we can mention the generators of Gaussian semigroups.

$$L(X) = VXV - \frac{1}{2}\{V^2, X\} \quad V \text{ s.a.},$$

and Poisson semigroups

$$L(X) = V^+XV - X \quad V \text{ unitary}.$$

Kossakowski [8] treated these cases in detail.

Markovian semigroups

THEORETICAL PHYSICS

On Necessary and Sufficient Conditions for a Generator of a Quantum Dynamical Semi-Group

by

A. KOSSAKOWSKI

Presented by W. RUBINOWICZ on June 14, 1972

Summary. A dynamical semi-group has been defined as a one-parameter contracting semi-group of trace preserving linear operators on the real Banach space $L^1(\mathcal{H})$ of self-adjoint trace class linear operators on a separable complex Hilbert space \mathcal{H} . It has been proved that a linear operator L with the domain $D(L)$ and the range $R(L)$ both in $L^1(\mathcal{H})$ generates dynamical semi-group $\mathcal{S}(\cdot) = \{S_t; t \geq 0\}$ iff the domain $D(L)$ is dense in $L^1(\mathcal{H})$, $R(I-L) = L^1(\mathcal{H})$, L is a dissipative operator in the sense of Lumer and Phillips, and $\text{Tr}(L\rho) = 0$ for all $\rho \in D(L)$. The resulting master equation $d_t/dt(S_t \rho) = L(S_t \rho)$, $\rho \in D(L)$, satisfies the positivity and normalization requirements.

Positive trace-preserving semigroups

Theorem 2.1. In order for a linear map $L : M(N) \rightarrow M(N)$ to be the generator of a dynamical semigroup of $M(N)$ it is necessary and sufficient that the conditions

$$\operatorname{tr}[P_r(LP_s)] \geq 0, \quad r \neq s = 1, 2, \dots, N \quad (2.1)$$

and

$$\sum_{r=1}^N \operatorname{tr}[P_r(LP_s)] = 0, \quad s = 1, 2, \dots, N \quad (2.2)$$

hold for all $\{P_1, P_2, \dots, P_N\} \in \rho_N$. Condition (2.2) is necessary and sufficient for L to generate a trace preserving semigroup, whereas (2.1) expresses the positivity requirement.

Completely positive dynamical semigroups of N -level systems*

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(Received 19 March 1975)

We establish the general form of the generator of a completely positive dynamical semigroup of an N -level quantum system, and we apply the result to derive explicit inequalities among the physical parameters characterizing the Markovian evolution of a 2-level system.

GKS 1976

Theorem 2.2. A linear operator $L : M(N) \rightarrow M(N)$ is the generator of a completely positive dynamical semi-group of $M(N)$ if it can be expressed in the form

$$L : \rho \rightarrow L\rho = -i[H, \rho] + \frac{1}{2} \sum_{i,j=1}^{N^2-1} c_{ij} \{ [F_i, \rho F_j^*] + [F_i \rho, F_j^*] \}, \quad \rho \in M(N), \quad (2.3)$$

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A Brief History of the GKLS Equation



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Possibility of the total thermodynamic entropy production rate of a finite-sized isolated quantum system to be negative for the Gorini-Kossakowski-Sudarshan-Lindblad-type Markovian dynamics of its subsystem

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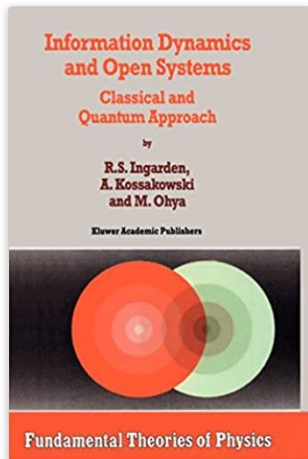
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We investigate a total thermodynamic entropy production rate of an isolated quantum system. In particular, we consider a quantum model of coupled harmonic oscillators in a star configuration, where a central harmonic oscillator (system) is coupled to a finite number of surrounding harmonic oscillators (bath). In this model, when the initial state of the total system is given by the tensor product of the Gibbs states of the system and the bath, every harmonic oscillator is always in a Gibbs state with a time-dependent temperature. This enables us to define time-dependent thermodynamic entropy for each harmonic oscillator and total nonequilibrium thermodynamic entropy as the summation of them. We analytically confirm that the total thermodynamic entropy satisfies the third law of thermodynamics. Our numerical solutions show that, even when the dynamics of the system is well approximated by the Gorini-Kossakowski-Sudarshan-Lindblad (GKSL)-type Markovian master equation, the total thermodynamic entropy production rate can be negative, while the total thermodynamic entropy satisfies the second law of thermodynamics. This result is a counterexample to the common belief that the total entropy production rate is non-negative when the system is under the GKSL-type Markovian dynamics.

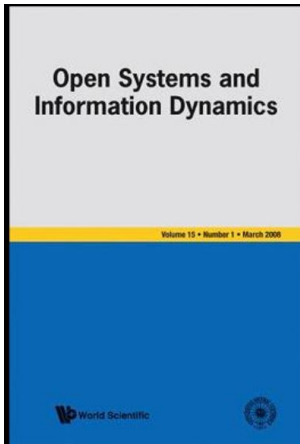
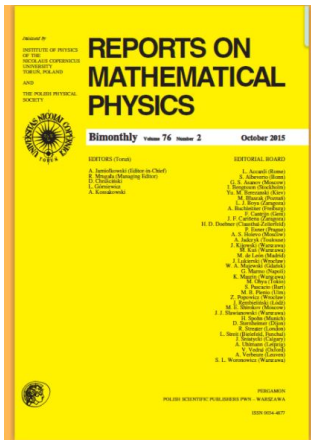
Ingarden-Kossakowski-Sudarshan-Gorini



Ingarden-Kossakowski-Ohya



ROMP & OSID





40 SMP (Toruń 2008)



40 SMP (Toruń 2008)





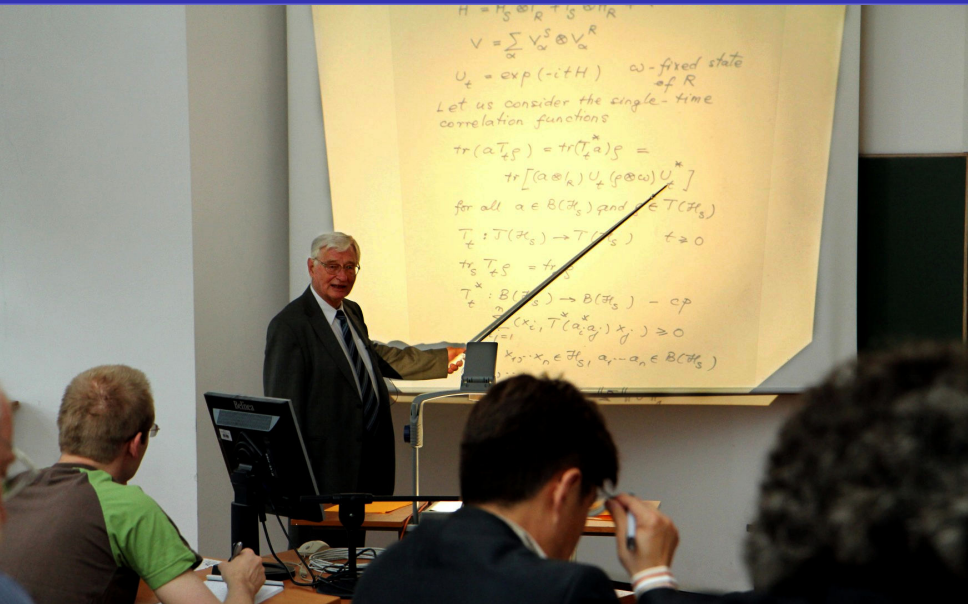








44 SMP (Toruń 2012)



$$H = H_S \otimes I_R + I_S \otimes H_R$$

$$V = \sum_{\alpha} V_{\omega}^S \otimes V_{\alpha}^R$$

$$U_t = \exp(-itH) \quad \omega\text{-fixed state of } R$$

Let us consider the single-time correlation functions

$$\text{tr}(a T_t^* \xi) = \text{tr}(T_t^* a) \xi =$$

$$\text{tr}[(a \otimes I_k) U_t(\rho \otimes \omega) U_t^*]$$

for all $a \in B(\mathcal{H}_S)$ and $\xi \in T(\mathcal{H}_S)$

$$T_t^* : T(\mathcal{H}_S) \rightarrow T(\mathcal{H}_S) \quad t \geq 0$$

$$\text{tr}_S T_t^* \xi = \text{tr}_S \xi$$

$$T_t^* : B(\mathcal{H}_S) \rightarrow B(\mathcal{H}_S) \quad \text{cp}$$

$$\sum_{i=1}^n \langle x_i, T_t^*(a_i^* a_j) x_j \rangle \geq 0$$

$$x_i, x_n \in \mathcal{H}_S, a_i, \dots, a_n \in B(\mathcal{H}_S)$$

40 years after (Toruń 2016)



40 years after (Toruń 2016)





Polish Nobel Prize

PROF. ANDRZEJ KOSSAKOWSKI – FNP PRIZE
LAUREATE 2019







THE NOBEL PRIZE
IN LITERATURE
2018

THE NOBEL PRIZE
IN LITERATURE
2019

Illustrations: Niklas Elmehed



Olga Tokarczuk

“for a narrative imagination that with encyclopedic passion represents the crossing of boundaries as a form of life.”

Peter Handke

“for an influential work that with linguistic ingenuity has explored the periphery and the specificity of human experience.”

THE SWEDISH ACADEMY

Olga Tokarczuk "The Books of Jacob"



THE BOOKS OF JACOB,

OR:

A FANTASTIC JOURNEY
ACROSS SEVEN BORDERS,
FIVE LANGUAGES
AND THREE MAJOR RELIGIONS,
NOT COUNTING THE MINOR SECTS.

OLGA TOKARCZUK
TRANSLATED BY JENNIFER CROFT

TOLD BY THE DEAD,
SUPPLEMENTED BY THE AUTHOR,
DRAWING FROM A RANGE OF
BOOKS

AND AIDED BY
IMAGINATION, THE WHICH BEING
THE GREATEST NATURAL GIFT
OF ANY PERSON.

FITZCARRALDO EDITION

Honorary citizen of Toruń (2020)





P. A. M. DIRAC

Directions in Physics

**THE DEVELOPMENT OF QUANTUM MECHANICS
QUANTUM ELECTRODYNAMICS
MAGNETIC MONOPOLES
A POSITIVE-ENERGY RELATIVISTIC WAVE EQUATION
COSMOLOGY AND THE GRAVITATIONAL CONSTANT**

from "Directions in Physics" by Dirac (1978)

It was a game, a very interesting game one could play. Whenever one solved one of the little problems, one could write a paper about it. It was very easy in those days for any second-rate physicist to do first-rate work. There has not been such a glorious time since then. It is very difficult now for a first-rate physicist to do second-rate work.

Andrzej Kossakowski managed to do the first-rate work

