

Quantum frequency estimation with local decoherence of arbitrary type

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51st Symposium on Mathematical Physics, Toruń, 18.06.2019

**Quantum frequency
estimation**

with

**local decoherence
of arbitrary type**

**Quantum frequency
estimation**



Quantum metrology

with



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Open-system dynamics

**Quantum frequency
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**Quantum Fisher
Information (QFI)**

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**Time-local Master
Equations (TLME)**

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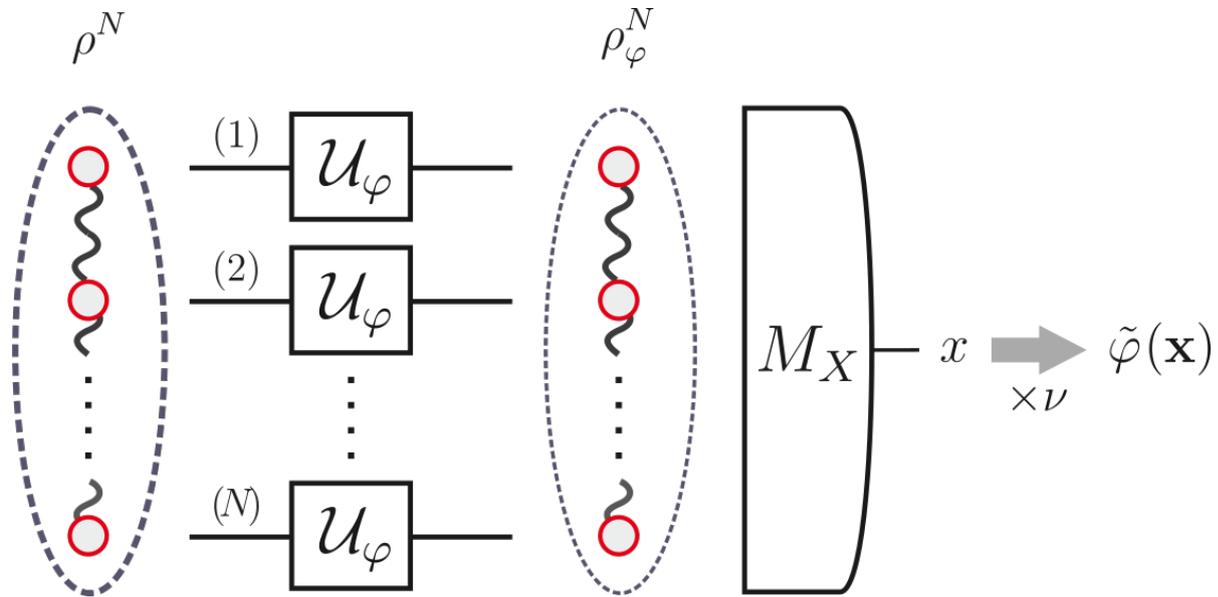
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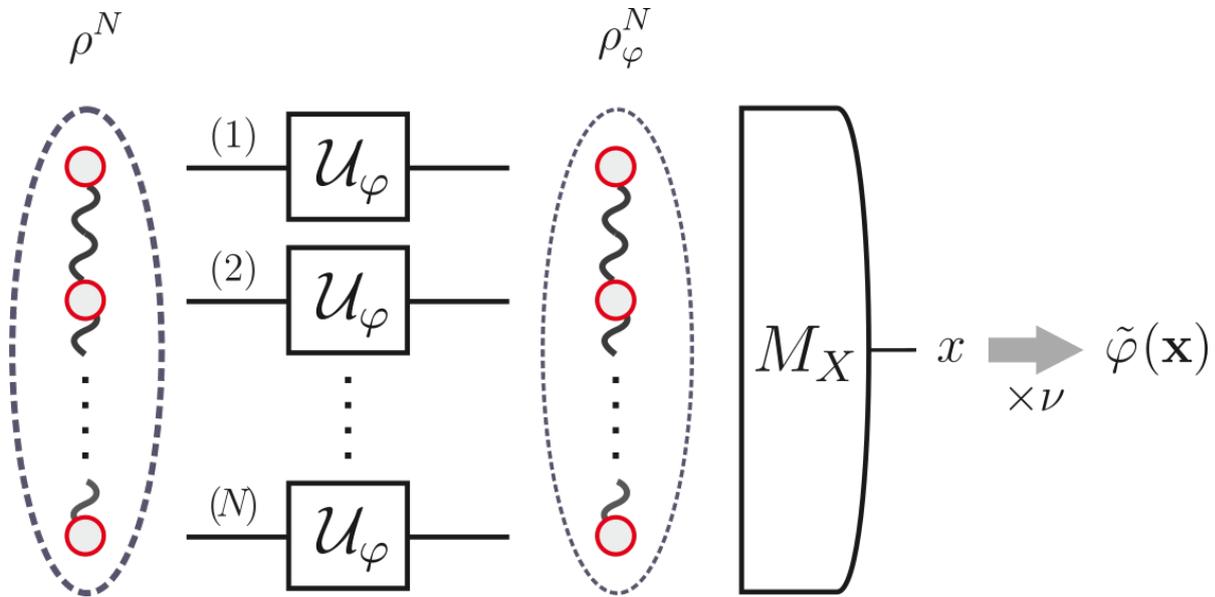
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Quantum Metrology

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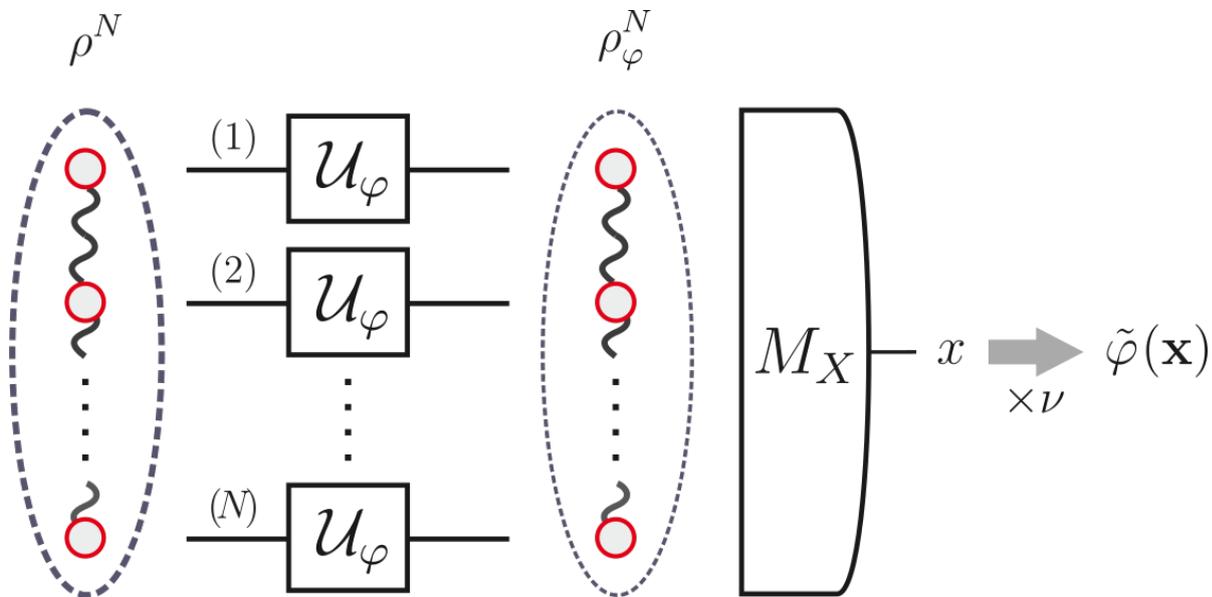
Quantum Metrology



minimise the
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$$\Delta^2 \tilde{\varphi} = \langle (\tilde{\varphi} - \varphi)^2 \rangle$$

Quantum Metrology



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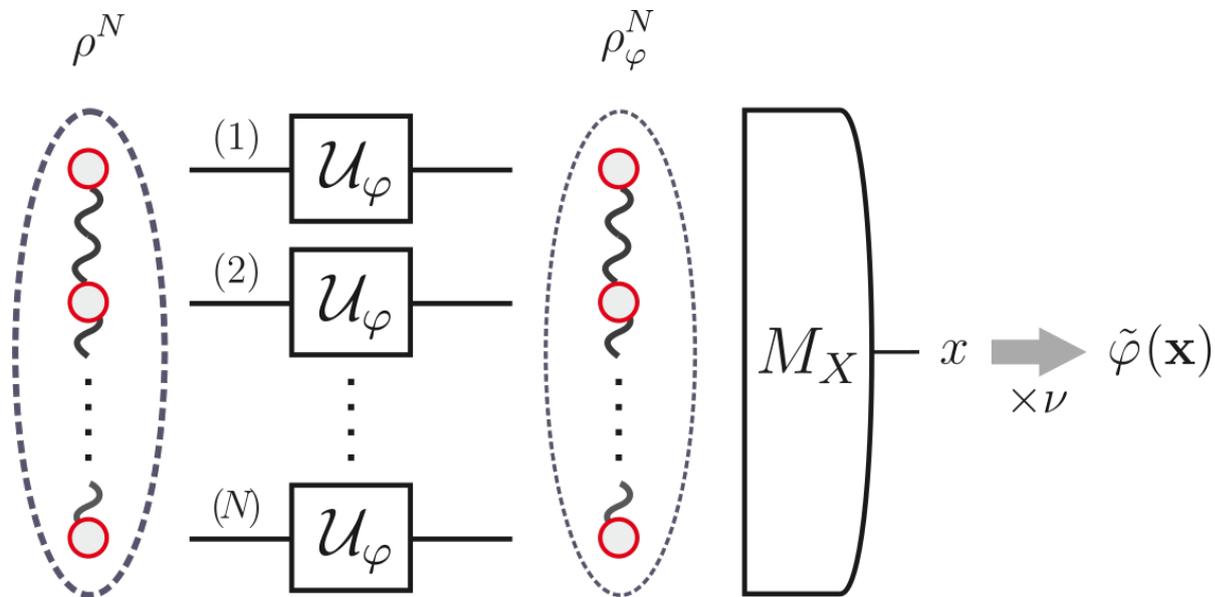
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Unitary encoding of the parameter in each probe:

$$\mathcal{U}_\varphi[\rho] = U_\varphi \rho U_\varphi^\dagger \quad \text{with} \quad U_\varphi = e^{-i\hat{h}\varphi}$$

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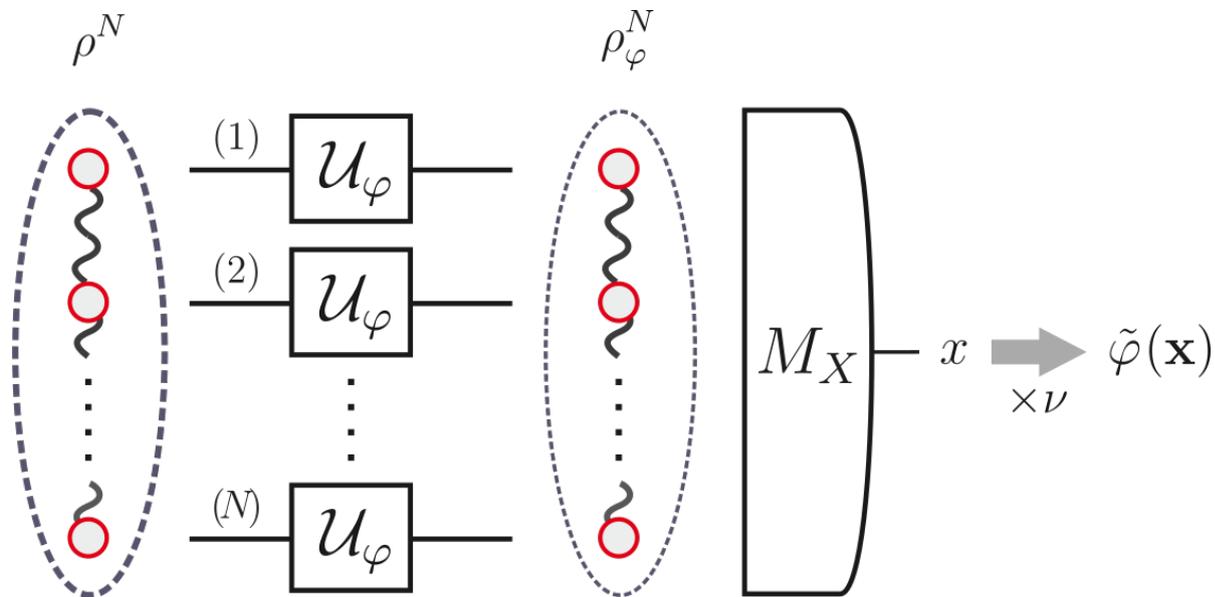
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Ultimate bound on the sensitivity to small fluctuations of the parameter (in the $\nu \rightarrow \infty$ limit):

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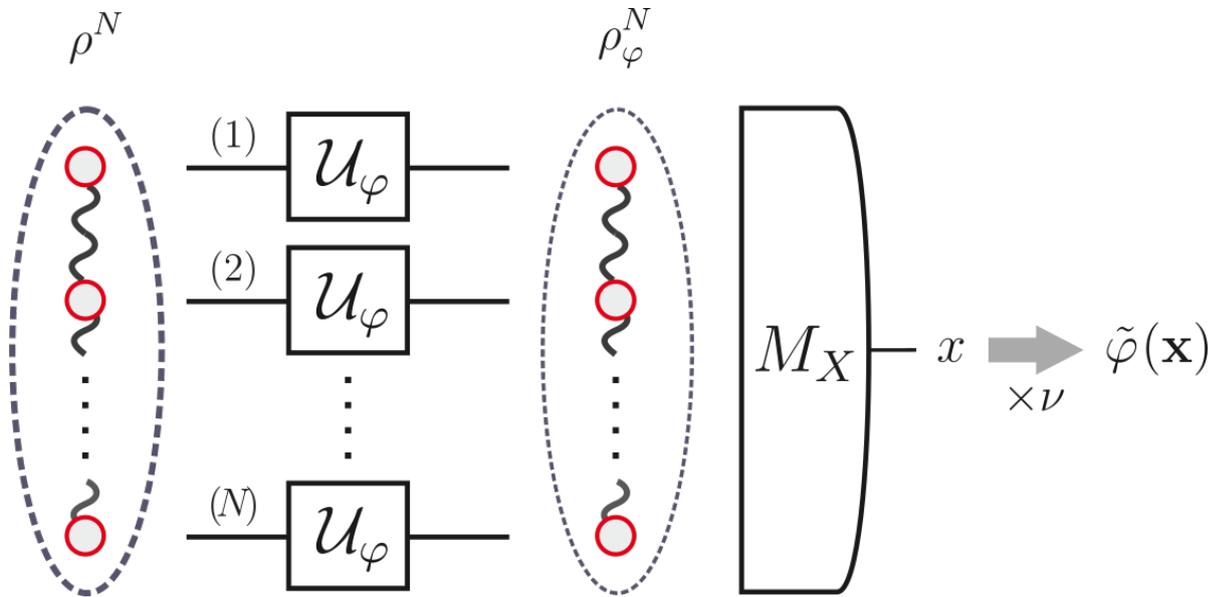
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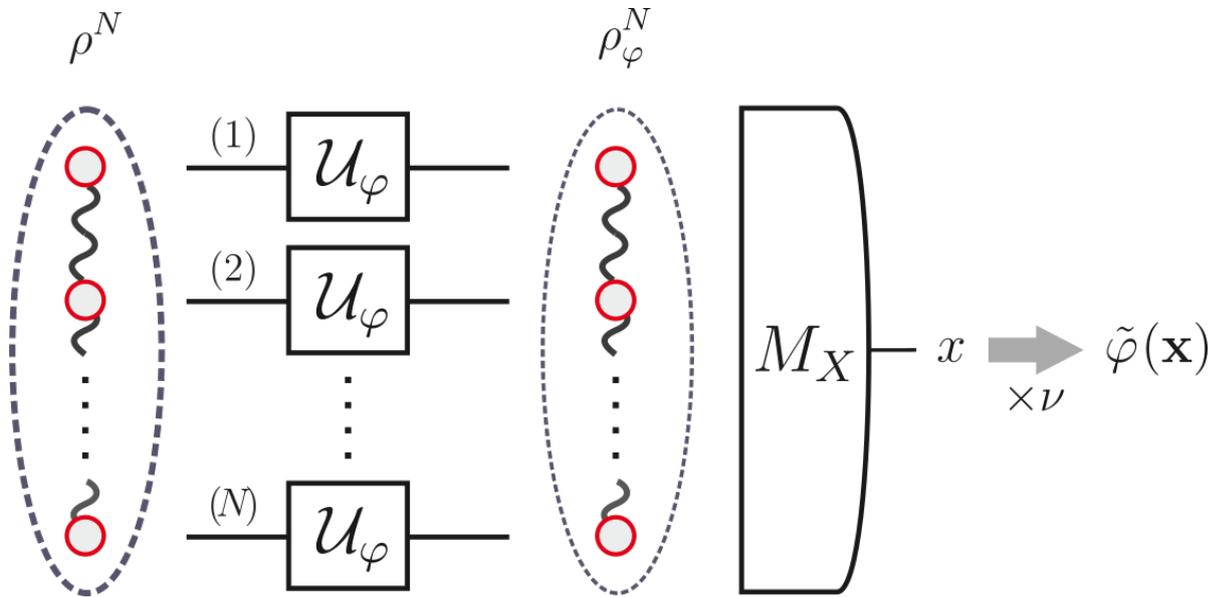
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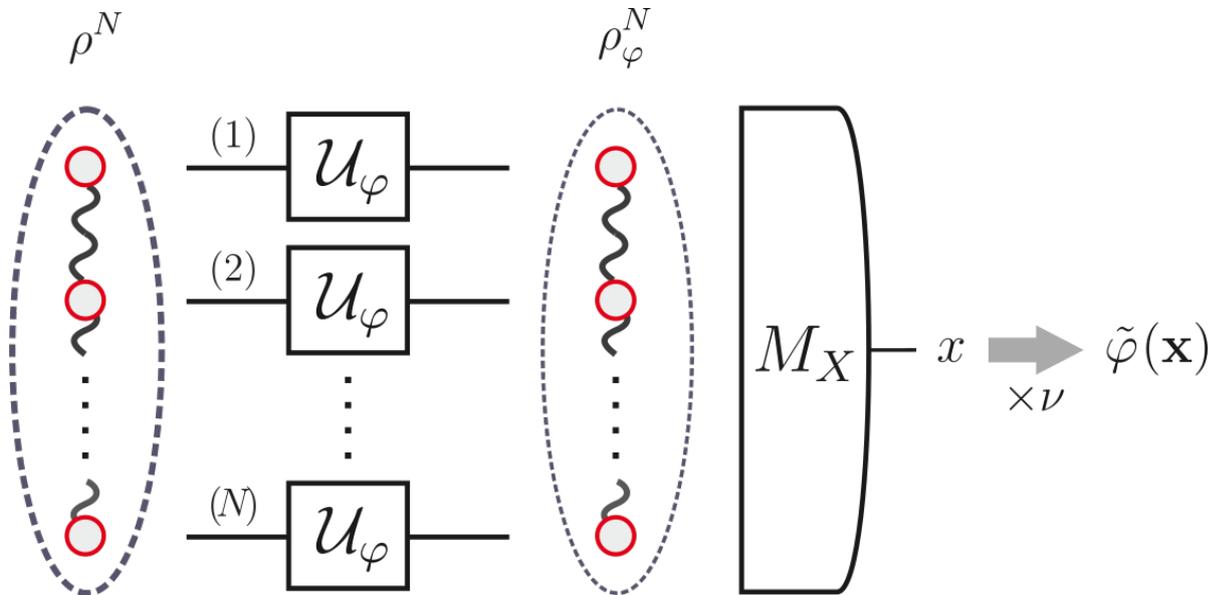
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Quantum Metrology



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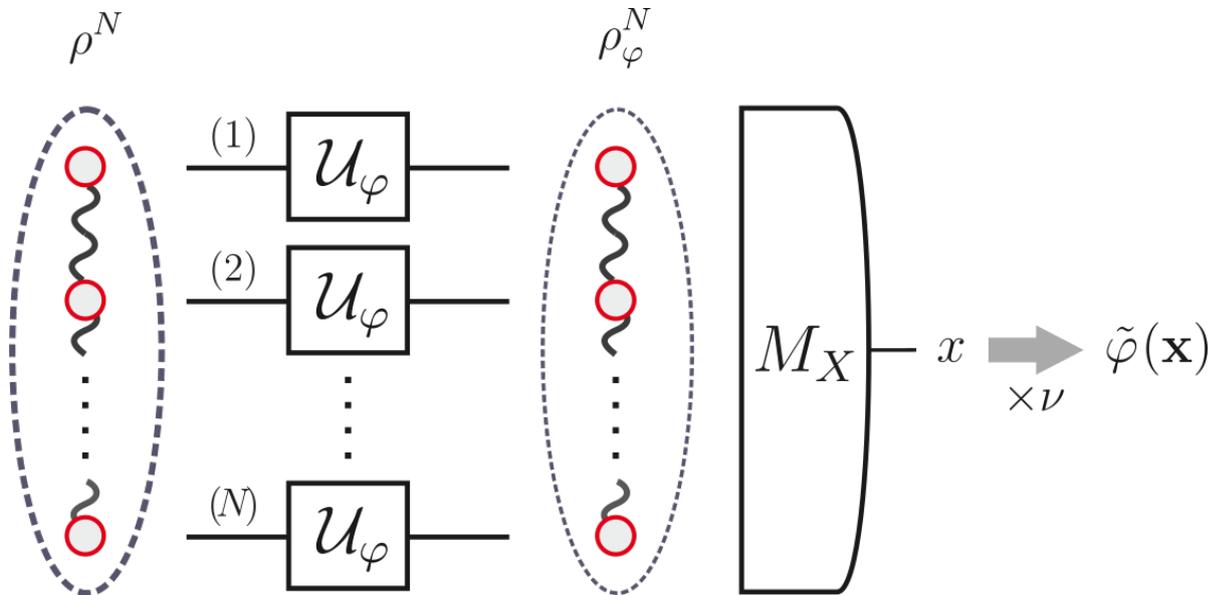
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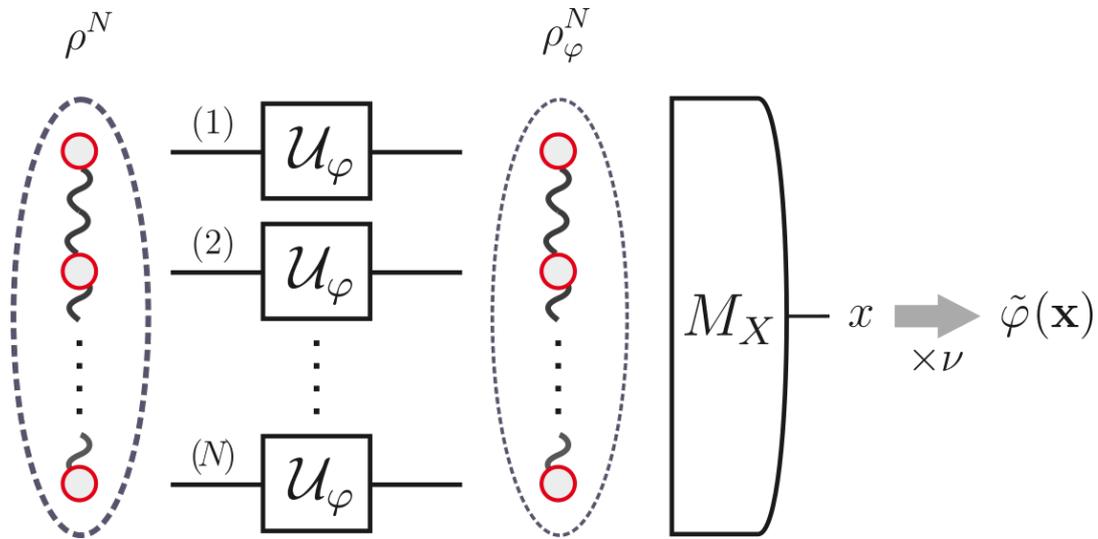
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- **Local (frequentist) estimation** with sufficiently large statistics (in contrast to the *Bayesian approach*).
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- **Parameter-value independence of QFI** due to unitary encoding. **Not true in general**.

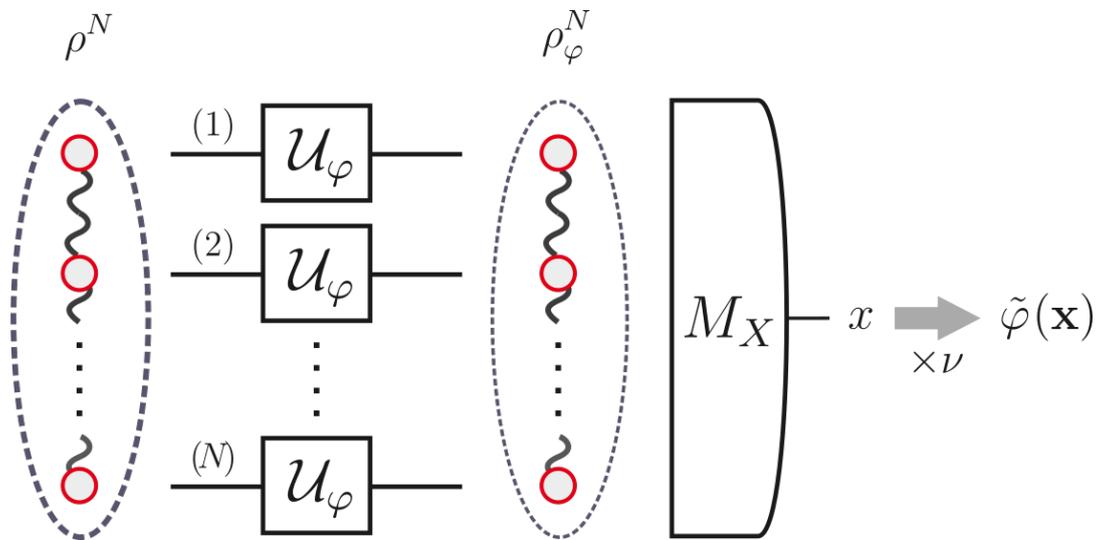
Quantum Metrology



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Quantum Metrology



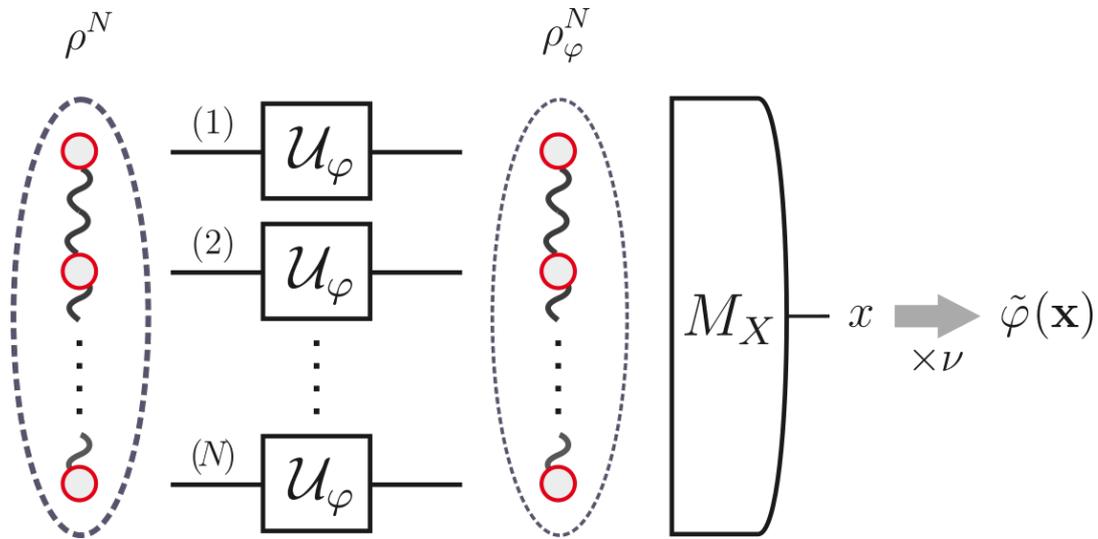
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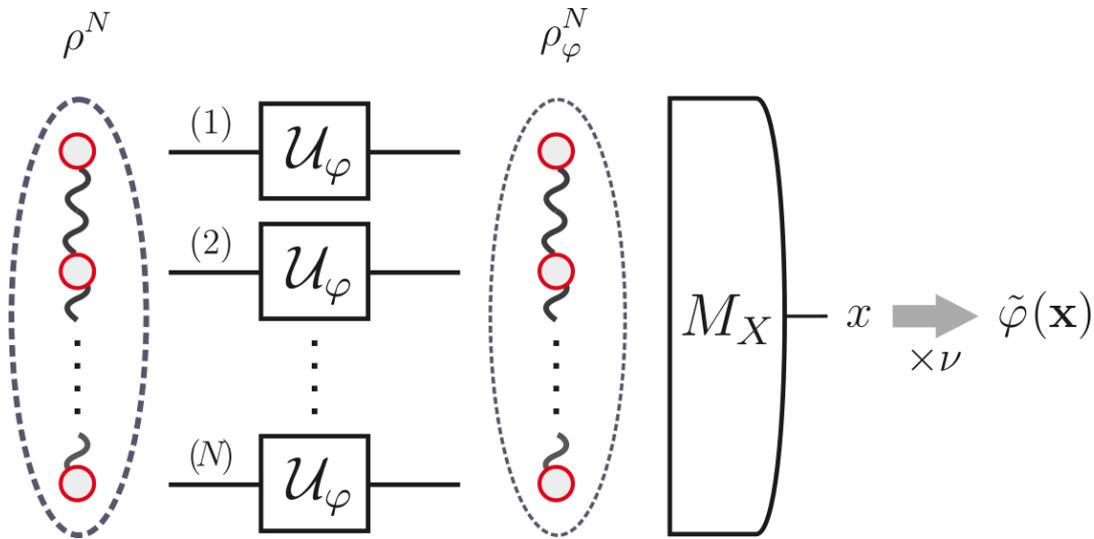
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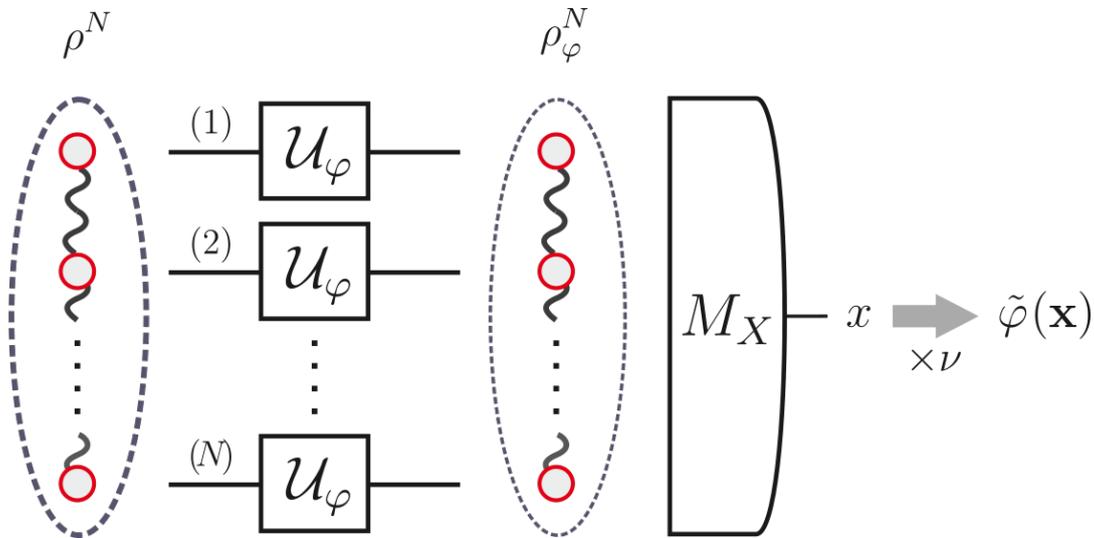
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$$\Delta^2 \hat{H}_{\text{sep}} \leq (s(\hat{H})^2/4) N \quad \Delta^2 \hat{H}_{\text{ent}} \leq (s(\hat{H})^2/4) N^2$$

Quantum Metrology



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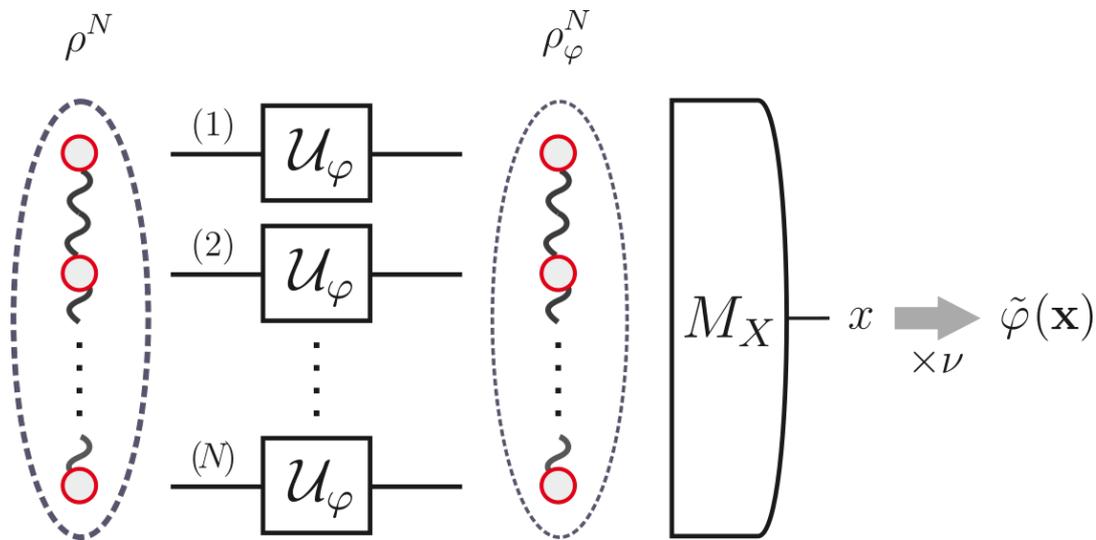
For **separable states** precision is bounded by the **Standard Quantum Limit**:

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Quantum Metrology



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For **entangled states** precision is bounded by the **Heisenberg Limit**:

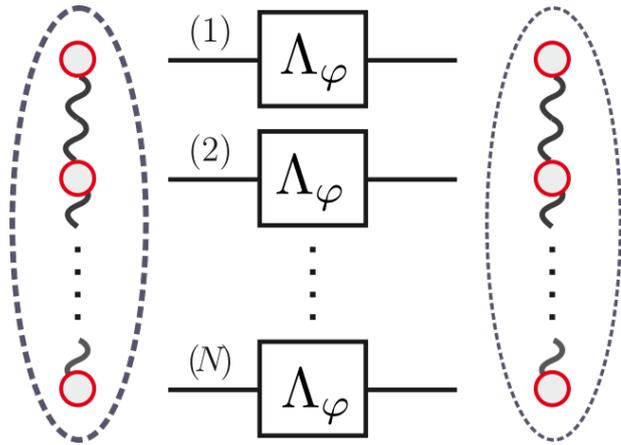
$$F_Q[\psi_{\text{GHZ}}^N] = N^2 \quad \Rightarrow \quad \Delta^2 \tilde{\varphi} \geq \frac{1}{N^2} \quad \text{HL}$$

GHZ (or NOON) state: $|\psi_{\text{GHZ}}^N\rangle := \frac{1}{\sqrt{2}}(|0^N\rangle + |1^N\rangle)$

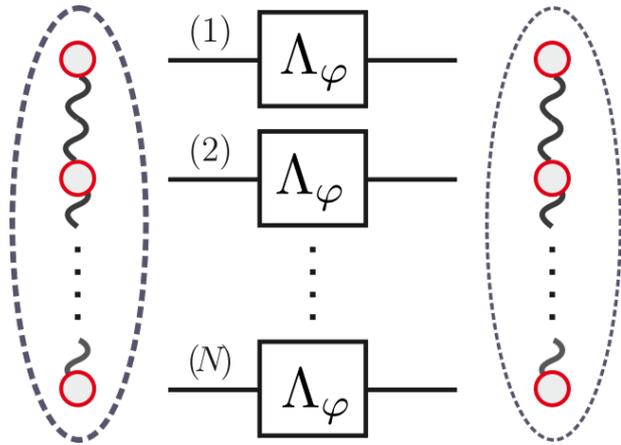
Noisy Quantum Metrology

[Akio Fujiwara and Hiroshi Imai, **J. Phys. A: Math. Theor.** **41** (2008)][Keiji
Matsumoto, **arXiv:1006.0300 [quant-ph]** (2010)]
[JK, Rafał Demkowicz-Dobrzański, **New J. Phys.** **15**, 073043 (2013)]

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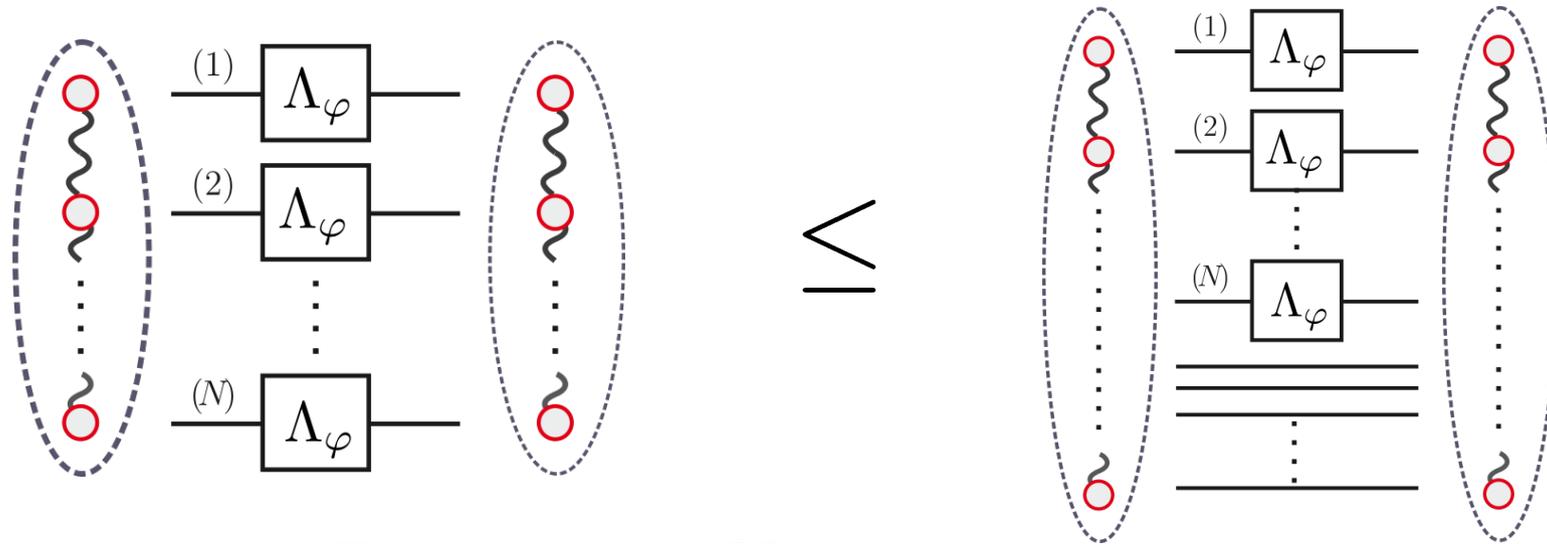


Noisy Quantum Metrology



The “**curse of SQL**”: $F_Q \leq \eta_{\text{noise}} N$

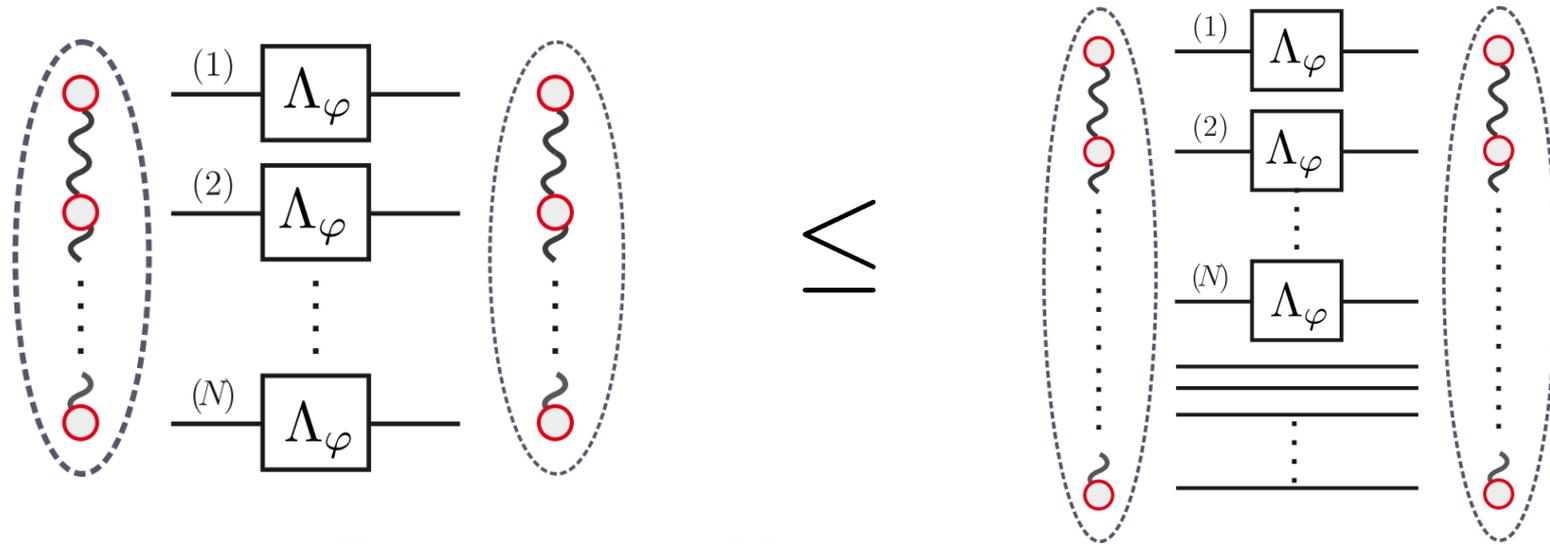
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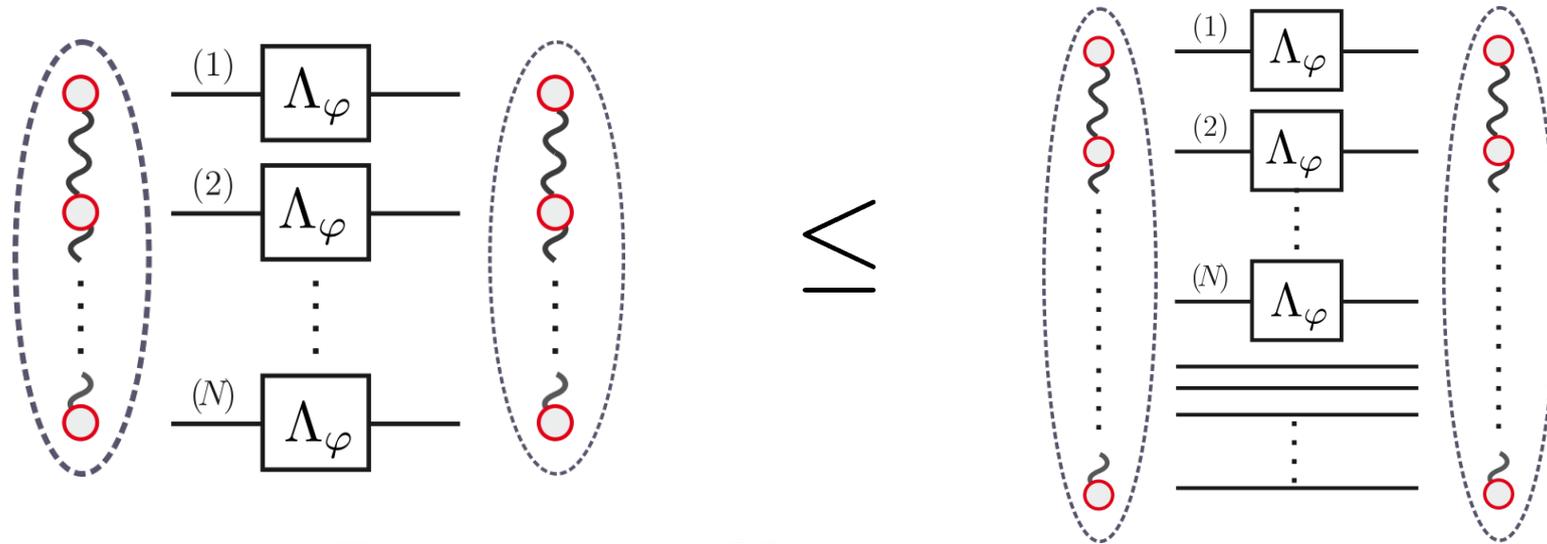


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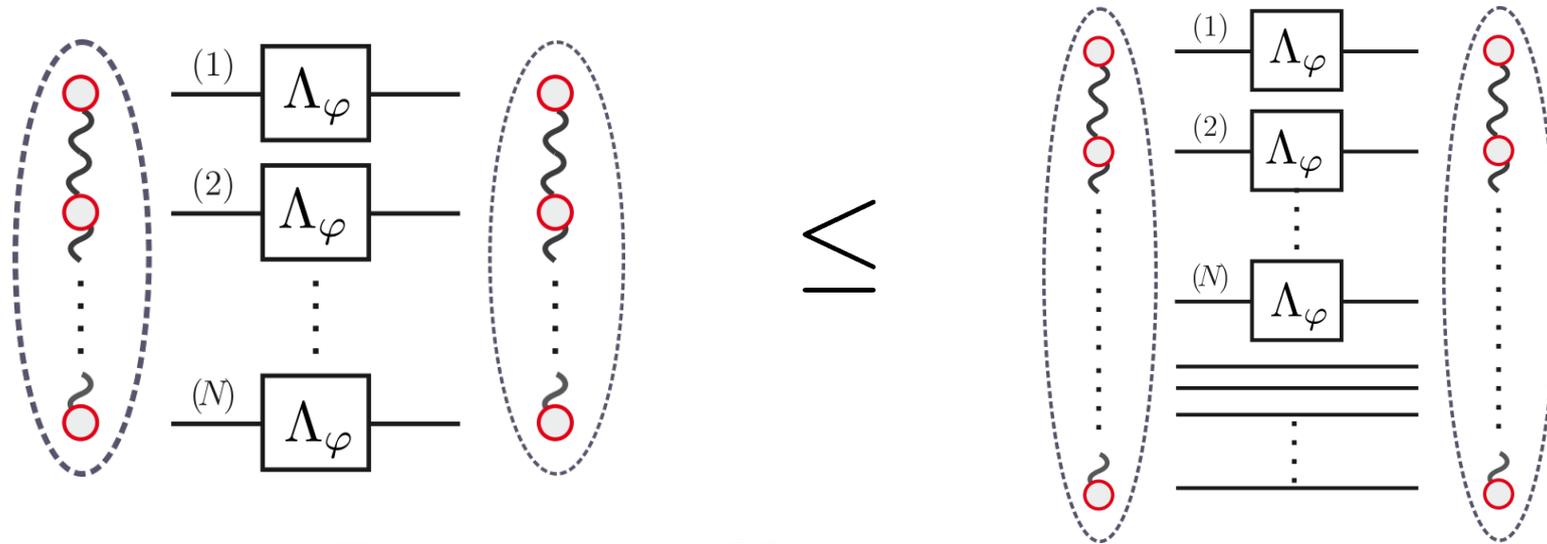
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Theorem: For almost all quantum channels (including φ -non-extremal and full-rank) one can find a Kraus representation such that:

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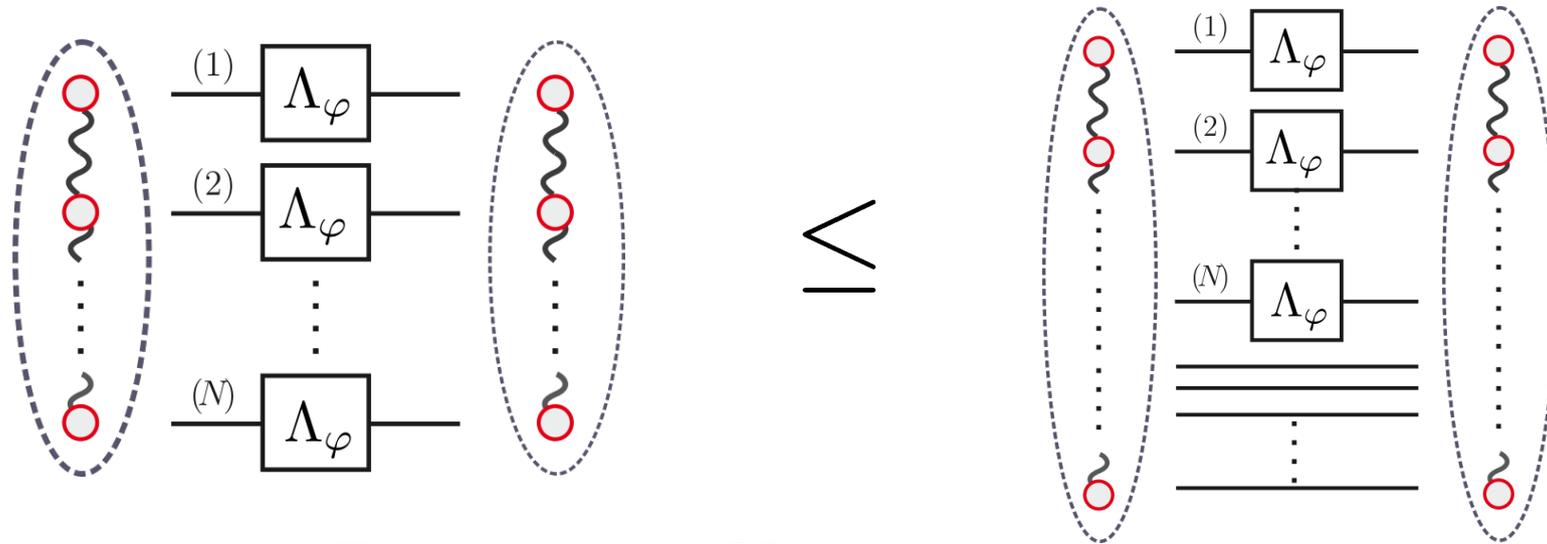
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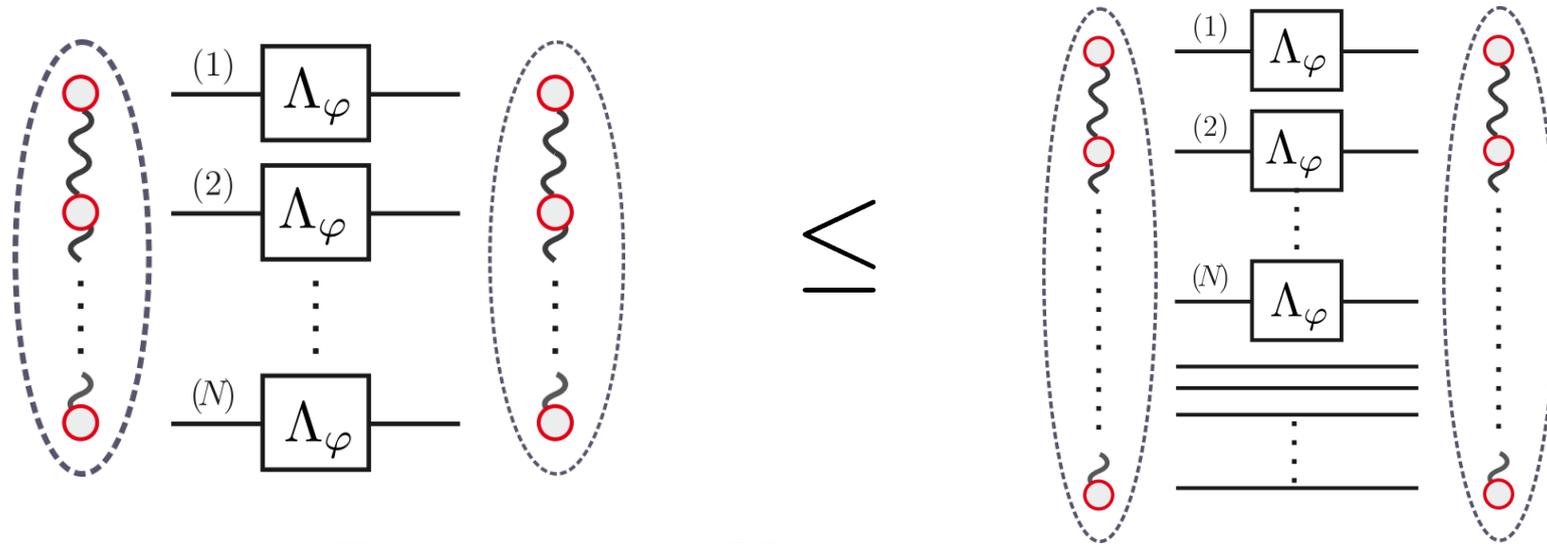
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[PHYSICAL MOTIVATION] Case of *noisy* unitary encoding:

What if you would like to split the channel into the **noiseless (unitary)** and **noisy (decoherence)** part?

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Theorem: For almost all quantum channels (including φ -non-extremal and full-rank) one can find a Kraus representation such that:

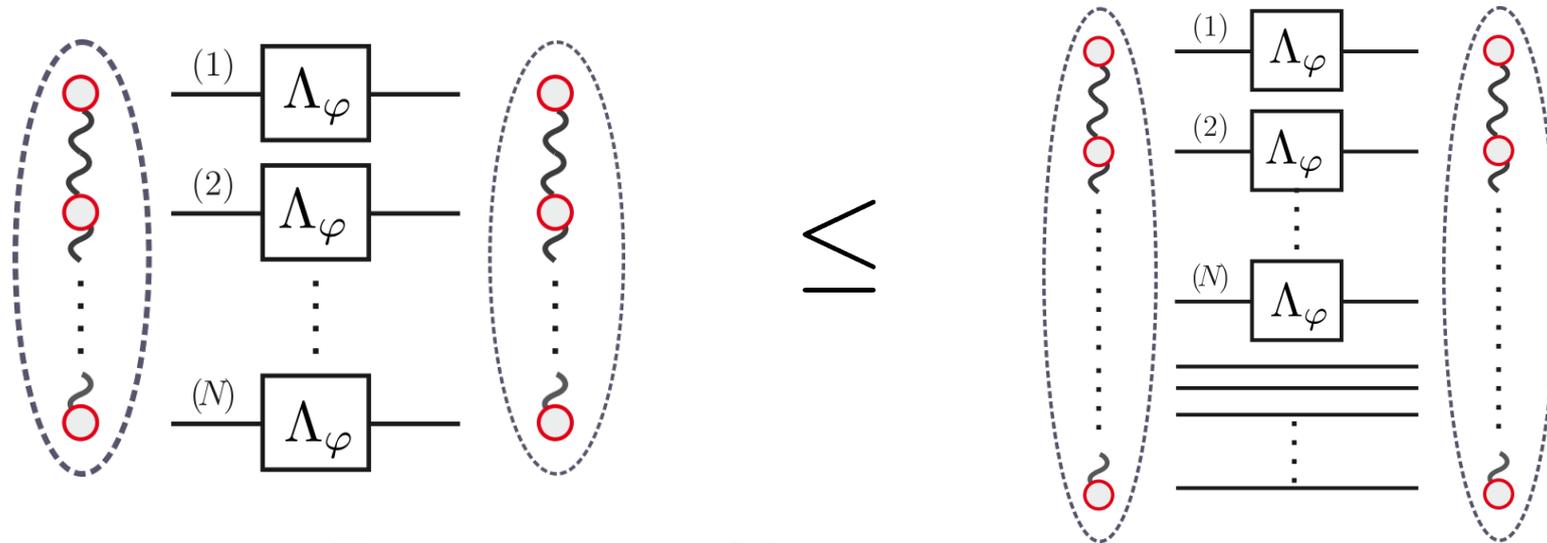
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[PHYSICAL MOTIVATION] Case of *noisy* unitary encoding:

What if you would like to split the channel into the **noiseless (unitary)** and **noisy (decoherence)** part?

$$\Lambda_\varphi = \Gamma_\eta \circ \mathcal{U}_\varphi = \mathcal{U}_\varphi \circ \Gamma_\eta$$

Noisy Quantum Metrology



The “**curse of SQL**”: $F_Q \leq \eta_{\text{noise}} N$

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Ok if commute, but in general need to **add time d.o.f.**

[Akio Fujiwara and Hiroshi Imai, *J. Phys. A: Math. Theor.* **41** (2008)][Keiji Matsumoto, *arXiv:1006.0300 [quant-ph]* (2010)]
 [JK, Rafał Demkowicz-Dobrzański, *New J. Phys.* **15**, 073043 (2013)]

Noisy Quantum Frequency Estimation

[*Huelga et al*, Phys. Rev. Lett. 79, 3865(1997)]

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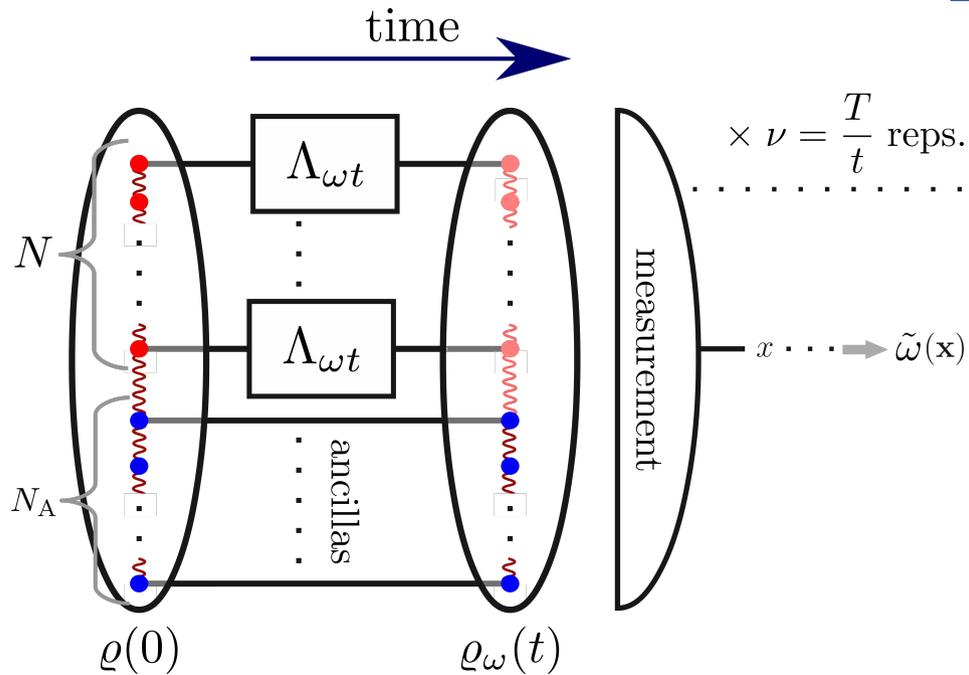
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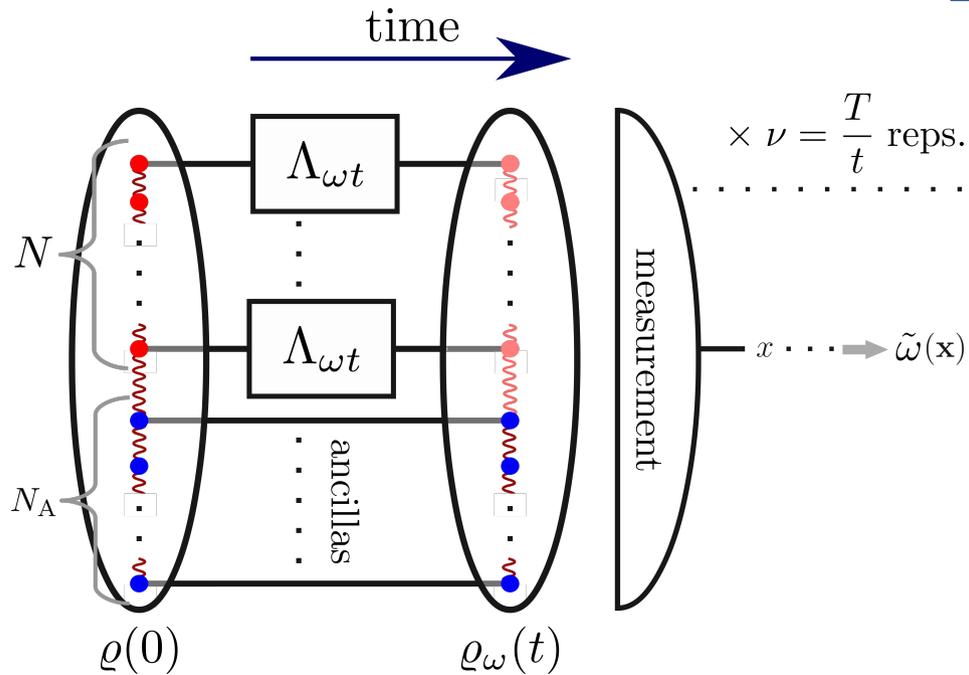
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$$\Delta^2 \tilde{\omega} \geq \frac{1}{\frac{T}{t} F_Q[\varrho_\omega(t)]}$$

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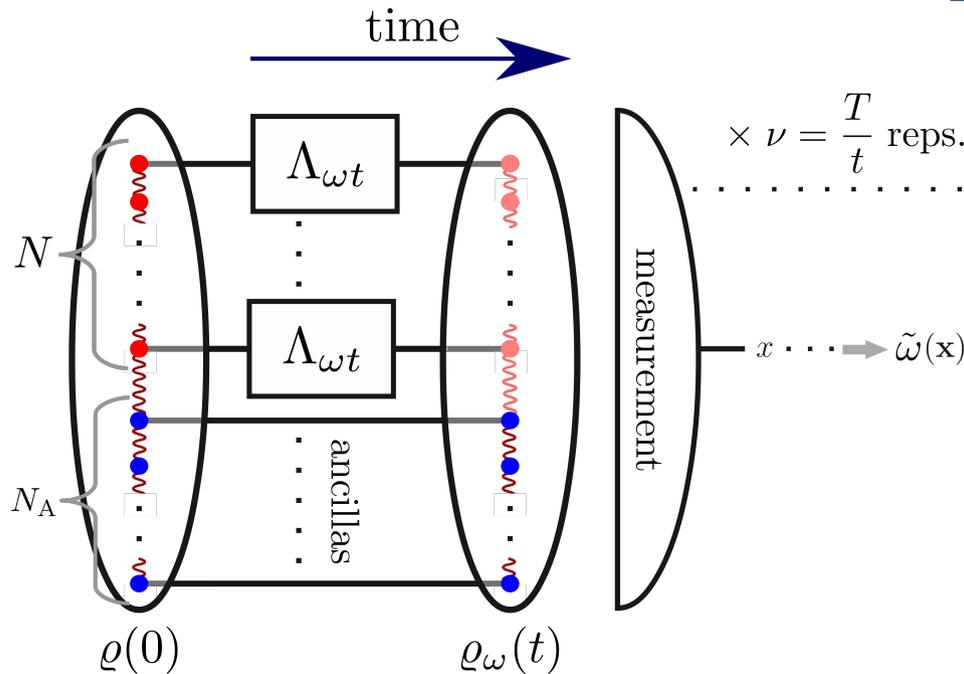
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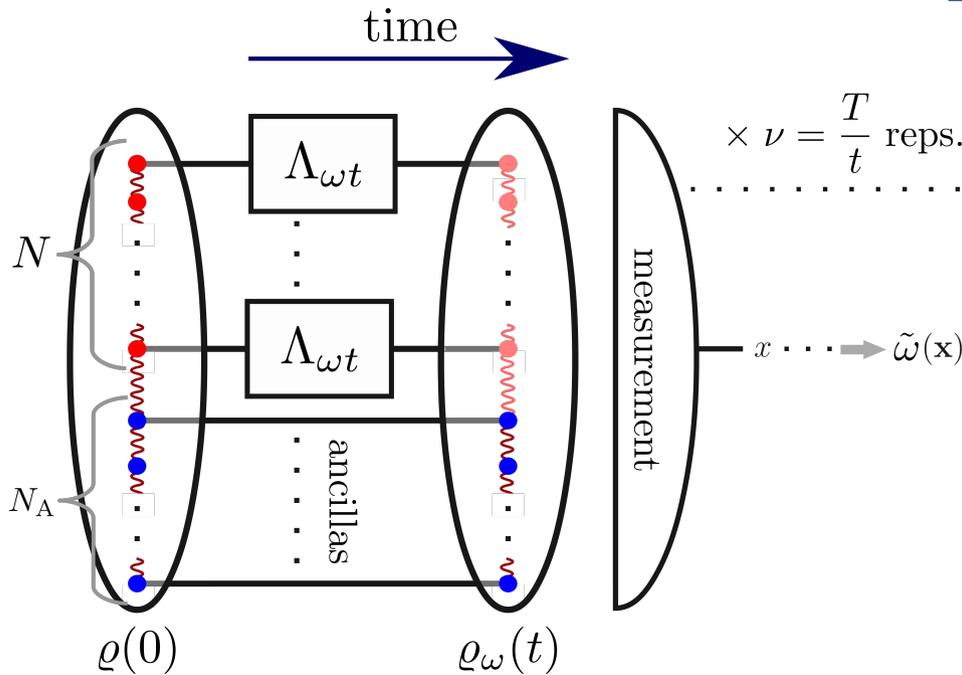
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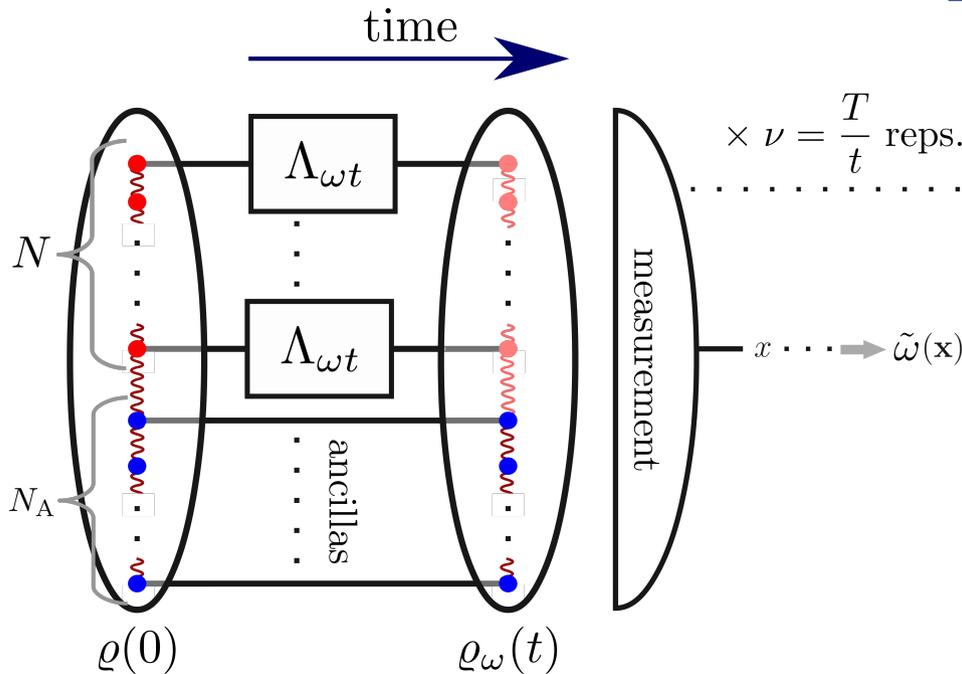
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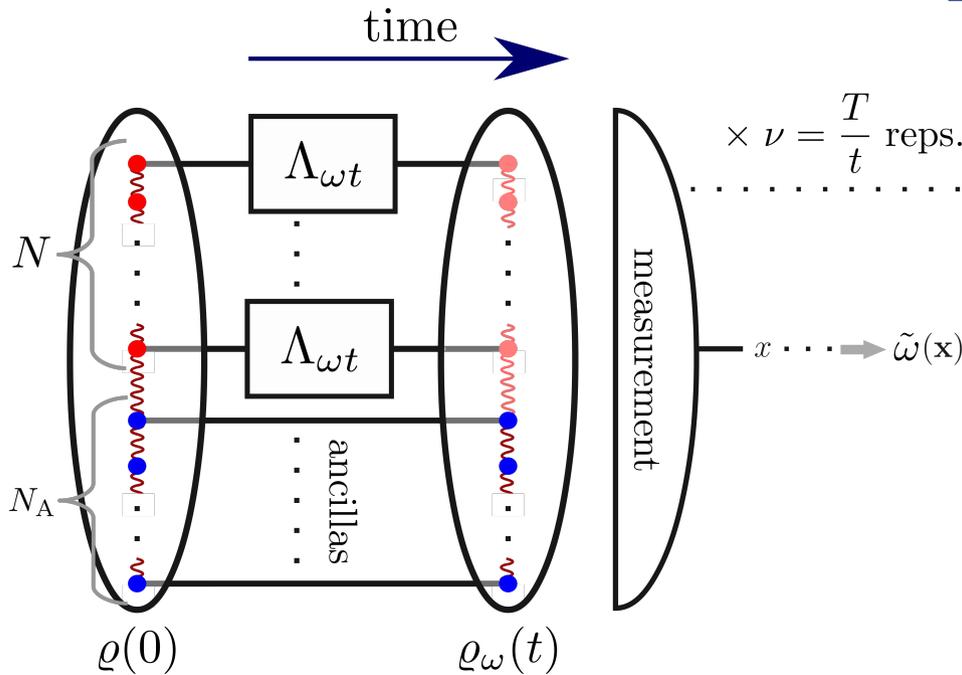
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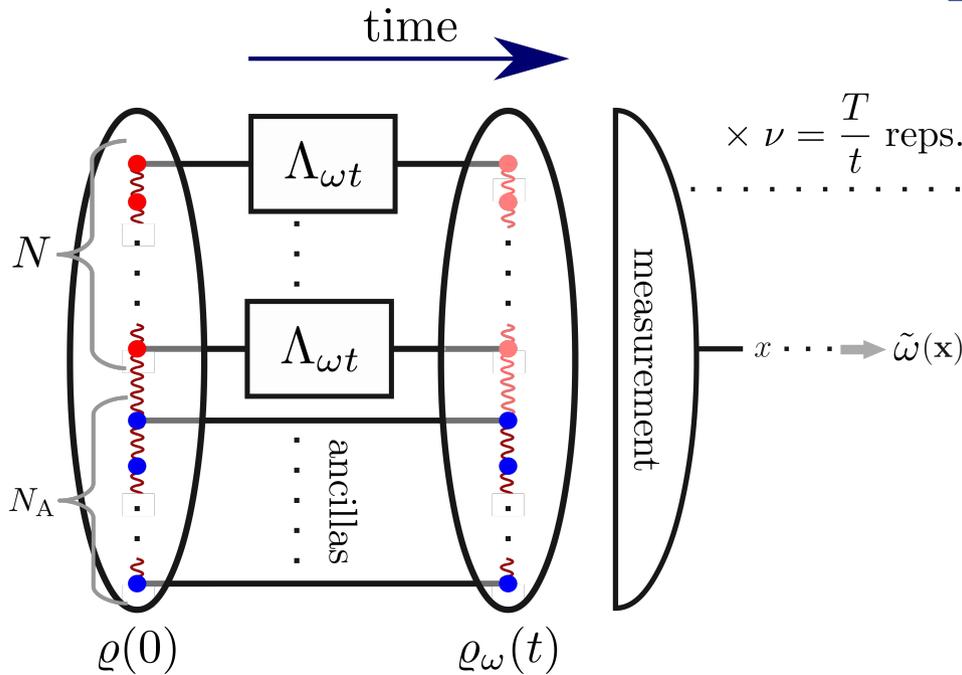
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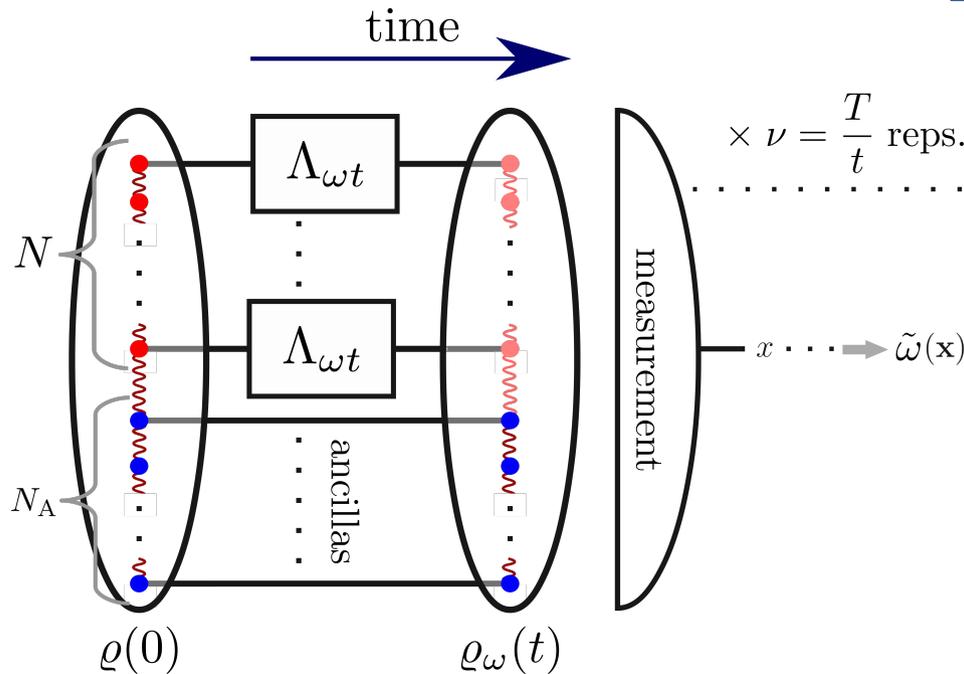
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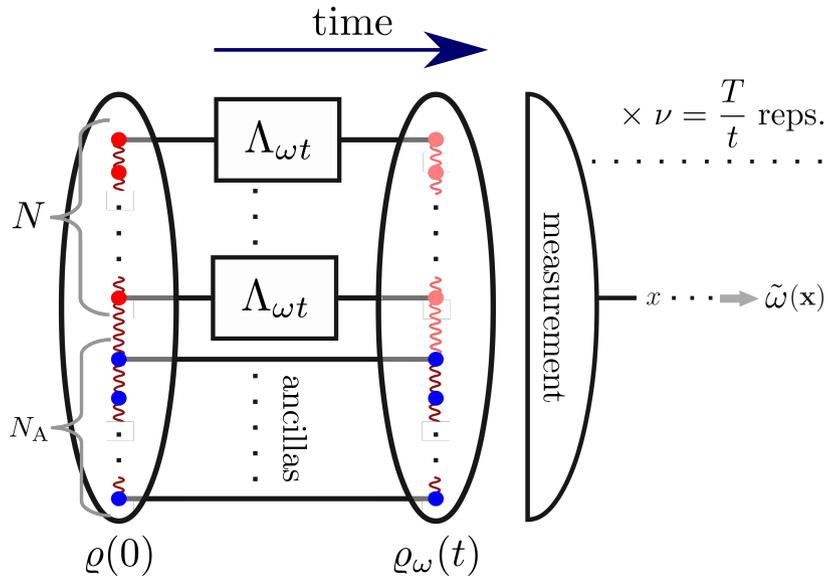
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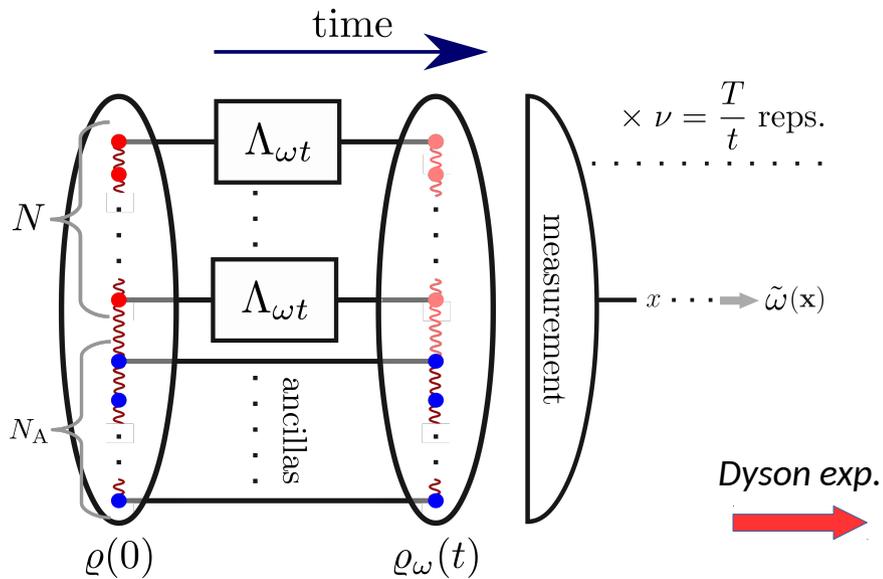
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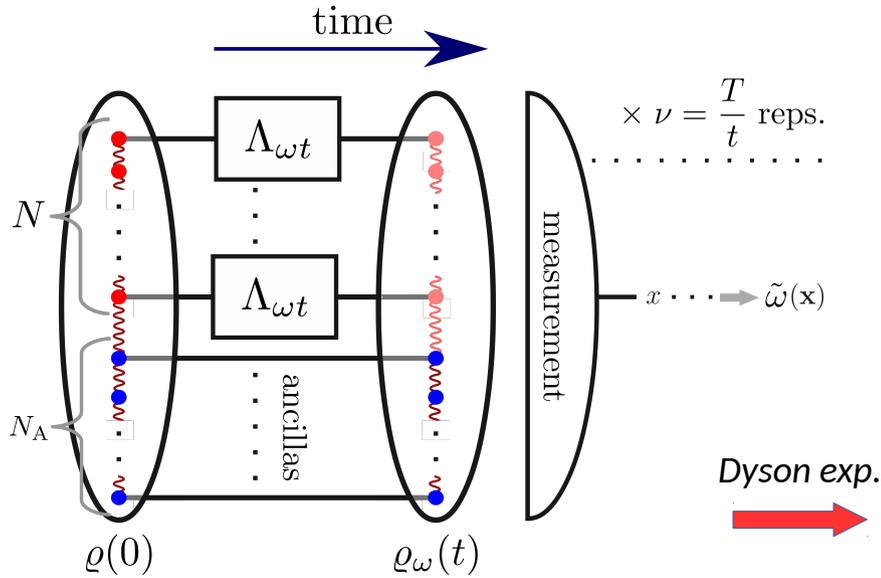
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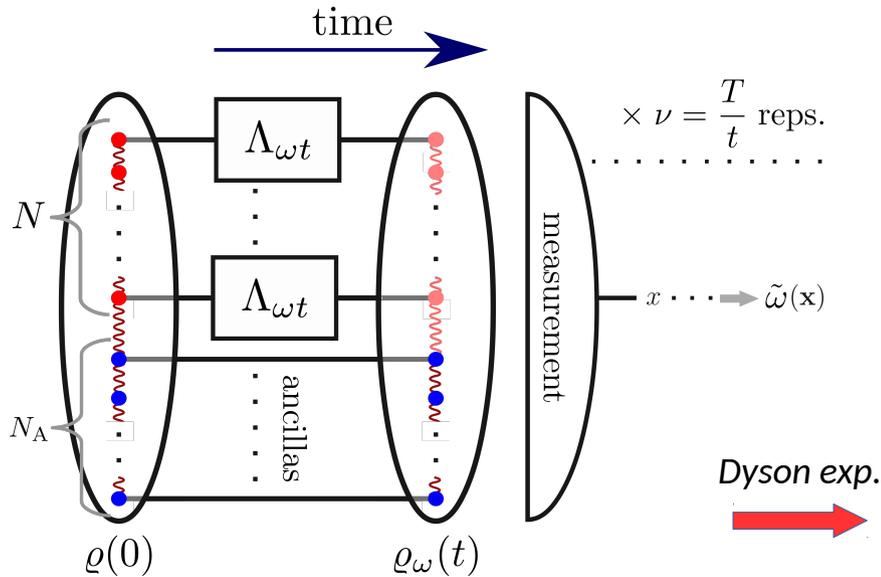
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$$\dot{\bullet} \equiv \frac{\partial}{\partial \omega} \bullet$$

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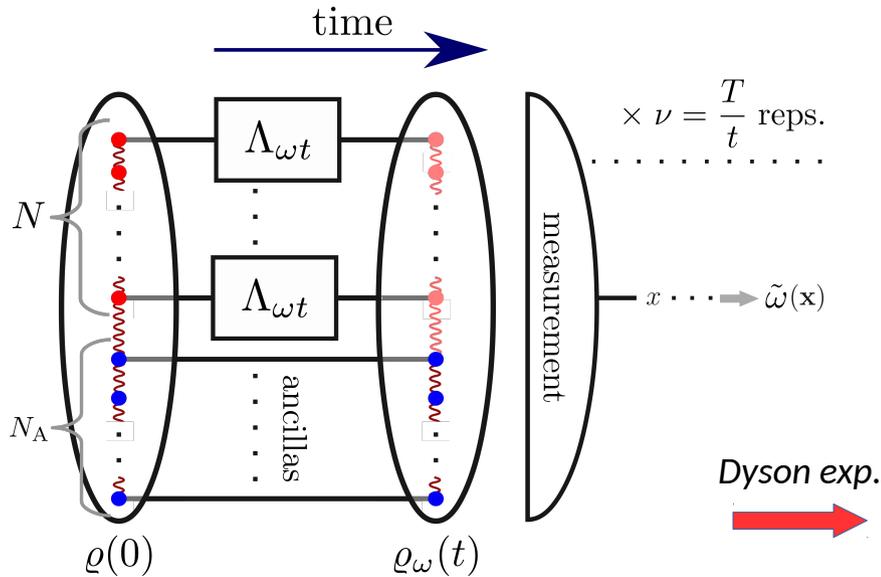
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But one should maximise precision over the round duration t , with the optimal $t_{opt}(N)$ determined now by noise (physics!).

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Quantum Cramer-Rao Bound

$$\Delta^2 \tilde{\omega} T \geq \frac{1}{f_{N,t}} \quad \text{with } f_{N,t} := \frac{F_Q[\varrho_\omega(t)]}{t}$$

QFI per round duration

$$\nu \gg 1 \implies T \gg t$$

$$\Lambda_{\omega t}[\rho] = \mathcal{T}_{\leftarrow} e^{\int_0^t d\tau \mathcal{L}_{\omega,\tau}[\rho]} = \sum_{i=0}^{r-1} K_i(t) \rho K_i^\dagger(t) = \mathbf{K} \rho \mathbf{K}^\dagger$$

Use the “channel” bound:

$$f_{N,t} \leq f_{N,t}^\uparrow := 4 \min_{\mathbf{h}(t)} \left\{ N \left\| \frac{\alpha_{\mathbf{h}}(t)}{t} \right\|^2 + N(N-1) \left\| \frac{\beta_{\mathbf{h}}(t)}{\sqrt{t}} \right\|^2 \right\} \quad \text{with}$$

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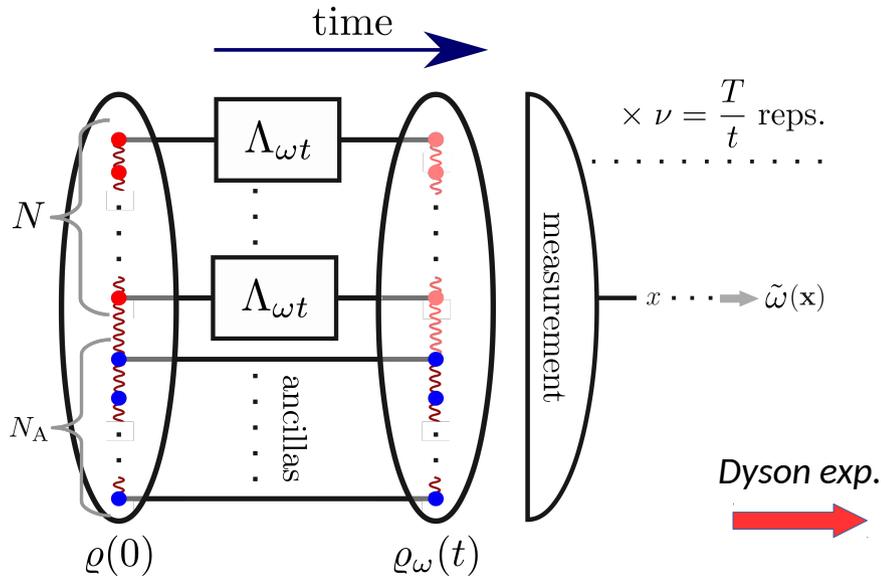
$$\dot{\bullet} \equiv \frac{\partial}{\partial \omega} \bullet$$

But one should maximise precision over the round duration t , with the optimal $t_{\text{opt}}(N)$ determined now by **noise (physics!)**.

$$f_{N,t} \leq f_N^\uparrow := \max_{0 \leq t \leq T} f_{N,t}^\uparrow \quad \text{and} \quad t_{\text{opt}}(N) \rightarrow 0, \text{ as } N \rightarrow \infty$$

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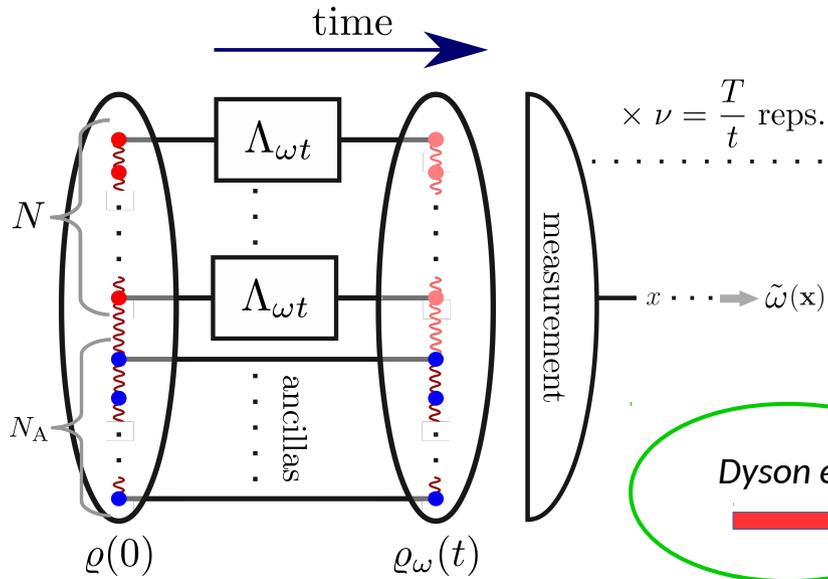
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Dyson exp.



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Catch...

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[JK, Rafał Demkowicz-Dobrzański, *New J. Phys.* **15**, 073043 (2013)]

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“Zeno regime” $\ell = 1$

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3) **perpendicular semigroup dephasing** - bound: *numeric*; “saturability” (GHZ states+parity): *analytic*.

[Chaves et al., ..., JK, ..., *Phys. Rev. Lett.* **111**, 120401 (2013)]

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[JK, Rafał Demkowicz-Dobrzański, *New J. Phys.* **15**, 073043 (2013)]

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[Smirne et al., ..., JK, ..., *Phys. Rev. Lett.* **116**, 120801 (2016)]

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[Haase et al., ..., JK, ..., *New J. Phys.* **20**, 053009 (2018)]

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$$f_{N,t} \leq \hat{f}_{N,t} := 4 \min_{\mathbf{h}(t)} \left\{ N \left\| \frac{\alpha_{\mathbf{h}}(t)}{t} \right\| + N(N-1) \left\| \frac{\beta_{\mathbf{h}}(t)}{\sqrt{t}} \right\|^2 \right\} \quad \begin{aligned} \alpha_{\mathbf{h}}(t) &= \dot{\tilde{\mathbf{K}}}^\dagger \dot{\tilde{\mathbf{K}}}, & \beta_{\mathbf{h}}(t) &= -i \dot{\tilde{\mathbf{K}}}^\dagger \tilde{\mathbf{K}}, \\ \tilde{\mathbf{K}} &= \mathbf{K}, & \dot{\tilde{\mathbf{K}}} &= \dot{\mathbf{K}} - i\mathbf{h}\mathbf{K}, \end{aligned}$$

(ignoring t-dependence)

$$\mathbf{h} = \left(\begin{array}{c|c} h_{00} & \vec{h}^\dagger \\ \hline \vec{h} & \mathfrak{h} \end{array} \right) \quad \Rightarrow$$

$$\begin{aligned} \alpha_{\mathbf{h}}(t) &= [h_{00}^2 + \vec{h}^\dagger \vec{h}] \mathbf{1} + 2[(h_{00} \vec{h}^\dagger + \vec{h}^\dagger \mathfrak{h}) \vec{J}^{(1/2)}]^\text{H} t^{\frac{1}{2}} + [2h_{00} H + 2(h_{00}^2 + \vec{h}^\dagger \vec{h})(J_0^{(1)})^\text{H} + \\ &\quad + (\vec{J}^{(1/2)})^\dagger (\vec{h} \vec{h}^\dagger + \mathfrak{h}^2) \vec{J}^{(1/2)} + 2[(h_{00} \vec{h}^\dagger + \vec{h}^\dagger \mathfrak{h}) \vec{J}^{(1)}]^\text{H} + 2i(\vec{h}^\dagger \dot{\vec{J}}^{(1)})^\text{AH}] t + O(t^{\frac{3}{2}}) \\ \beta_{\mathbf{h}}(t) &= h_{00} \mathbf{1} + 2(\vec{h}^\dagger \vec{J}^{(1/2)})^\text{H} t^{\frac{1}{2}} + \left[H + 2h_{00} (J_0^{(1)})^\text{H} + 2(\vec{h}^\dagger \vec{J}^{(1)})^\text{H} + (\vec{J}^{(1/2)})^\dagger \mathfrak{h} \vec{J}^{(1/2)} \right] t \\ &\quad + 2\left[((J_0^{(1)})^\dagger (\vec{h}^\dagger \vec{J}^{(1/2)}))^\text{H} + i\omega (H (\vec{h}^\dagger \vec{J}^{(1/2)}))^\text{AH} + ((\vec{J}^{(1)})^\dagger \mathfrak{h} \vec{J}^{(1/2)})^\text{H} + (\vec{h}^\dagger \vec{J}^{(3/2)})^\text{H} \right] t^{\frac{3}{2}} \\ &\quad + \left[\omega^2 h_{00} H^2 + h_{00} (J_0^{(1)})^\dagger J_0^{(1)} + 2h_{00} (J_0^{(2)})^\text{H} + 2i h_{00} \omega (H J_0^{(1)})^\text{AH} + \right. \\ &\quad \left. + 2((\vec{J}^{(1)})^\dagger \vec{h}) J_0^{(1)} \right]^\text{H} + 2i\omega (H (\vec{h}^\dagger \vec{J}^{(1)}))^\text{AH} + (\vec{J}^{(1)})^\dagger \mathfrak{h} \vec{J}^{(1)} + \\ &\quad \left. + 2(\vec{h}^\dagger \vec{J}^{(2)})^\text{H} - i((\dot{\vec{J}}^{(3/2)})^\dagger \vec{J}^{(1/2)})^\text{AH} + 2((\vec{J}^{(3/2)})^\dagger \mathfrak{h} \vec{J}^{(1/2)})^\text{H} \right] t^2 + O(t^{\frac{5}{2}}), \end{aligned}$$

3) Find term in the “beta”-expansion such that you can cancel the Ht term and:

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We provide a general (“step by step”) recipe how to construct a meaningful bound for a given dynamics.

Summary of the procedure:

1) With help of the Dyson expansion construct Taylor-expansions of Kraus operators up to sufficient order:

$$\mathbf{K}(t) = \begin{bmatrix} K_0(t) \\ \vec{K}(t) \end{bmatrix} \Rightarrow \begin{aligned} K_0(t) &= -\left(i\omega H - J_0^{(1)}\right) t + \sum_{\ell=2,3,\dots} J_0^{(\ell)} t^\ell \Rightarrow \dot{K}_0(t) = -iHt + \sum_{\ell=2,3,\dots} \dot{J}_0^{(\ell)} t^\ell. \\ \vec{K}(t) &= \sum_{\ell=\frac{1}{2},1,\frac{3}{2},\dots} \vec{J}^{(\ell)} t^\ell \Rightarrow \dot{\vec{K}}(t) = \sum_{\ell=\frac{1}{2},\frac{3}{2},\dots} \dot{\vec{J}}^{(\ell)} t^\ell, \end{aligned}$$

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$$f_{N,t} \leq \hat{f}_{N,t} := 4 \min_{\mathbf{h}(t)} \left\{ N \left\| \frac{\alpha_{\mathbf{h}}(t)}{t} \right\| + N(N-1) \left\| \frac{\beta_{\mathbf{h}}(t)}{\sqrt{t}} \right\|^2 \right\} \quad \begin{aligned} \alpha_{\mathbf{h}}(t) &= \dot{\tilde{\mathbf{K}}}^\dagger \dot{\tilde{\mathbf{K}}}, & \beta_{\mathbf{h}}(t) &= -i \dot{\tilde{\mathbf{K}}}^\dagger \tilde{\mathbf{K}}, \\ \tilde{\mathbf{K}} &= \mathbf{K}, & \dot{\tilde{\mathbf{K}}} &= \dot{\mathbf{K}} - i\mathbf{h}\mathbf{K}, \end{aligned}$$

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$$\mathbf{h} = \left(\begin{array}{c|c} h_{00} & \vec{h}^\dagger \\ \hline \vec{h} & \mathfrak{h} \end{array} \right) \Rightarrow$$

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$$\beta_{\mathbf{h}}(t) = b_0 t + b_1 \chi t^b + O(t^{\max\{\frac{3}{2}, b + \frac{1}{2}\}})$$

Noisy Quantum Frequency Estimation

Can one provide an analytic derivation for general time-local dissipator with **analytic rates and Lindblad operators**?

$$\mathcal{D}_t[\rho(t)] = \sum_{j=1}^{l \leq r-1} \gamma_j(t) \left(L_j(t) \rho(t) L_j^\dagger(t) - \frac{1}{2} \left\{ \rho(t), L_j^\dagger(t) L_j(t) \right\} \right) \text{ with } \gamma_j(t) = \sum_{\ell=0,1,2,\dots}^{\infty} \gamma_j^{(\ell)} t^\ell, \quad L_j(t) = \sum_{\ell=0,1,2,\dots}^{\infty} L_j^{(\ell)} t^\ell$$

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$$f_{N,t} \leq f_{N,t}^\dagger := 4 \min_{\mathbf{h}(t)} \left\{ N \left\| \frac{\alpha_{\mathbf{h}}(t)}{t} \right\| + N(N-1) \left\| \frac{\beta_{\mathbf{h}}(t)}{\sqrt{t}} \right\|^2 \right\}$$

$$\|\alpha_{\mathbf{h}}(t)\| = (a_0 + a_1 \chi + a_2 \chi^2) t^a + O(t^{a+\frac{1}{2}})$$

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$$\|\alpha_{\mathbf{h}}(t)\| = (a_0 + a_1\chi + a_2\chi^2) t^a + O(t^{a+\frac{1}{2}})$$

$$\|\beta_{\mathbf{h}}(t)\|^2 = \left(b_0 t + b_1 \chi t^b + O(t^{\max\{\frac{3}{2}, b+\frac{1}{2}\}}) \right)^2$$

$$f_{N,t}^\dagger \leq 4N (a_0 + a_1\chi + a_2\chi^2) t^{a-1} + 4N(N-1) \frac{[b_0 t + b_1 \chi t^b]^2}{t}$$

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5) Minimise the inverse of the so-obtained bound over t to obtain the optimal $t_{\text{opt}}(N)$.

$$t_{\text{opt}}(N) \underset{N \gg 1}{=} \left(\frac{1}{2b-a-1} \frac{a_2}{b_1^2} \frac{1}{N} \right)^{\frac{1}{2b-a}}$$

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6) Substitute for the optimal $t_{\text{opt}}(N)$ and compute the final form of the bound in the asymptotic N limit.

$$f_N^\dagger \underset{N \gg 1}{\leq} 4 b_0^2 \frac{2b-a-1}{2b-a} \left(\frac{1}{2b-a-1} \frac{a_2}{b_1^2} \right)^{\frac{1}{2b-a}} N^{2-\frac{1}{2b-a}}$$

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6) Substitute for the optimal $t_{\text{opt}}(N)$ and compute the final form of the bound in the asymptotic N limit.

$$f_N^\dagger \underset{N \gg 1}{\leq} 4 b_0^2 \frac{2b-a-1}{2b-a} \left(\frac{1}{2b-a-1} \frac{a_2}{b_1^2} \right)^{\frac{1}{2b-a}} N^{2-\frac{1}{2b-a}}$$

This way we can analytically construct all the aforementioned asymptotic bounds (and beyond), including the proportionality constants, that scale in between the **SQL** and the **HL**.

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Can one provide an analytic derivation for general time-local dissipator with **analytic rates and Lindblad operators**?

$$\mathcal{D}_t[\rho(t)] = \sum_{j=1}^{l \leq r-1} \gamma_j(t) \left(L_j(t) \rho(t) L_j^\dagger(t) - \frac{1}{2} \left\{ \rho(t), L_j^\dagger(t) L_j(t) \right\} \right) \quad \text{with} \quad \gamma_j(t) = \sum_{\ell=0,1,2,\dots}^{\infty} \gamma_j^{(\ell)} t^\ell, \quad L_j(t) = \sum_{\ell=0,1,2,\dots}^{\infty} L_j^{(\ell)} t^\ell$$

We provide a general ("step by step") recipe how to construct a meaningful bound for a given dynamics.

Summary of the procedure:

4) Minimise the resulting bound by the \mathbf{h} -induced parameter χ : χ $f_{N,t} \leq f_{N,t}^\dagger := 4 \min_{\mathbf{h}(t)} \left\{ N \left\| \frac{\alpha_{\mathbf{h}}(t)}{t} \right\| + N(N-1) \left\| \frac{\beta_{\mathbf{h}}(t)}{\sqrt{t}} \right\|^2 \right\}$

$$\|\alpha_{\mathbf{h}}(t)\| = (a_0 + a_1 \chi + a_2 \chi^2) t^a + O(t^{a+\frac{1}{2}}) \quad \|\beta_{\mathbf{h}}(t)\|^2 = (b_0 t + b_1 \chi t^b + O(t^{\max\{\frac{3}{2}, b+\frac{1}{2}\}}))$$

$$f_{N,t}^\dagger \leq 4N (a_0 + a_1 \chi + a_2 \chi^2) t^{a-1} + 4N(N-1) \frac{[b_0 t + b_1 \chi t^b]^2}{t} \quad \xrightarrow{b>1} \quad = 4N t^{a-1} \left(a_2 \frac{b_0^2}{b_1^2} \frac{1}{t^{2(b-1)}} + O\left(\frac{1}{t^{2b-1}}\right) \right) - \frac{t^{2a-4b+1}}{b_1^4} (2a_2 b_0 + O(t^{b-1}))^2$$

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OPEN QUESTION:

Can it be proven for *any* dynamics that at **lowest orders** one must always deal with a **CP-divisible evolution**?

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How about models that are not CP-divisible at *any* $t > 0$?

e.g.: **qubit eternally non-Markovian dynamics**

$$\mathcal{D}_t^{\text{RU}}[\rho(t)] = \sum_{i=x,y,x} \gamma_i (\sigma_i \rho(t) \sigma_i - \rho(t))$$

“without Zeno”: $(\Gamma \geq 1, \Omega \geq 0)$

“with Zeno”: $(\Gamma^2 \geq \Omega \geq 0)$

$$\boxed{\gamma_x = \gamma_y = \Gamma \frac{\Omega}{2}}$$

$$\boxed{\gamma_x = \Gamma t = \Gamma t}$$

$$\gamma_z = -\Gamma \frac{\Omega}{2} \tanh(\Omega t) = -\frac{1}{2} (\Gamma \Omega^2) t + O(t^2)$$

$$\gamma_z = -\Omega \frac{t^3}{1-t^4} = -\Omega^2 t^3 + O(t^5)$$

$$\boxed{\ell = 0 \implies f_{N,t} \underset{N \rightarrow \infty}{\lesssim} N \quad \text{SQL}}$$

$$\boxed{\ell = 1 \implies f_{N,t} \underset{N \rightarrow \infty}{\lesssim} N^{3/2} \quad \text{>SQL}}$$

OPEN QUESTION:

Can it be proven for *any* dynamics that at **lowest orders** one must always deal with a **CP-divisible evolution**?

Thank you very much for your attention!