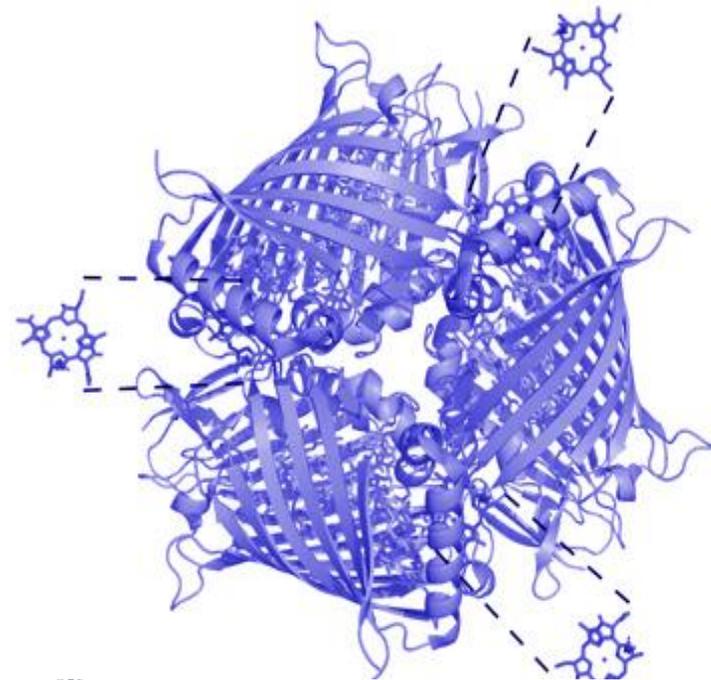


# Coherent Effects in Photosynthetic Complexes

Principles of noise assisted transport and the origin of long-lived coherences

Susana Huelga (Universität Ulm, Germany)



ulm university universität  
**uulm**

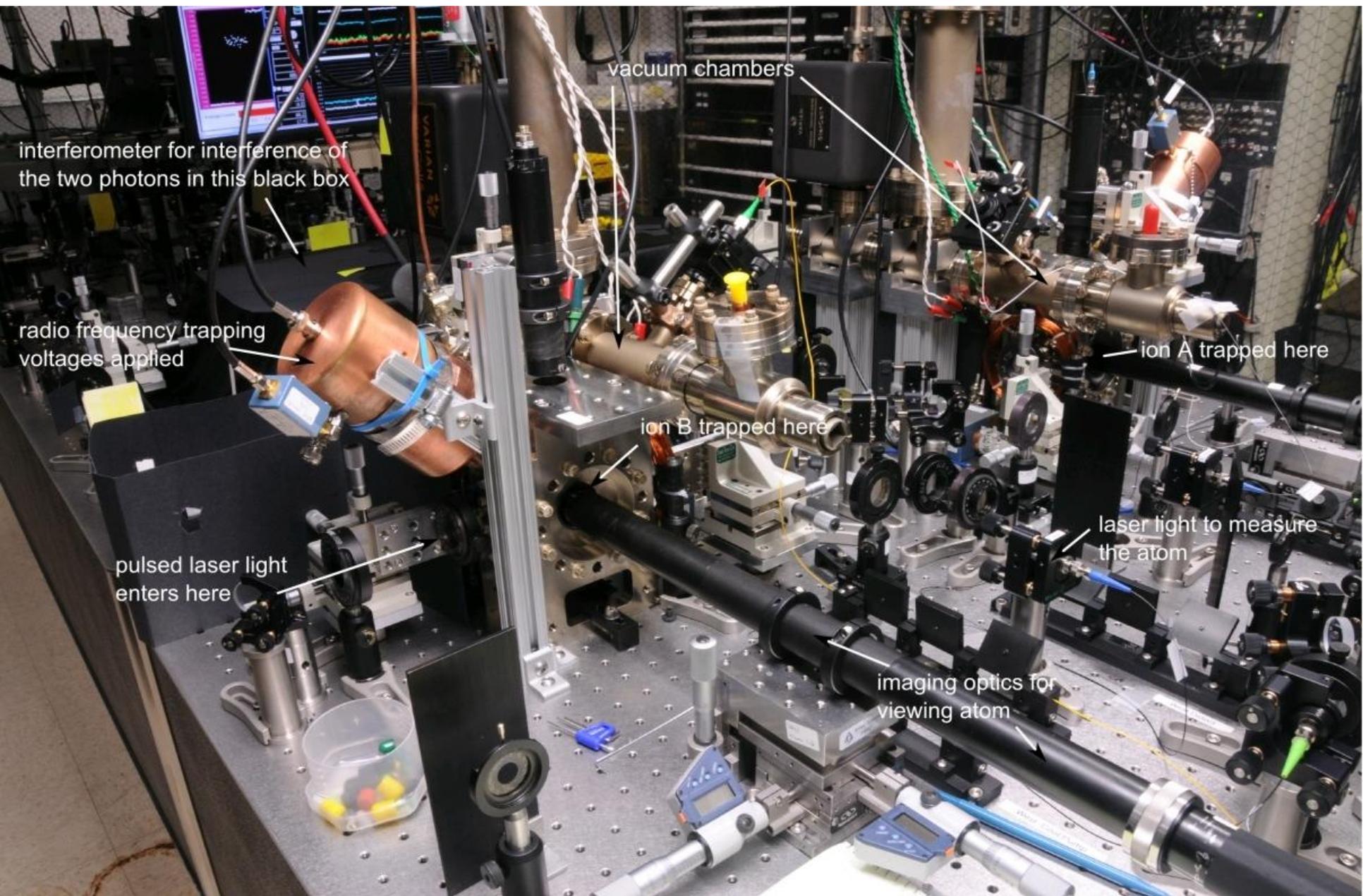
# (Non-trivial) Quantum Effects in Biology??



NJP 13 115002 (2011)

**Focus on quantum effects and noise in biomolecules**

GR Fleming, SF Huelga and MB Plenio Eds



vacuum chambers

interferometer for interference of the two photons in this black box

radio frequency trapping voltages applied

ion A trapped here

ion B trapped here

laser light to measure the atom

pulsed laser light enters here

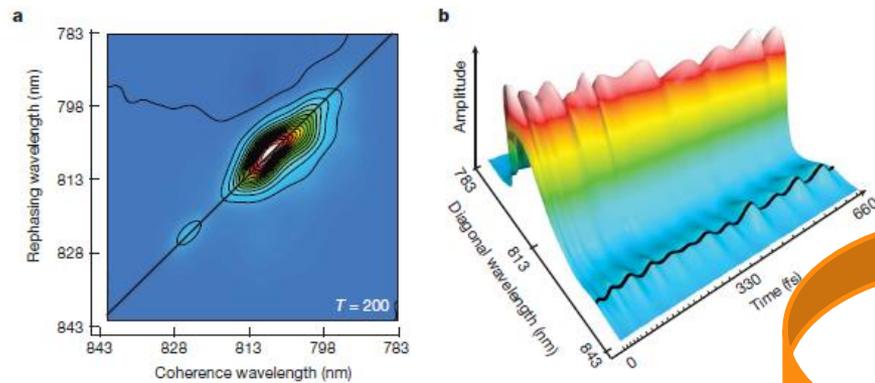
imaging optics for viewing atom



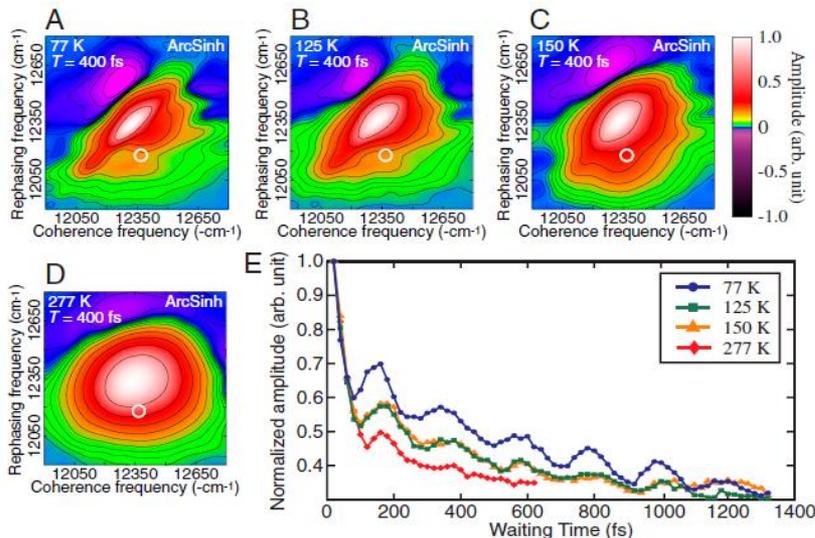
## LETTERS

# Evidence for wavelike energy transfer through quantum coherence in photosynthetic systems

Gregory S. Engel<sup>1,2</sup>, Tessa R. Calhoun<sup>1,2</sup>, Elizabeth L. Read<sup>1,2</sup>, Tae-Kyu Ahn<sup>1,2</sup>, Tomáš Mančal<sup>1,2,†</sup>, Yuan-Chung Cheng<sup>1,2</sup>, Robert E. Blankenship<sup>3,4</sup> & Graham R. Fleming<sup>1,2</sup>



Experimental evidence at cryogenic and physiological temperatures indicates the presence of quantum coherence



PNAS

[www.pnas.org/cgi/doi/10.1073/pnas.1005484107](http://www.pnas.org/cgi/doi/10.1073/pnas.1005484107)

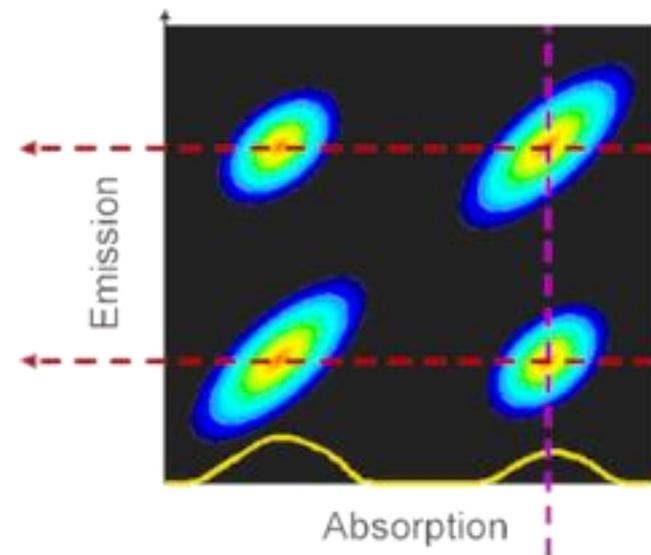
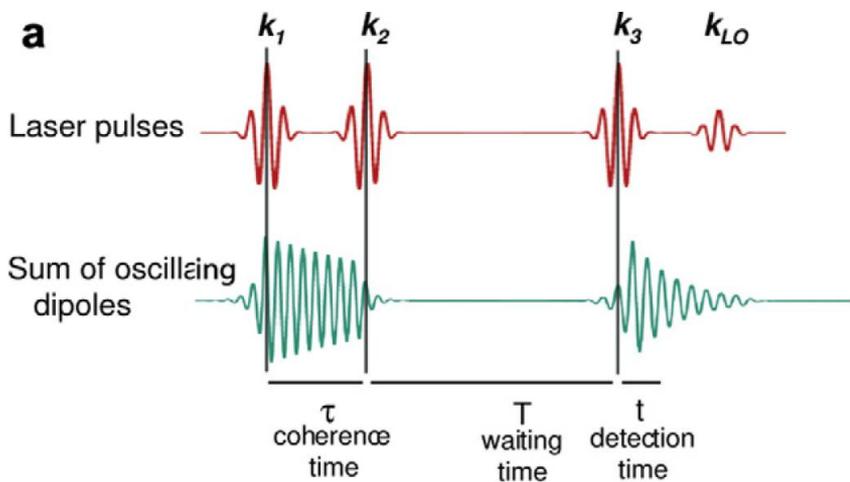
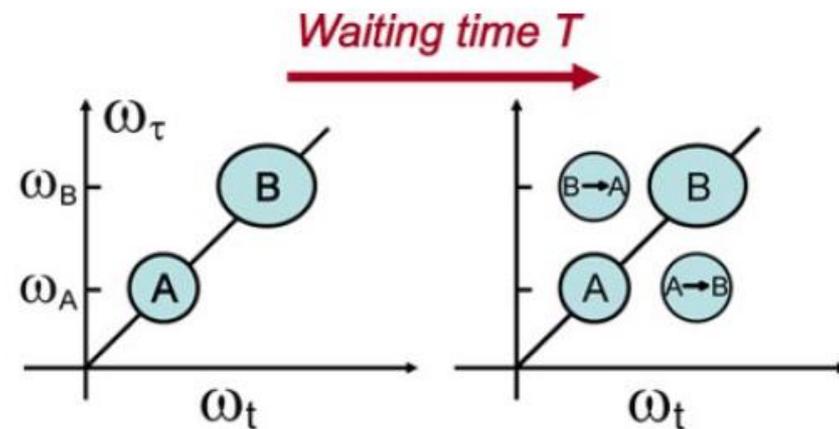
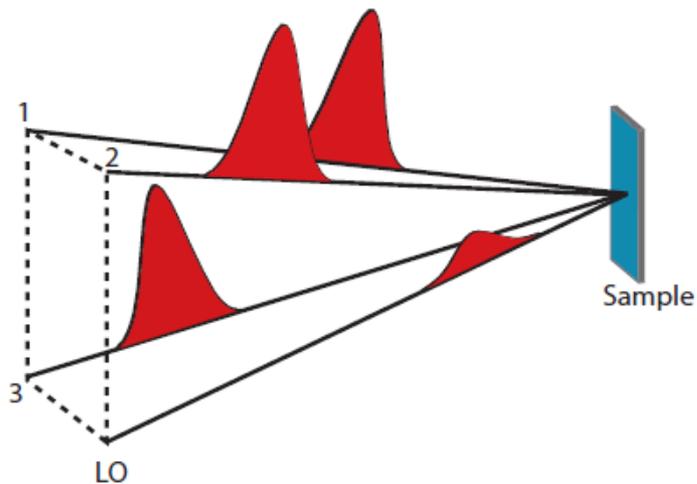
## Long-lived quantum coherence in photosynthetic complexes at physiological temperature

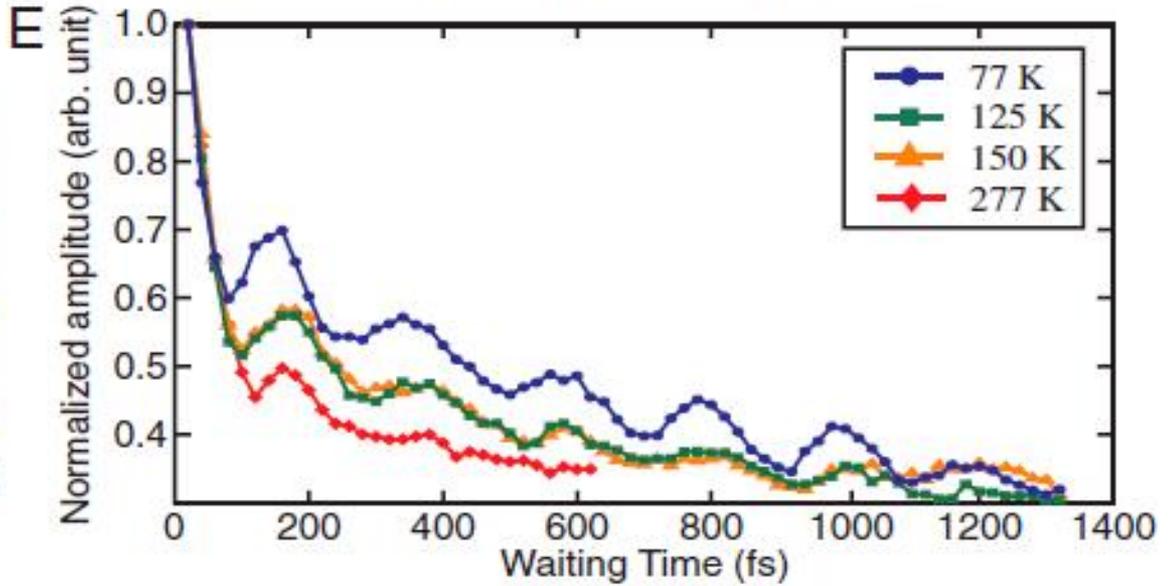
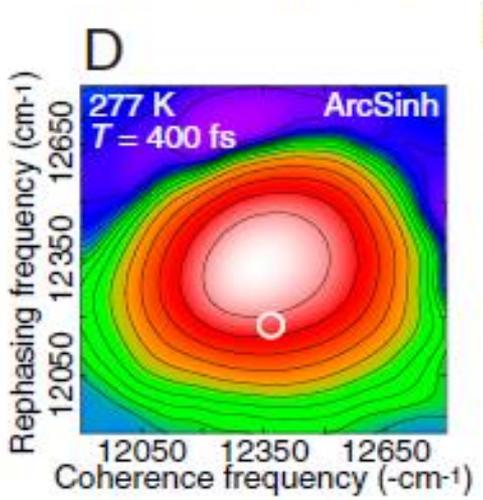
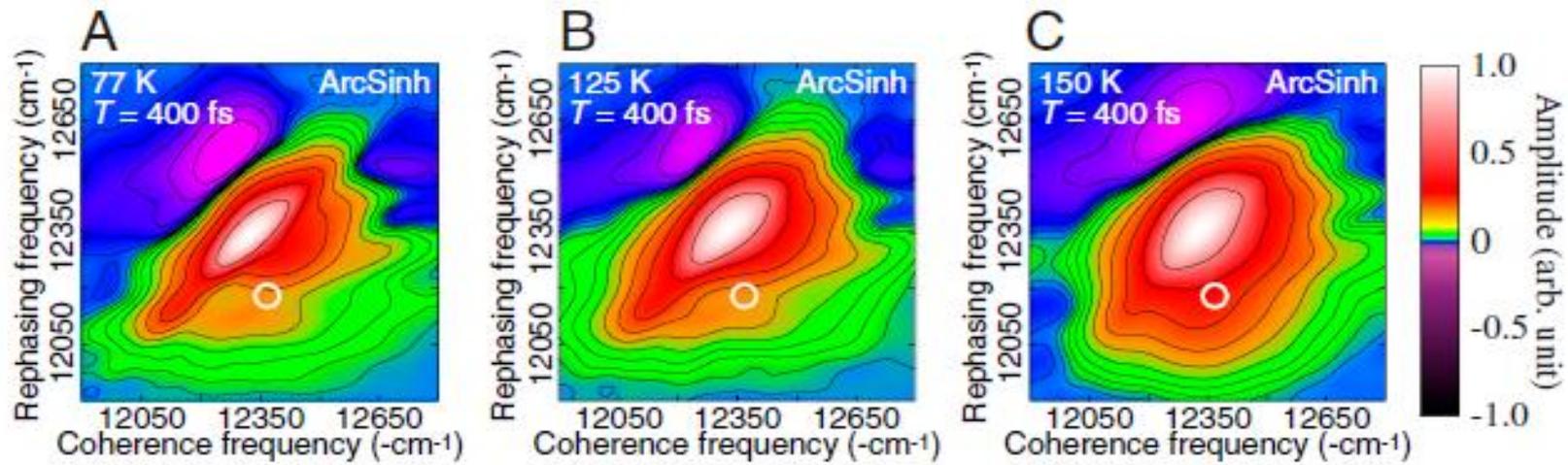
Gitt Panitchayangkoon<sup>a</sup>, Dugan Hayes<sup>a</sup>, Kelly A. Fransted<sup>a</sup>, Justin R. Caram<sup>a</sup>, Elad Harel<sup>a</sup>, Jianzhong Wen<sup>b</sup>, Robert E. Blankenship<sup>b</sup>, and Gregory S. Engel<sup>a,1</sup>

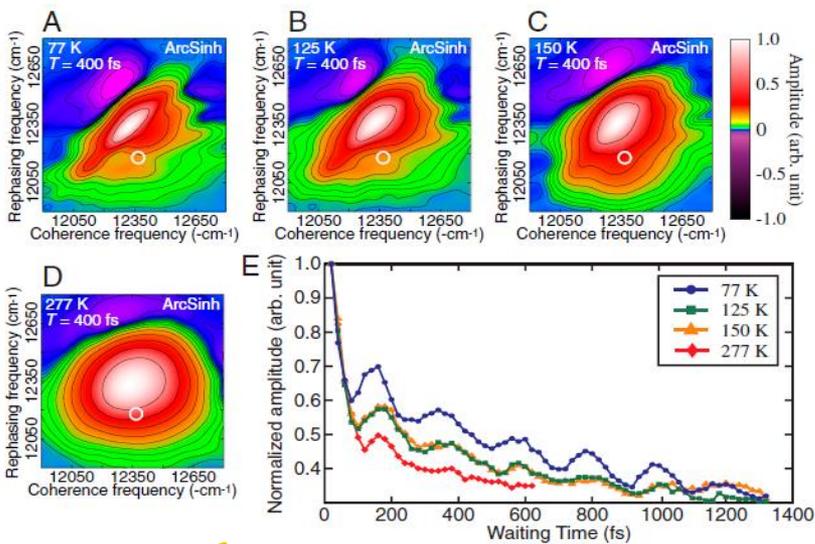
<sup>a</sup>Department of Chemistry and The James Franck Institute, University of Chicago, Chicago, IL 60637; and <sup>b</sup>Departments of Biology and Chemistry, Washington University, St. Louis, MO 63130

Other results: Scholes (Toronto), van Hulst (ICFO)...

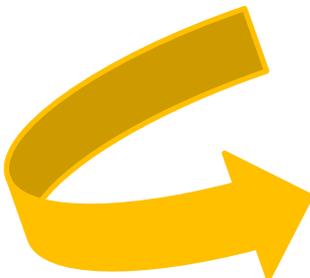
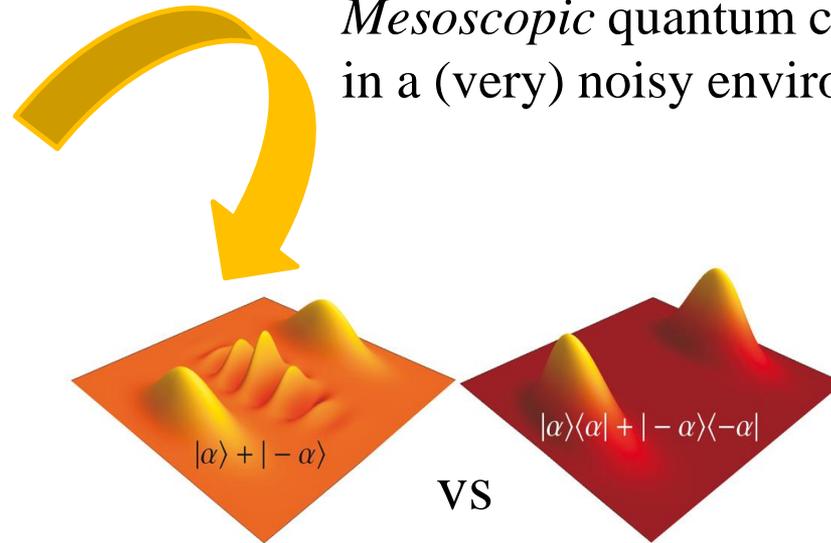
# 2D Spectroscopy







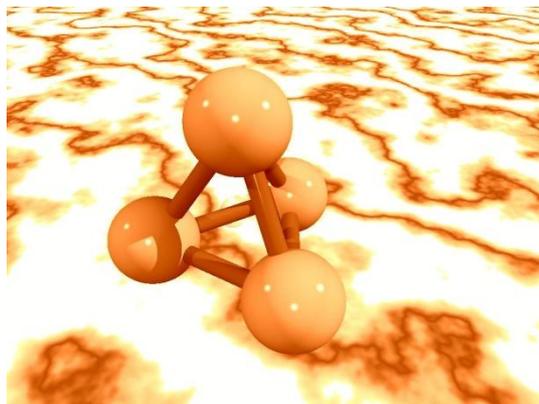
*Mesoscopic quantum coherence in a (very) noisy environment*



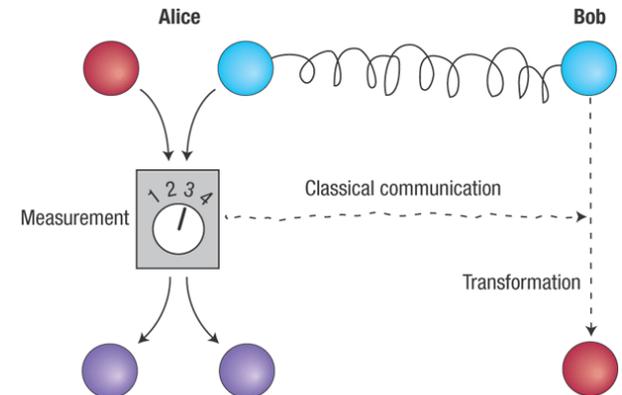
And beyond

A fundamental question arises: Is there any link between coherence and biological function?

Do/could quantum features **facilitate** processes of biological relevance?

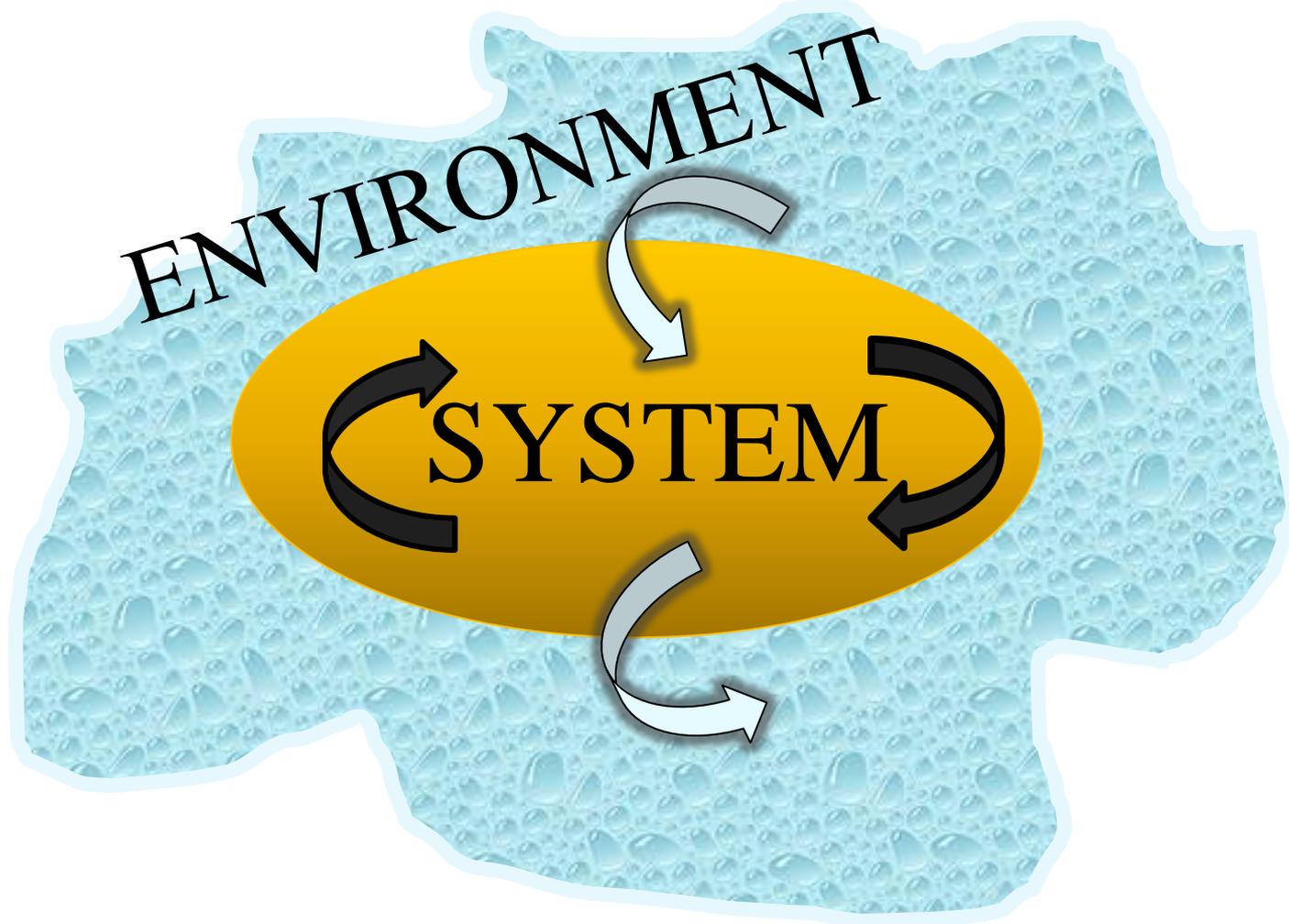


VS



# More prosaic:

New perspectives for non-equilibrium open quantum systems



Careful **interplay between coherent and *noisy* interactions** is key for keeping quantum traits and for them to survive (even) in the steady state

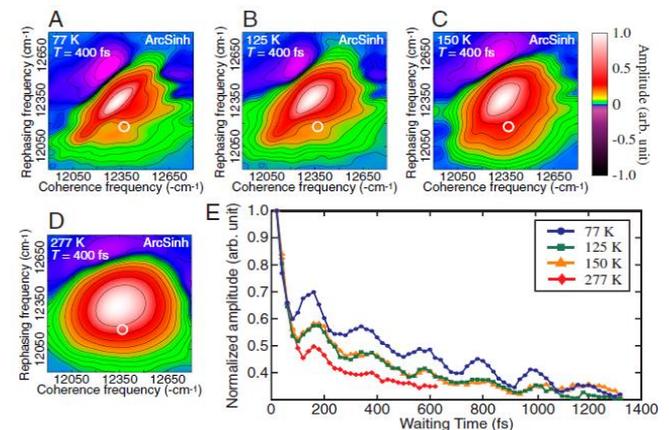
# Outline

## Noise assisted transport: Fundamental mechanisms

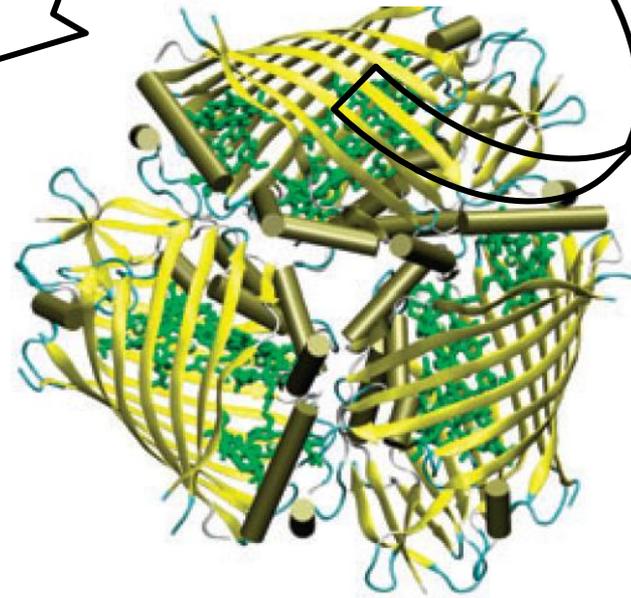
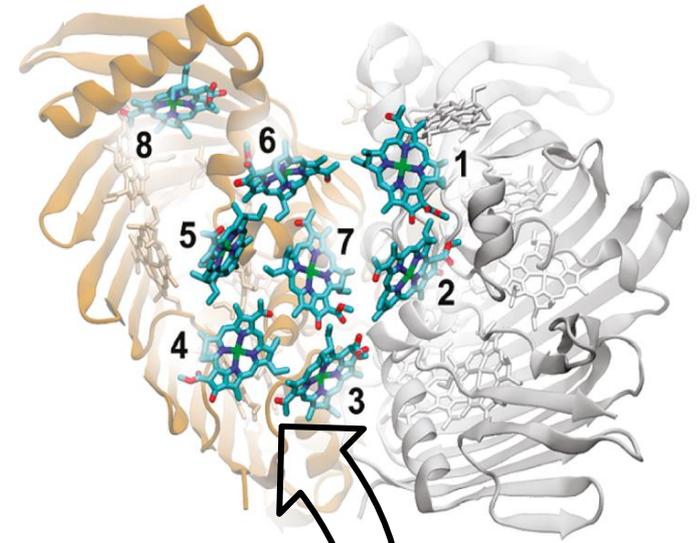
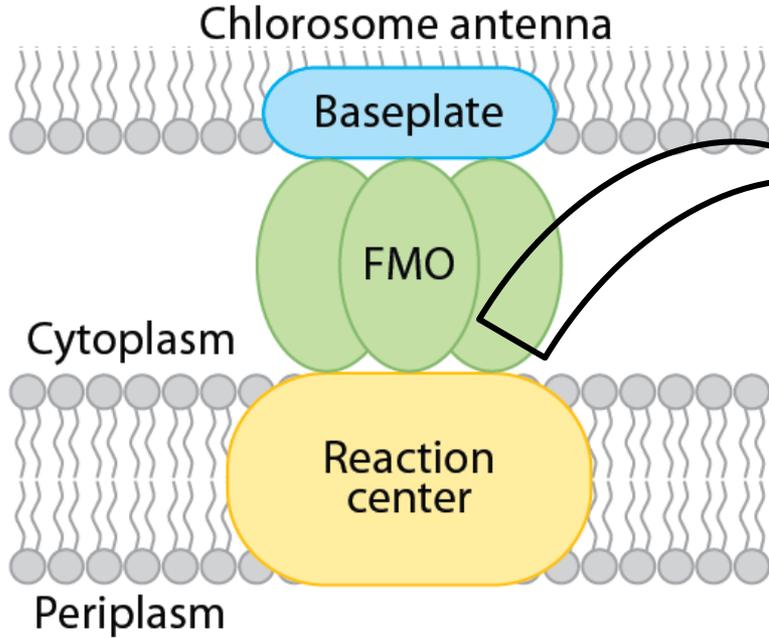
The concept of *phonon antenna*

## Long lived coherence

How can quantum coherence persist in a *wet* and *hot* environment?

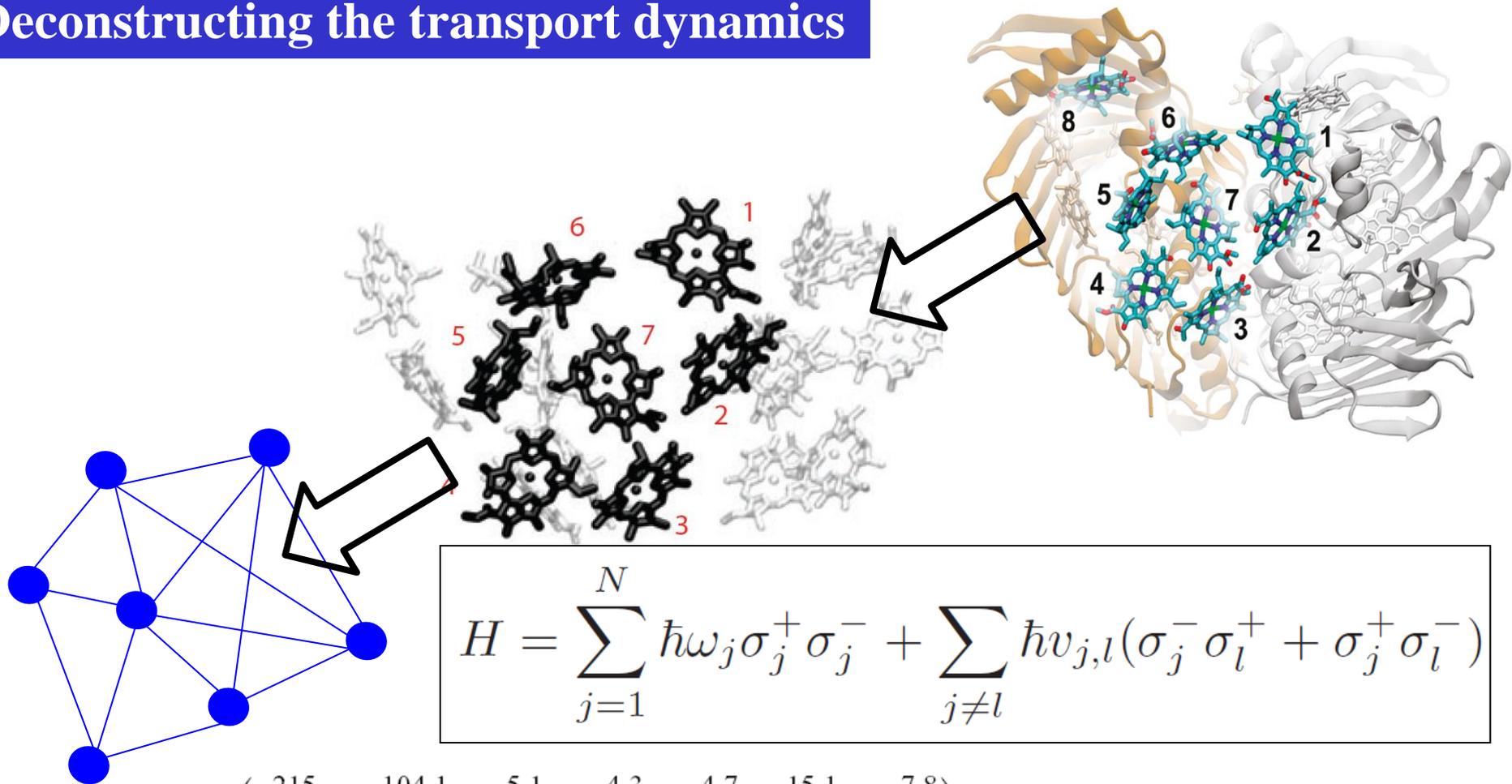


# Fenna-Matthew-Olson (FMO) complex



- R.E. Fenna and B.W. Matthews, *Nature* **258** 573 (1975)
- Y.-C. Cheng & G.R. Fleming, *Annual Reviews of Phys. Chem.* 2009
- Schmidt am Busch et al, *The Eighth Bacteriochlorophyll Completes the Excitation Energy Funnel in the FMO Protein.* *J. Phys. Chem. Lett.* **2**, 93 (2011)

# Deconstructing the transport dynamics

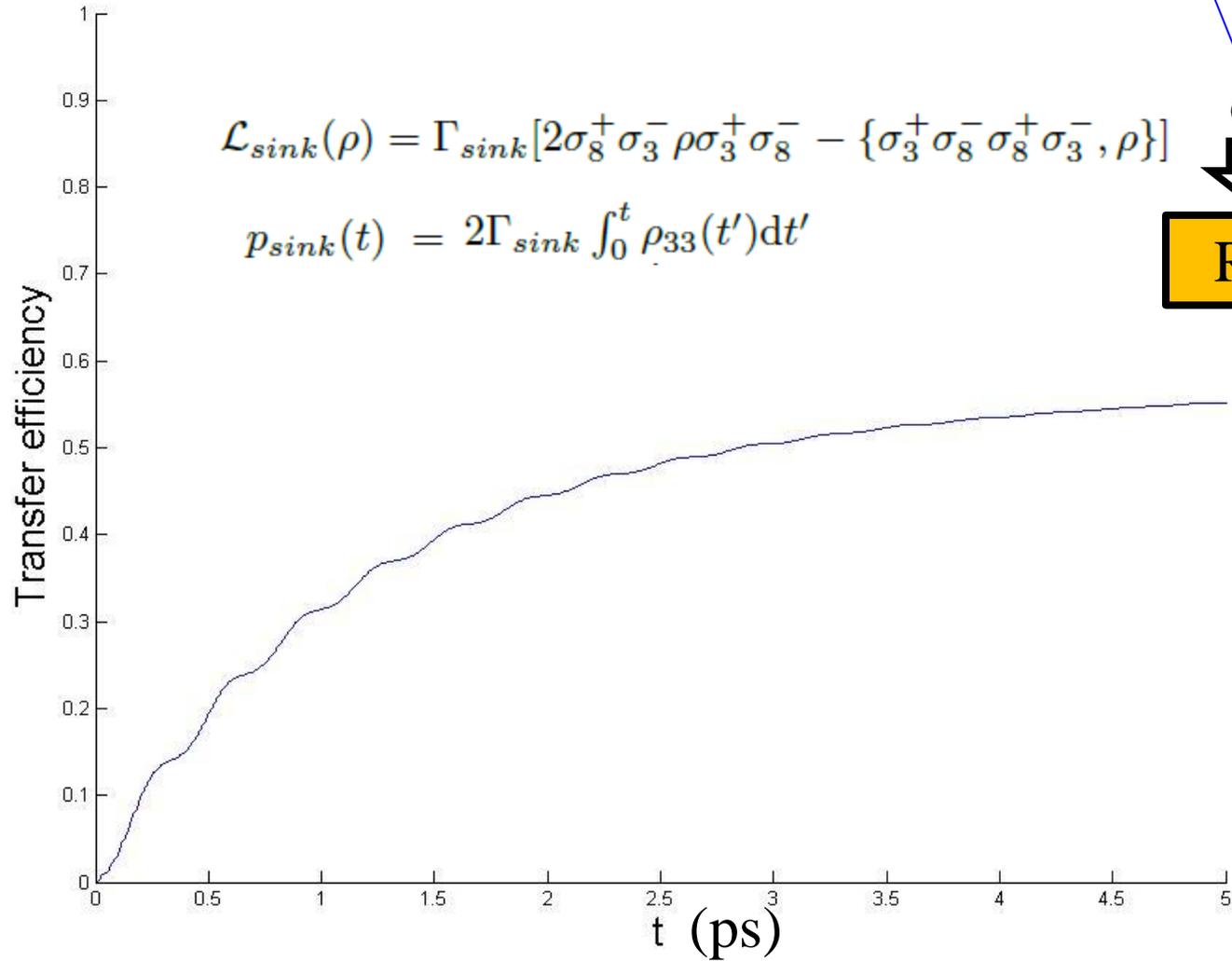


$$H = \sum_{j=1}^N \hbar\omega_j \sigma_j^+ \sigma_j^- + \sum_{j \neq l} \hbar v_{j,l} (\sigma_j^- \sigma_l^+ + \sigma_j^+ \sigma_l^-)$$

$$H = \begin{pmatrix} 215 & -104.1 & 5.1 & -4.3 & 4.7 & -15.1 & -7.8 \\ -104.1 & 220.0 & 32.6 & 7.1 & 5.4 & 8.3 & 0.8 \\ 5.1 & 32.6 & 0.0 & -46.8 & 1.0 & -8.1 & 5.1 \\ -4.3 & 7.1 & -46.8 & 125.0 & -70.7 & -14.7 & -61.5 \\ 4.7 & 5.4 & 1.0 & -70.7 & 450.0 & 89.7 & -2.5 \\ -15.1 & 8.3 & -8.1 & -14.7 & 89.7 & 330.0 & 32.7 \\ -7.8 & 0.8 & 5.1 & -61.5 & -2.5 & 32.7 & 280.0 \end{pmatrix}$$

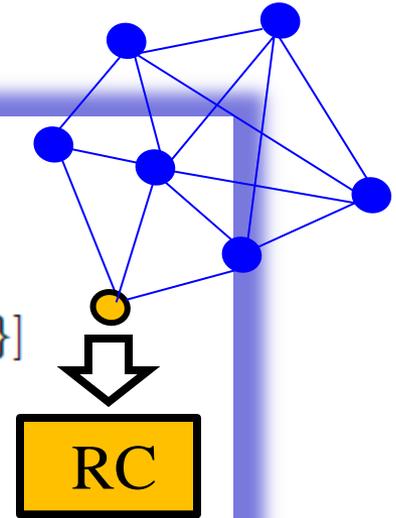
Adapted from Adolphs&Renger 2006  
Plenio&Huelga, New J. Phys. **10**, 113019 (2008)

(Units  $1.2414 \cdot 10^{-4}$  eV)

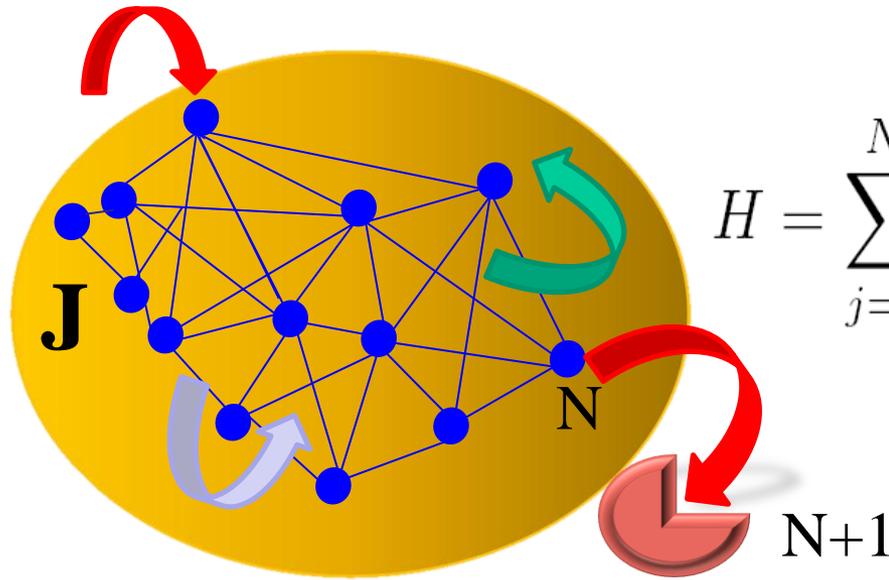


$$\mathcal{L}_{sink}(\rho) = \Gamma_{sink} [2\sigma_8^+ \sigma_3^- \rho \sigma_3^+ \sigma_8^- - \{\sigma_3^+ \sigma_8^- \sigma_8^+ \sigma_3^-, \rho\}]$$

$$p_{sink}(t) = 2\Gamma_{sink} \int_0^t \rho_{33}(t') dt'$$



## Further simplification: Fully Connected Networks (FCN)



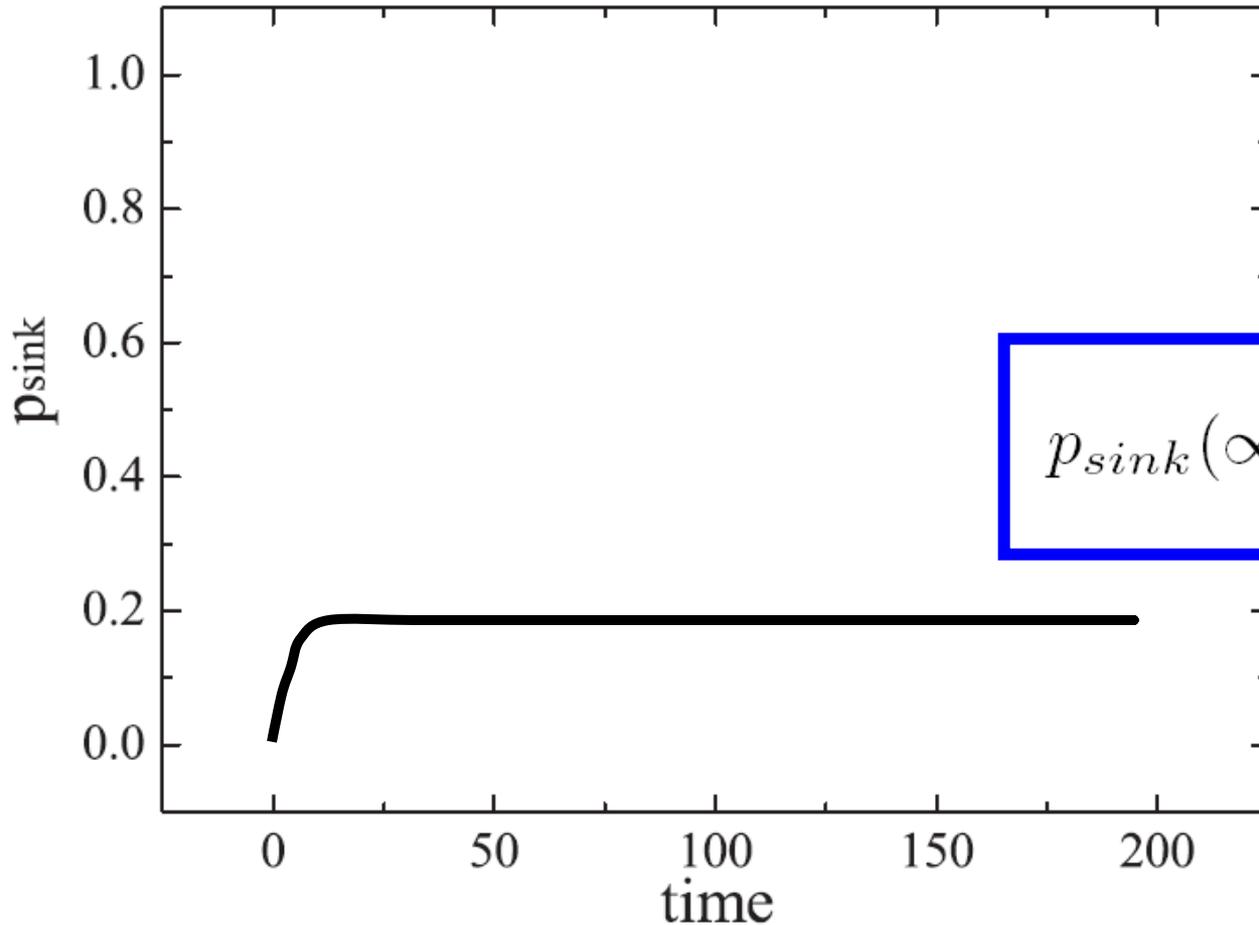
$$H = \sum_{j=1}^N \hbar \omega_j \sigma_j^+ \sigma_j^- + \sum_{j \neq l} \hbar \mathbf{J} (\sigma_j^- \sigma_l^+ + \sigma_j^+ \sigma_l^-)$$

Single irreversible process: population transfer from site  $N$  to the sink  $N+1$

$$\mathcal{L}_{sink}(\rho) = \Gamma_{N+1} [2\sigma_{N+1}^+ \sigma_k^- \rho \sigma_k^+ \sigma_{N+1}^- - \{\sigma_k^+ \sigma_{N+1}^- \sigma_{N+1}^+ \sigma_k^-, \rho\}]$$

In a given time  $T$ , how much of the initial population in site 1 can be transferred to site  $N+1$  (**trapping site**) and how is the transport affected by noise?

$$p_{sink}(t) = 2\Gamma_{N+1} \int_0^t \rho_{kk}(t') dt'$$



$$p_{sink}(\infty) = \frac{1}{N-1}$$

$$\Gamma = 0, J = 1, \Gamma_{N+1} = 1$$

$$\gamma_i = \omega_i = 0$$

# Deconstructing the transport dynamics

Opinion: Well, this is obvious

$$H = \begin{pmatrix} 215 & -104.1 & 5.1 & -4.3 & 4.7 & -15.1 & -7.8 \\ -104.1 & 220.0 & 32.6 & 7.1 & 5.4 & 8.3 & 0.8 \\ 5.1 & 32.6 & 0.0 & -46.8 & 1.0 & -8.1 & 5.1 \\ -4.3 & 7.1 & -46.8 & 125.0 & -70.7 & -14.7 & -61.5 \\ 4.7 & 5.4 & 1.0 & -70.7 & 450.0 & 89.7 & -2.5 \\ -15.1 & 8.3 & -8.1 & -14.7 & 89.7 & 330.0 & 32.7 \\ -7.8 & 0.8 & 5.1 & -61.5 & -2.5 & 32.7 & 280.0 \end{pmatrix}$$

**Site basis**

# Deconstructing the dynamics

Opinion: This is obvious

$$H = \begin{pmatrix} 215 & -104.1 & 5.1 & -4.3 & 4.7 & -15.1 & -7.8 \\ -104.1 & 220.0 & 32.6 & 7.1 & 5.4 & 8.3 & 0.8 \\ 5.1 & 32.6 & 0.0 & -46.8 & 1.0 & -8.1 & 5.1 \\ -4.3 & 7.1 & -46.8 & 125.0 & -70.7 & -14.7 & -61.5 \\ 4.7 & 5.4 & 1.0 & -70.7 & 450.0 & 89.7 & -2.5 \\ -15.1 & 8.3 & -8.1 & -14.7 & 89.7 & 330.0 & 32.7 \\ -7.8 & 0.8 & 5.1 & -61.5 & -2.5 & 32.7 & 280.0 \end{pmatrix}$$

**Site basis**

$$H = \begin{pmatrix} 513 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 332 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 307 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 268 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 121 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 102 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -23 \end{pmatrix}$$

**Exciton basis**



No coherent dynamics left



Noise supports transport

# Deconstructing the transport dynamics

Opinion: This is obvious

$$H = \begin{pmatrix} 215 & -104.1 & 5.1 & -4.3 & 4.7 & -15.1 & -7.8 \\ -104.1 & 220.0 & 32.6 & 7.1 & 5.4 & 8.3 & 0.8 \\ 5.1 & 32.6 & 0.0 & -46.8 & 1.0 & -8.1 & 5.1 \\ -4.3 & 7.1 & -46.8 & 125.0 & -70.7 & -14.7 & -61.5 \\ 4.7 & 5.4 & 1.0 & -70.7 & 450.0 & 89.7 & -2.5 \\ -15.1 & 8.3 & -8.1 & -14.7 & 89.7 & 330.0 & 32.7 \\ -7.8 & 0.8 & 5.1 & -61.5 & -2.5 & 32.7 & 280.0 \end{pmatrix}$$

Site basis

**But: Answer depends on structure of network**

Sometimes noise helps  
sometimes it hinders transport

$$H = \begin{pmatrix} 513 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 332 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 307 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 268 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 121 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 102 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -23 \end{pmatrix}$$

**Exciton basis**



No coherent dynamics left



Noise supports transport

**Moreover: Is exciton relaxation entirely incoherent?**

# How can *noise* assist transport?

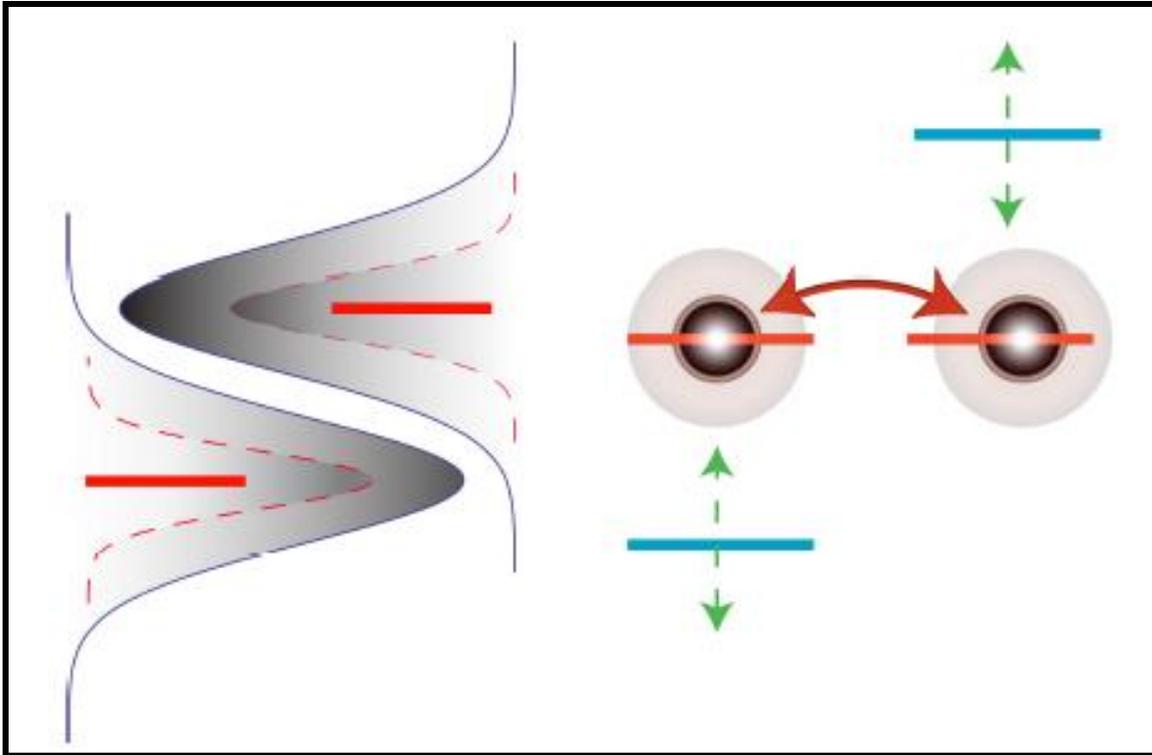
Identifying the building blocks for

**NAT (noise assisted transport)**

**(or ENAQT (environment assisted  
quantum transport) )**

# Basic Mechanisms underpinning noise assisted transport

## (1) Bridging energy gaps and blocking unfavourable (coherent) transport paths



Typical **line broadening** mechanisms

Excess noise may not always be detrimental

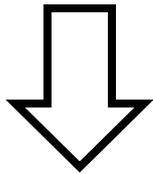
# Basic Mechanisms underpinning noise assisted transport

## (2) Interference effects

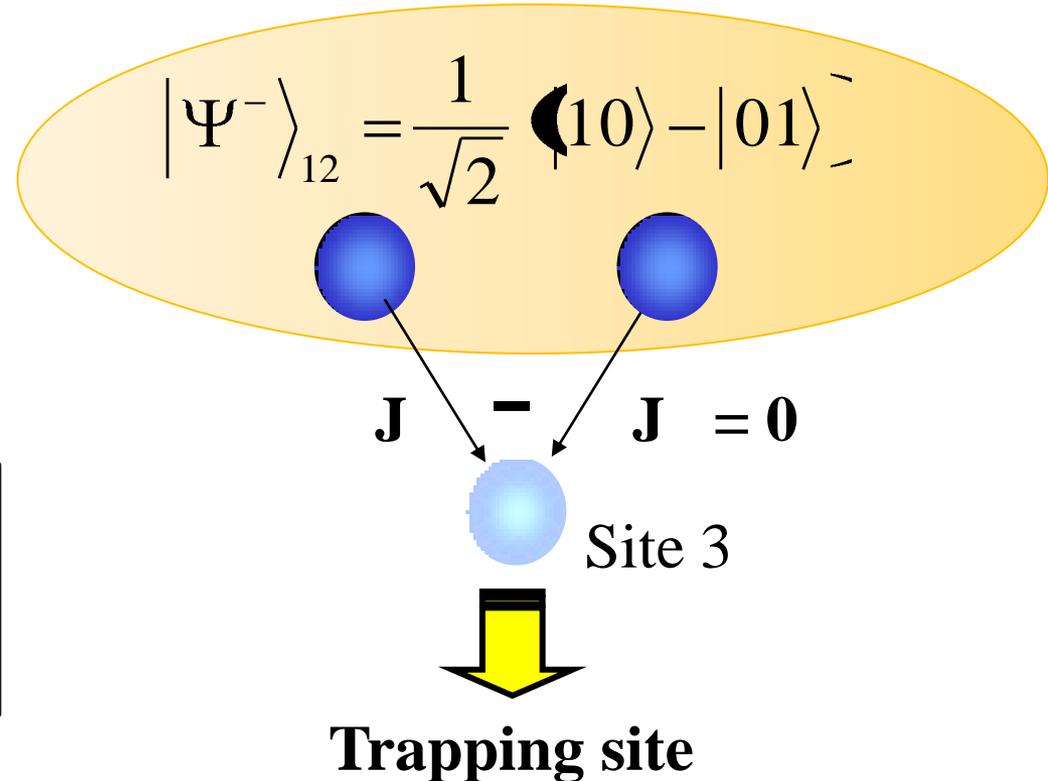
Simplest case: N=3 network

$$H = \sum_{k=1}^3 E_i |i\rangle \langle i| + \sum_{k=1}^2 J_{k3} (|k\rangle \langle 3| + h.c.)$$

$$J_{13} = J_{23}$$



Destructive interference of tunneling amplitudes.  
Transport to site 3 is inhibited

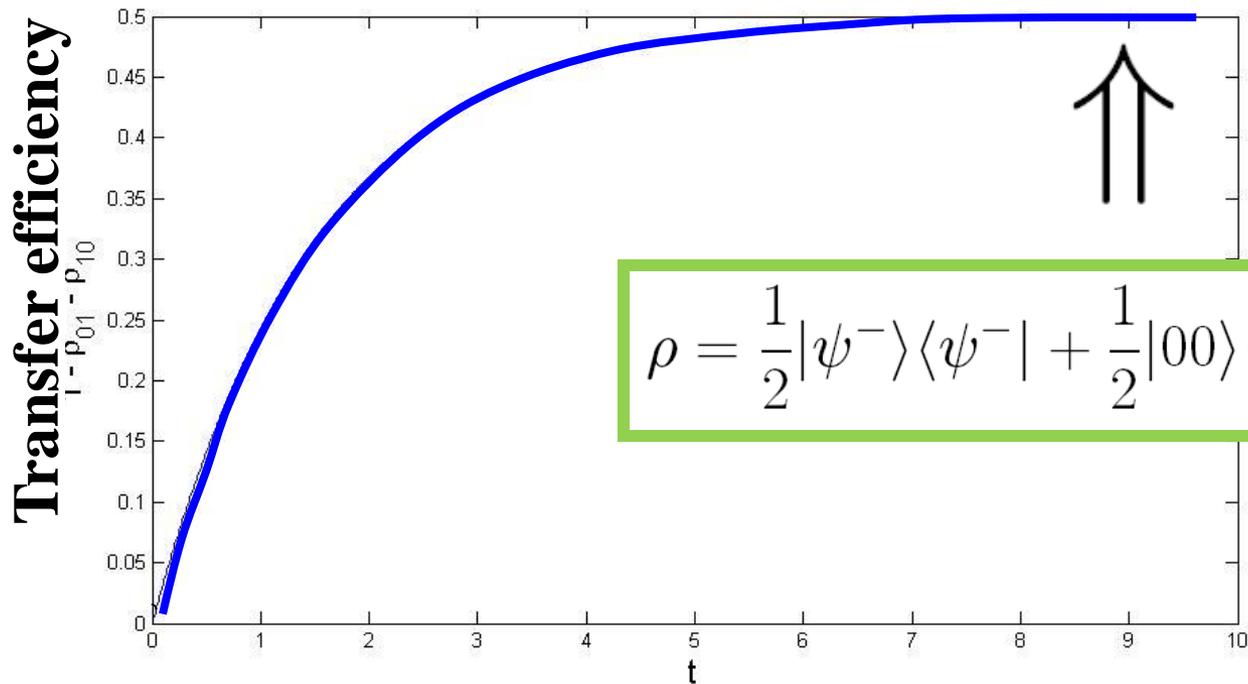


# Basic Mechanisms underpinning noise assisted transport

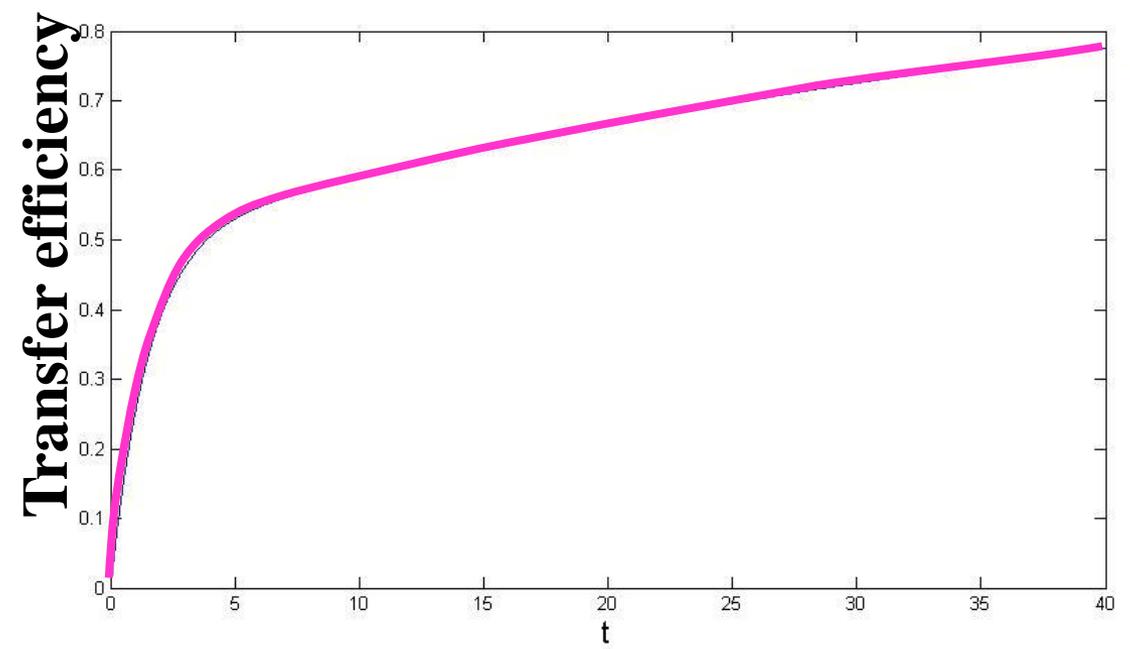
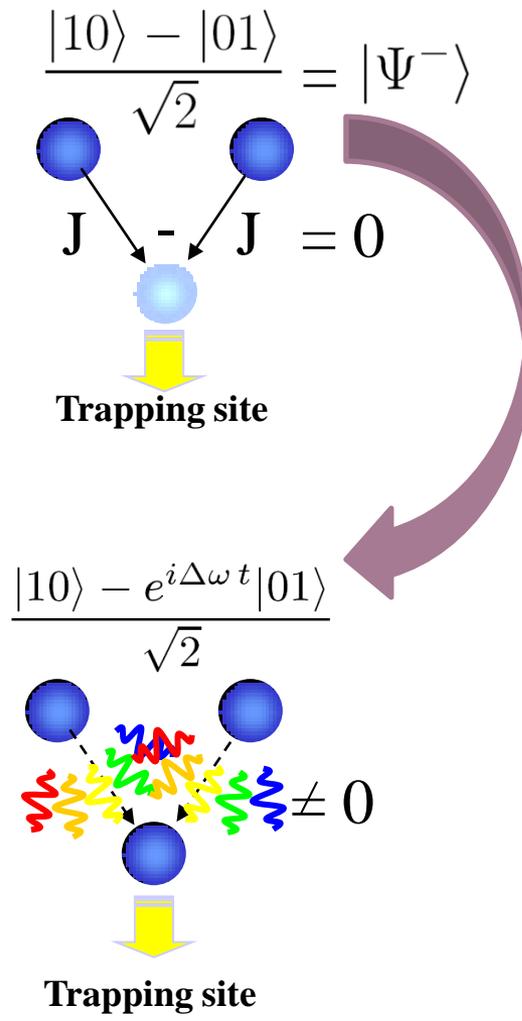
## (2) Interference effects

Selective excitation:

$$|01\rangle = \frac{1}{\sqrt{2}} \left[ \frac{|01\rangle - |10\rangle}{\sqrt{2}} + \frac{|01\rangle + |10\rangle}{\sqrt{2}} \right]$$

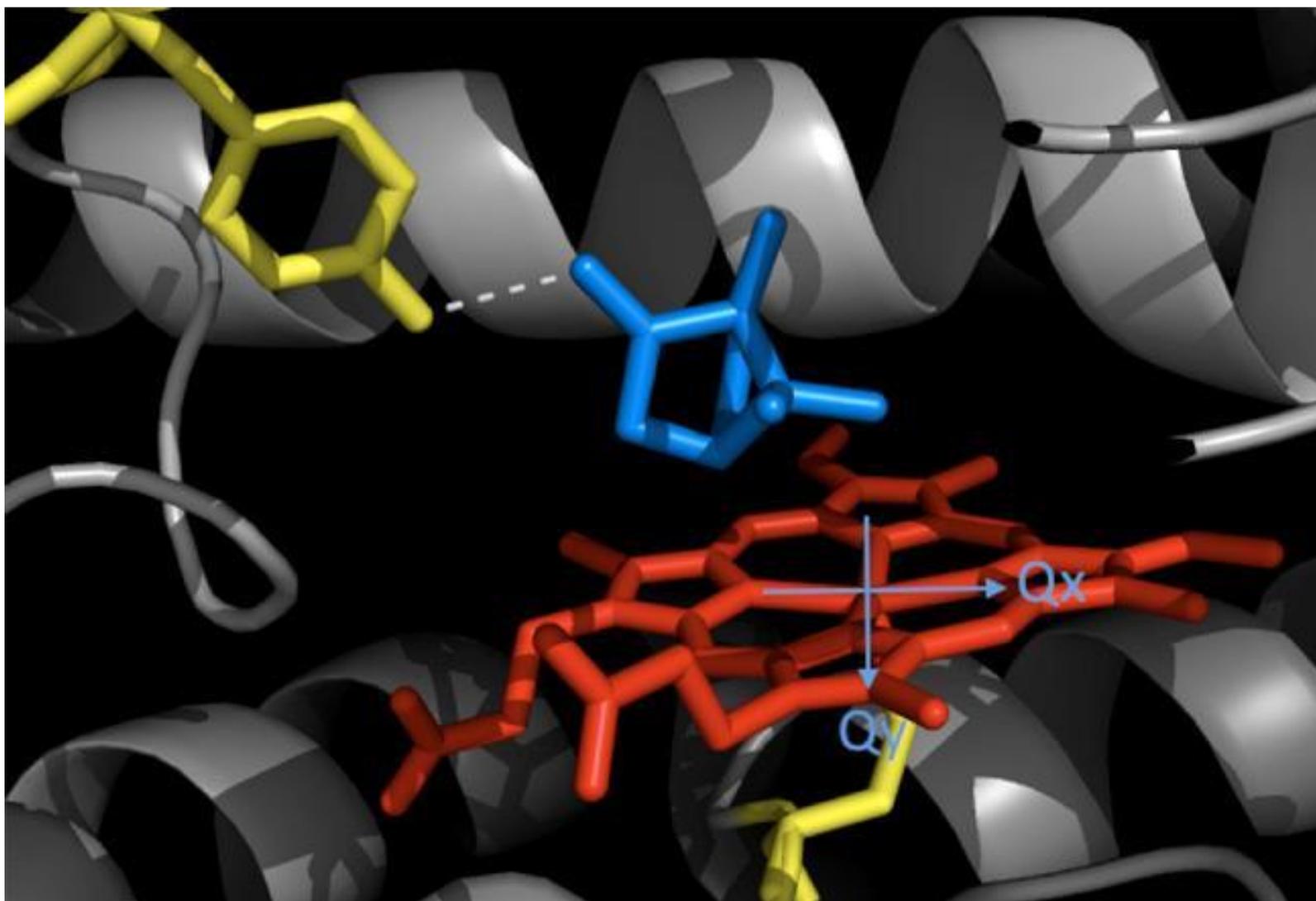


# What effects can release a non-propagating state ?



Reduction of destructive interference effects  $\longrightarrow$  Transport

# Is this relevant in the context of EET?

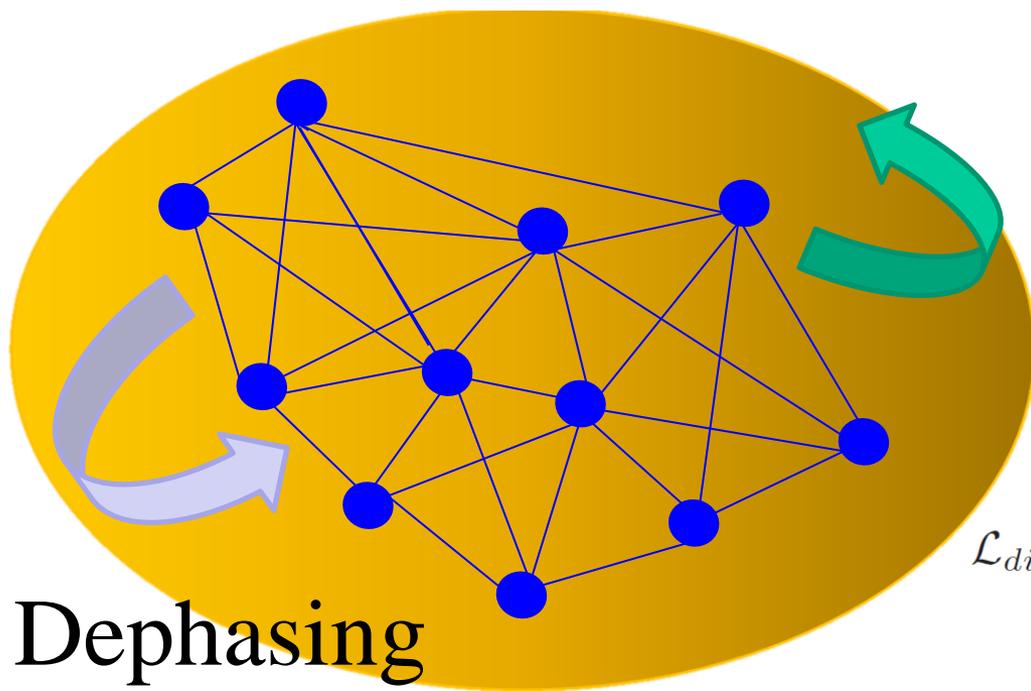


Dominant noise source: Pure dephasing due to protein environment

# Simplest Theoretical Model

$$\frac{d\rho}{dt} = -i[H, \rho] + \mathcal{L}_{deph}(\rho) + \mathcal{L}_{diss}(\rho)$$

$$H = \sum_{j=1}^N \hbar\omega_j \sigma_j^+ \sigma_j^- + \sum_{j \neq l} \hbar v_{j,l} (\sigma_j^- \sigma_l^+ + \sigma_j^+ \sigma_l^-)$$



Relaxation  
Energy exchange

$$\mathcal{L}_{diss}(\rho) = \sum_{j=1}^N \Gamma_j [-\{\sigma_j^+ \sigma_j^-, \rho\} + 2\sigma_j^- \rho \sigma_j^+]$$

(ps range)

Dephasing

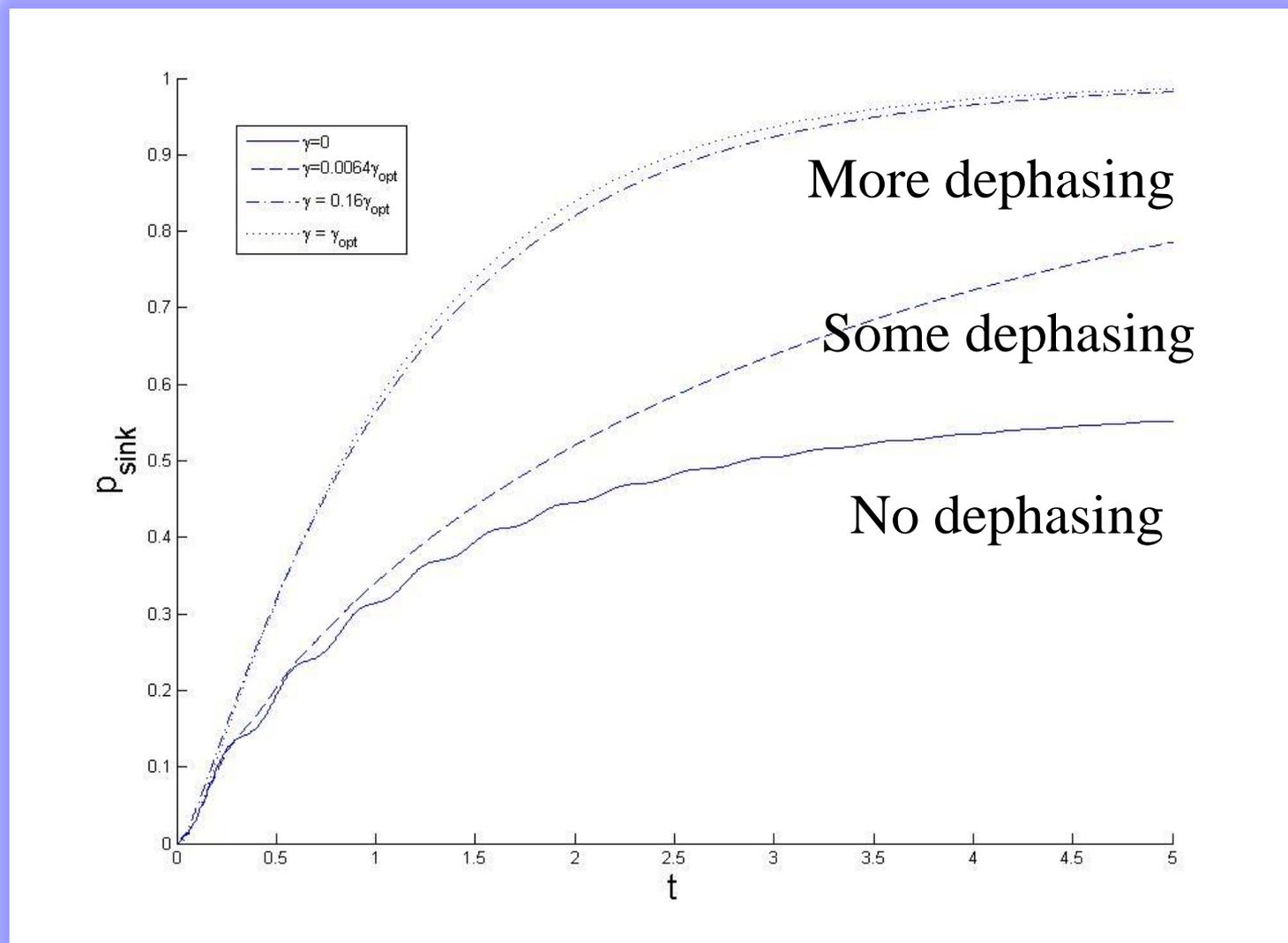
Phase randomization

$$\mathcal{L}_{deph}(\rho) = \sum_{j=1}^N \gamma_j [-\{\sigma_j^+ \sigma_j^-, \rho\} + 2\sigma_j^+ \sigma_j^- \rho \sigma_j^+ \sigma_j^-]$$

Faster time scale (fs)

Dynamical Map is CPTP

# Does it work for transport across FMO?



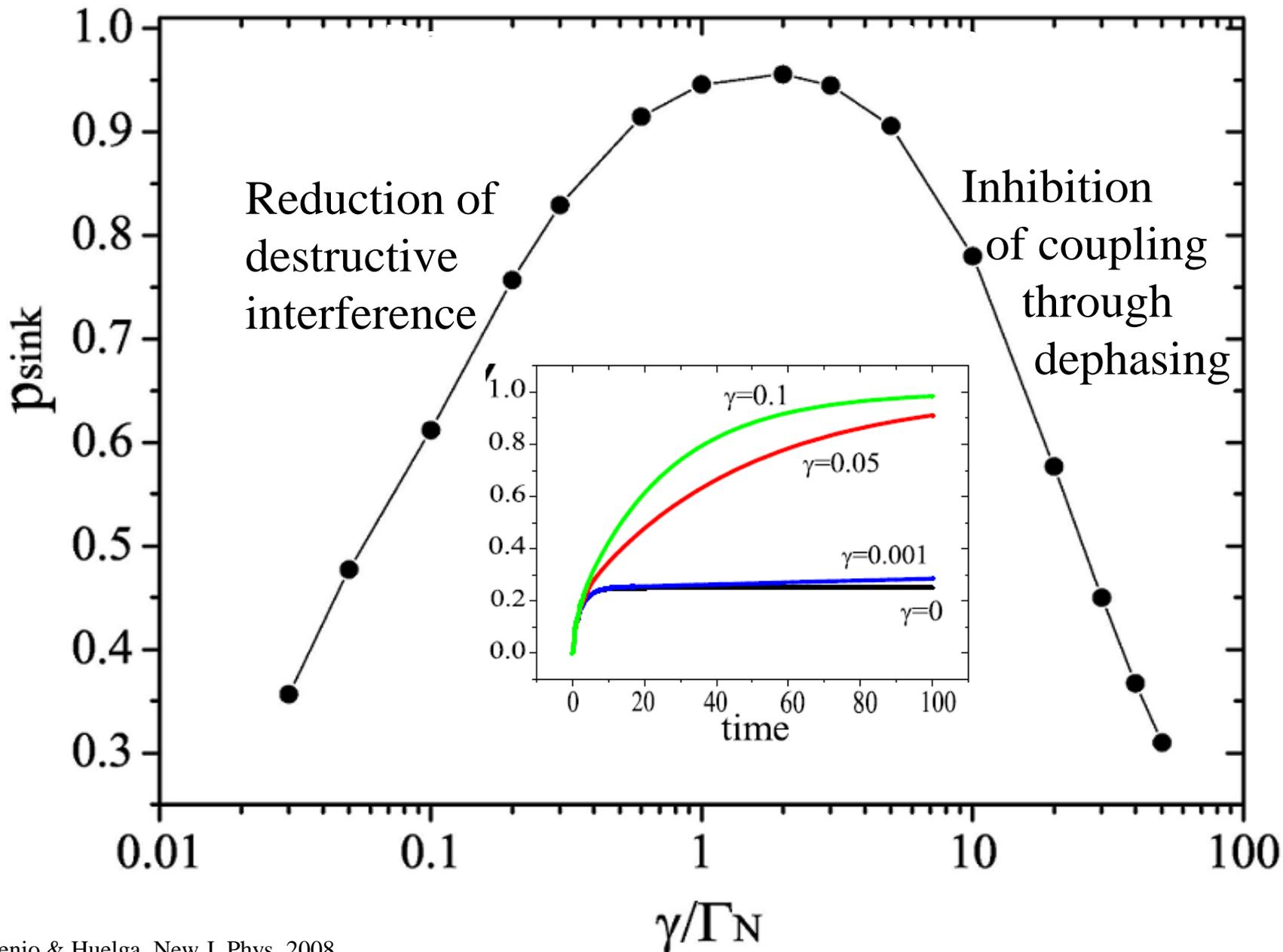
Plenio & Huelga, New J. Phys. 2008

Caruso, Chin, Datta, Huelga, Plenio, J. Chem. Phys. 2009

Chin, Caruso, Datta, Huelga, Plenio, New J. Phys. 2010

Mohseni, Rebentrost, Lloyd, Aspuru-Guzik, J. Phys. Chem. 2008

Rebentrost, Mohseni, Kassal, Lloyd, Aspuru-Guzik, New J. Phys. 2009



Plenio & Huelga, New J. Phys. 2008

Caruso, Chin, Datta, Huelga, Plenio, J. Chem. Phys. 2009

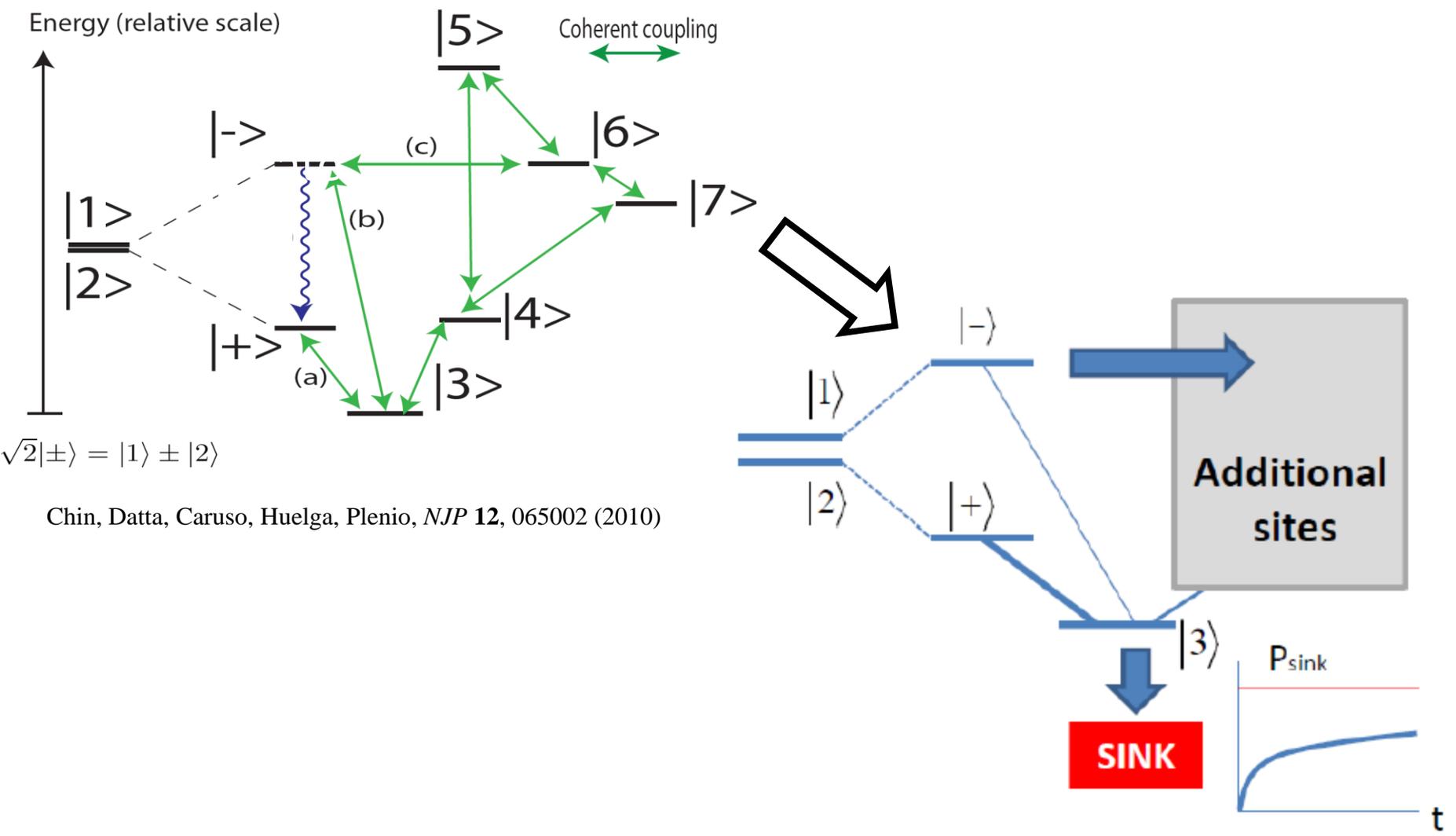
Chin, Caruso, Datta, Huelga, Plenio, New J. Phys. 2010

Mohseni, Rebentrost, Lloyd, Aspuru-Guzik, J. Phys. Chem. 2008

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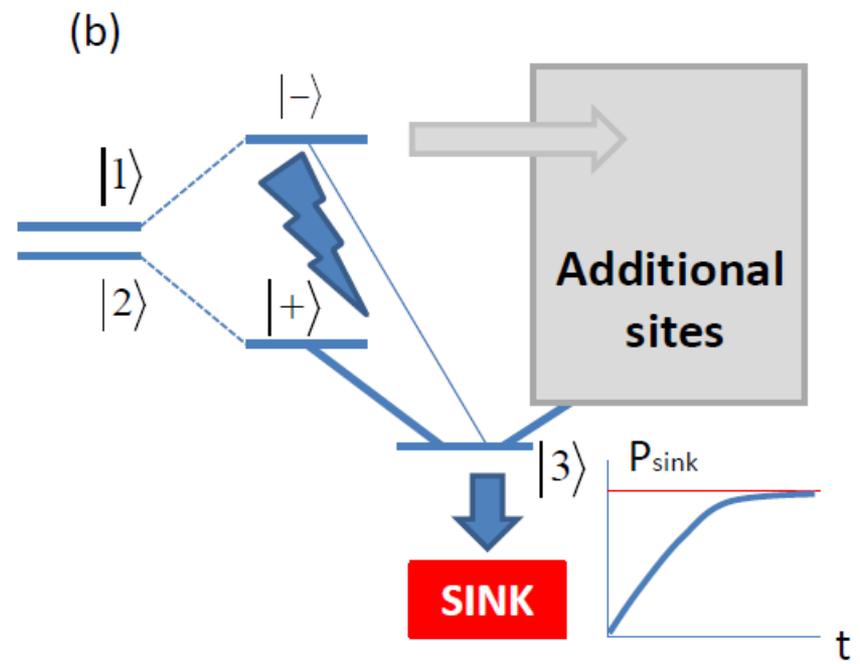
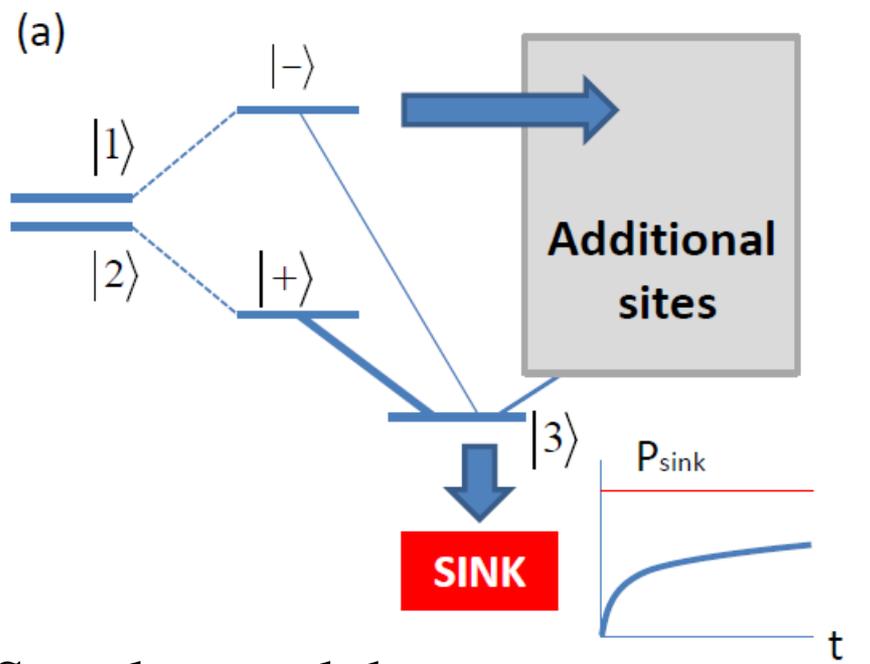
# Transport dynamics in FMO

(microscopic details)



# Transport dynamics in FMO

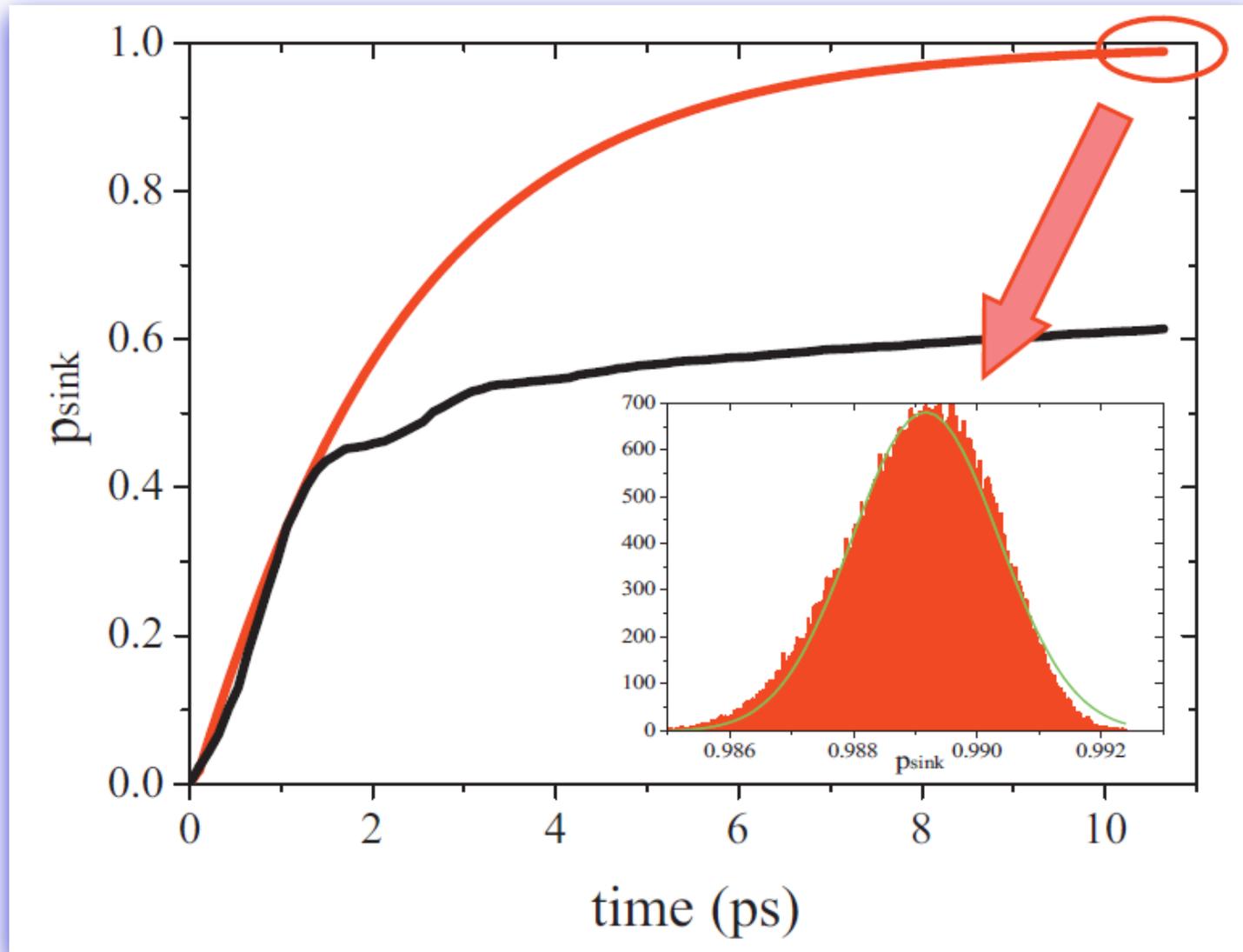
(microscopic details)



See also work by  
Aspuru-Guzik et al  
Ishizaki, Fleming et al  
Whaley et al  
Thorwart et al,....

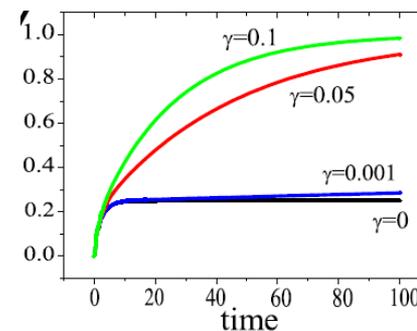
Caruso, Chin, Datta, Huelga, Plenio, J Chem Phys 2009  
Chin, Caruso, Datta, Huelga, Plenio, NJP 2010

# Robustness



**Noise leads to the creation of transport paths that were forbidden under a purely quantum evolution**

**How?**



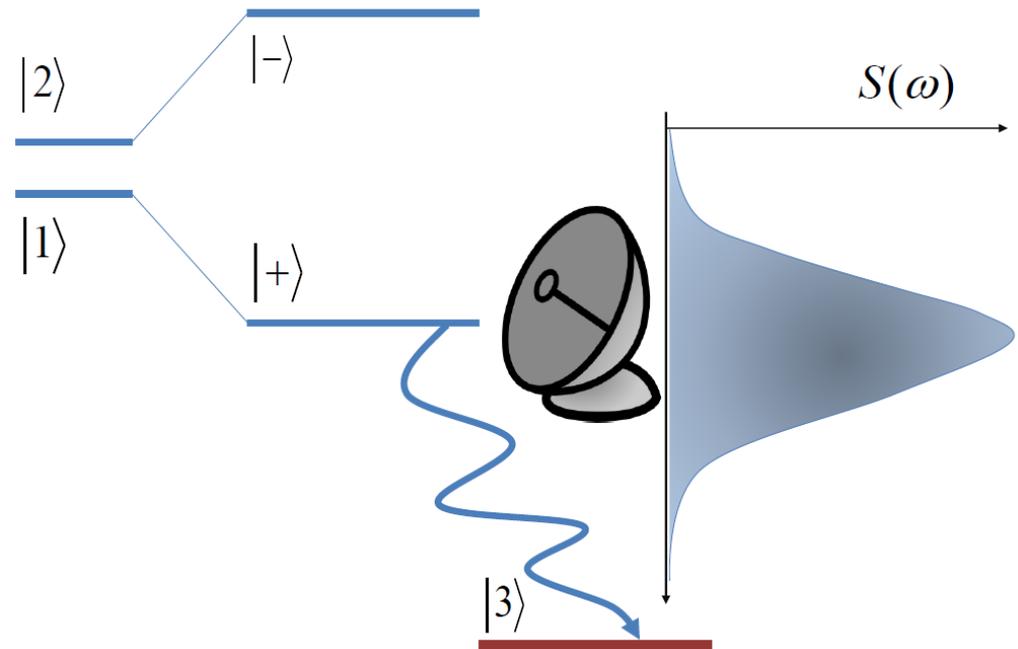
**Removal of destructive interference via dephasing of tunnelling amplitudes**  
**Activation of energetically unfavourable transition via line broadening**  
**Suppression of inefficient coherent paths**

**Moreover, the presence of coherent couplings can allow the system to efficiently harvest *noise* via the creation of a *phonon antennae***

Possible structure in the environment has been ignored so far  
Exciton relaxation may have a coherent component

## (3) Splitting of energy levels—Phonon Antennae

Electronic coupling may facilitate the creation of a **phonon antenna**



# (3) Splitting of energy levels—Phonon Antennae

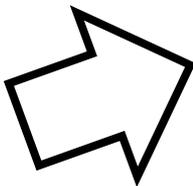
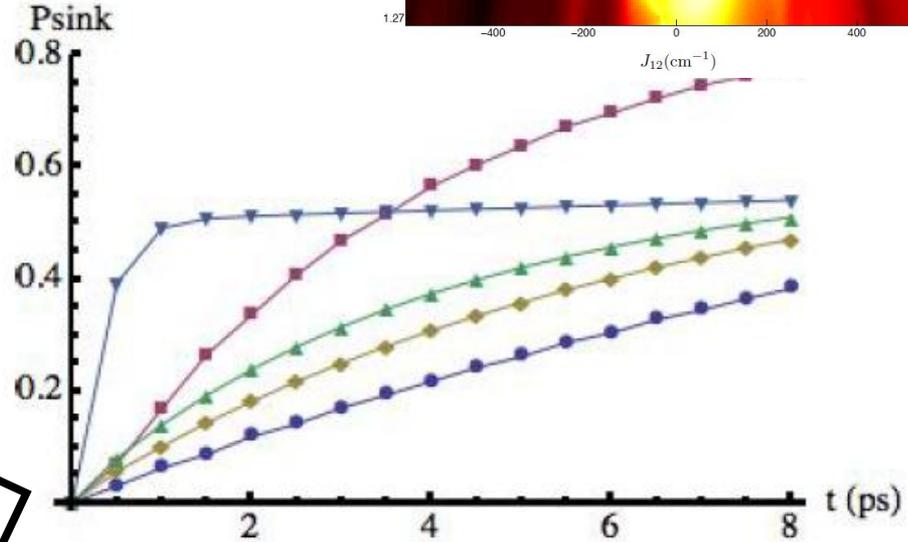
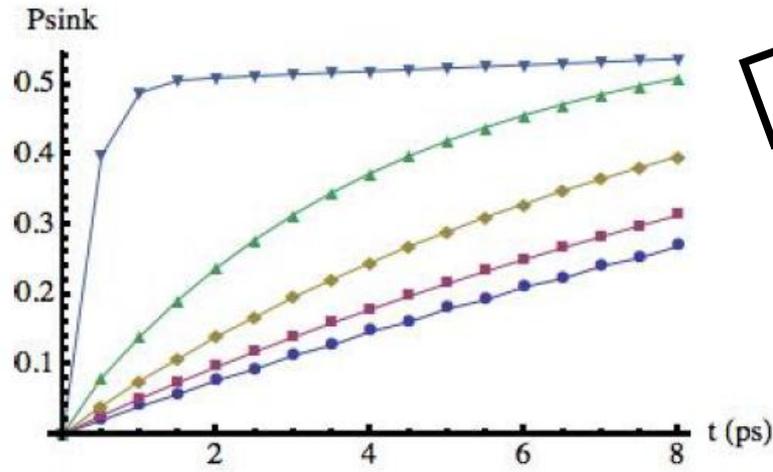
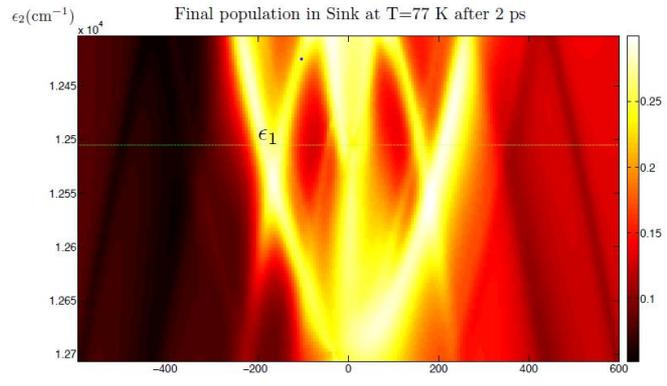
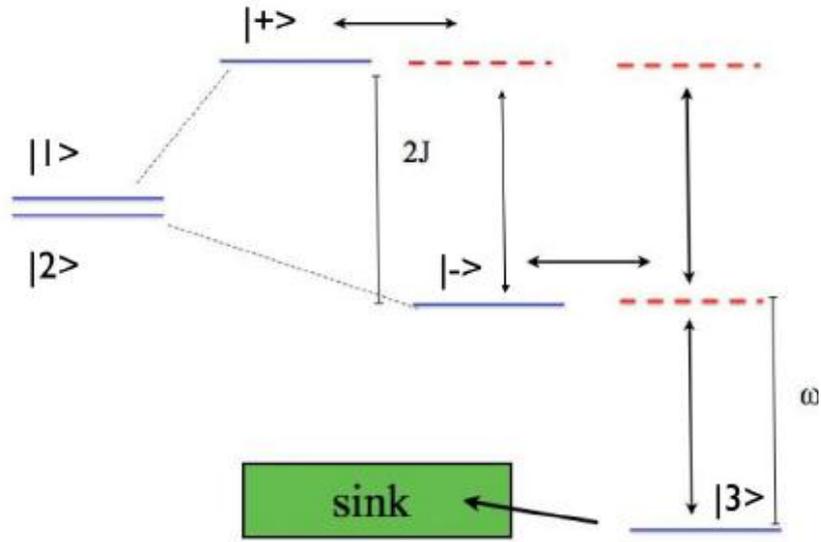
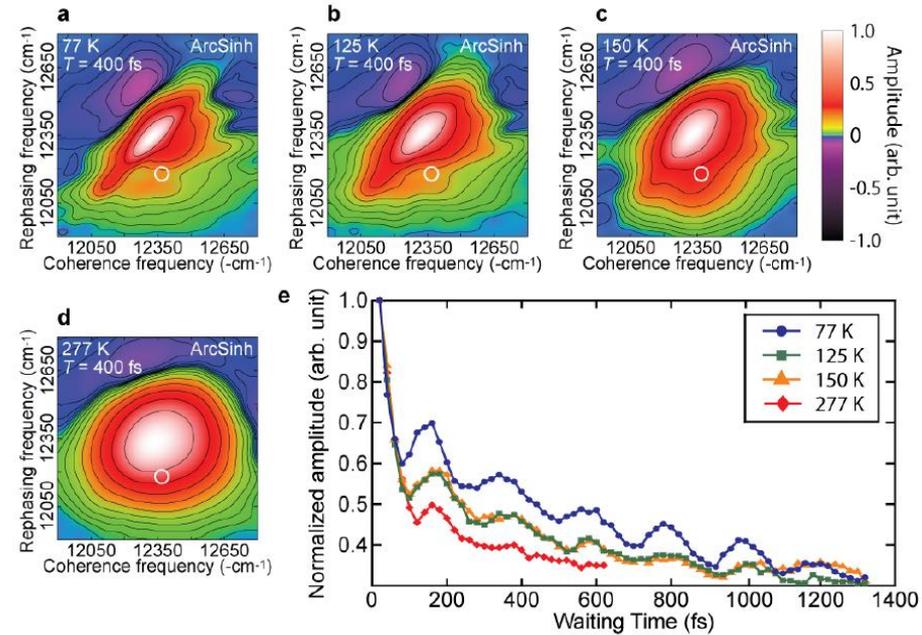


FIG. 6.  $P_{sink}$  as a function of time for a three-site system without any coupling to a mode. Weak transport is now seen for  $J_{12} = 100$  cm<sup>-1</sup> (squares). The other curves are  $J_{12} = 150$  cm<sup>-1</sup> (diamonds),  $J_{12} = 200$  cm<sup>-1</sup> (triangles),  $J_{12} = 300$  cm<sup>-1</sup> (inverted squares) and  $J_{12} = 50$  cm<sup>-1</sup> (dots). For all curves  $J_{23} = 30$  cm<sup>-1</sup> and the initial state of the vibrational mode is a thermal state at 77 K.

# Beyond transport

Engel et al, Science 2007  
Panitchayangkoon et al, PNAS 2010

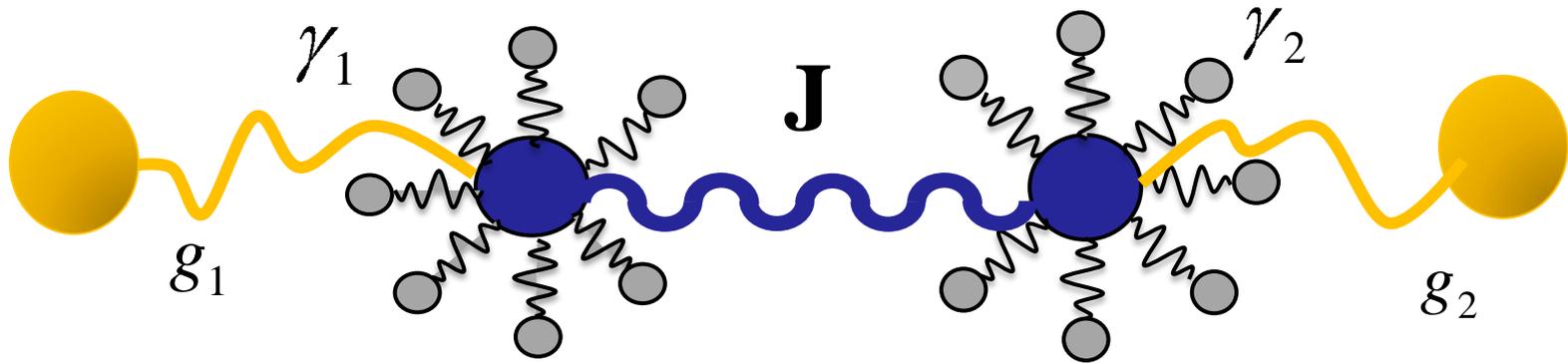


# What about the *long lasting* coherences?

Experimental evidence from traditional spectroscopy supports the existence of localized vibrational modes (in particular modes quasis resonant with excitonic transitions)

← Crucial??

# Dimer system subject to the action of damped local modes



$$H = \sum_{j=1}^2 \omega_j \sigma_j^+ \sigma_j^- + J (\sigma_1^- \sigma_2^+ + \sigma_1^+ \sigma_2^-)$$

$$H_{s-m} = g_j \sigma_z^j (a_j + a_j^\dagger)$$

$$\mathcal{L}_{deph}(\rho) = \sum_{j=1}^N \gamma_j [-\{\sigma_j^+ \sigma_j^-, \rho\} + 2\sigma_j^+ \sigma_j^- \rho \sigma_j^+ \sigma_j^-]$$

System experiences **dephasing in the site basis**  
**Populations are preserved**

# Exciton versus site coherence

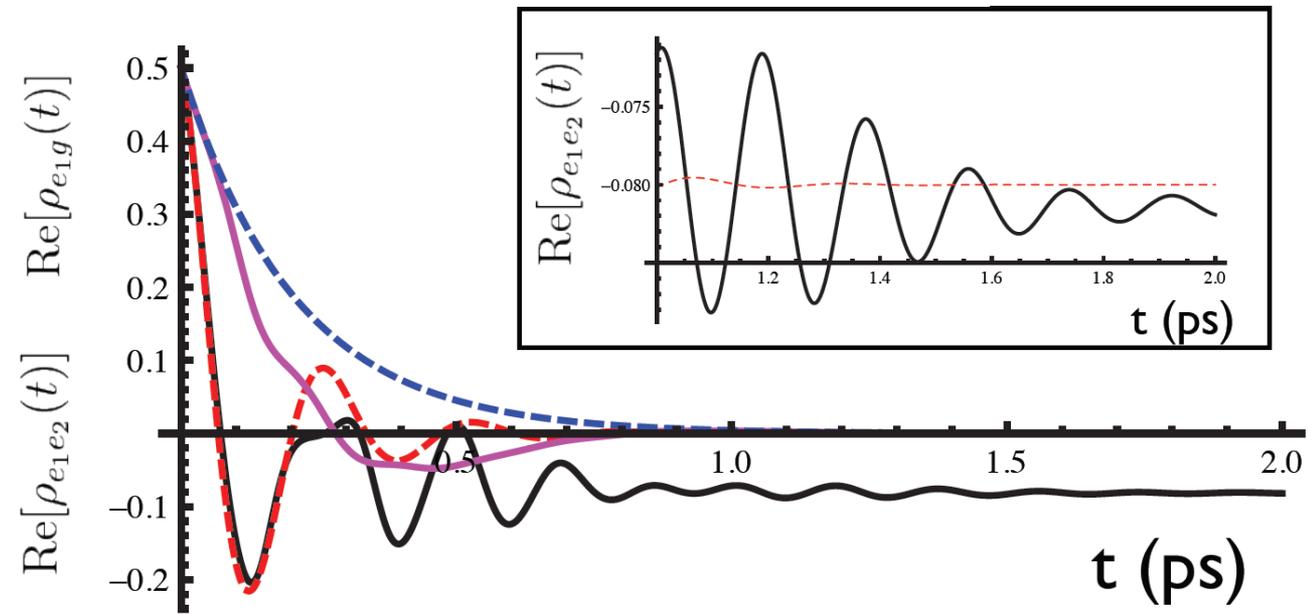
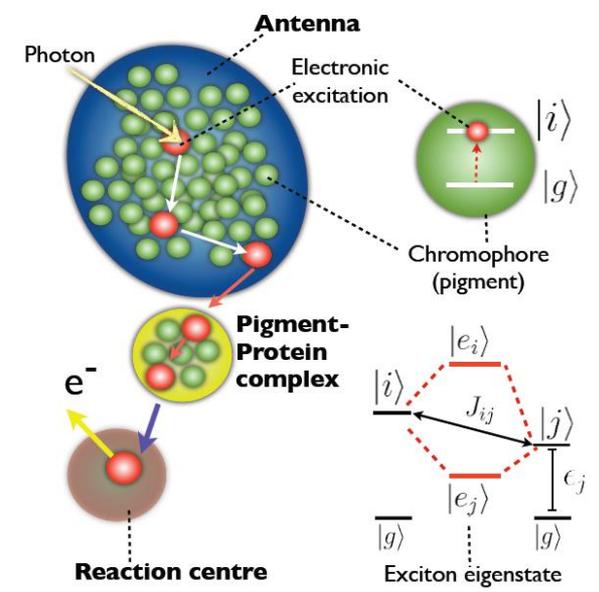
$$H = \sum_i ((\epsilon_i + X_i)|i\rangle\langle i| + H_i^B) + \sum_{i \neq j} J_{ij}|i\rangle\langle j|$$

$$X_i = \sum_k g_{ik}(a_{ik} + a_{ik}^\dagger)$$

$$H_I = \frac{1}{2} \sum_{n,m} (Q_{nm}|e_n\rangle\langle e_m| + h.c.),$$

$$Q_{nm} = \sum_{ik} \sqrt{S_k} \omega_k C_n^i C_m^i (a_{ik} + a_{ik}^\dagger)$$

$$H_{driving} \approx \frac{1}{2} \sum_{n \neq m} (\langle Q_{nm} \rangle(t) |e_n\rangle\langle e_m| + h.c.)$$



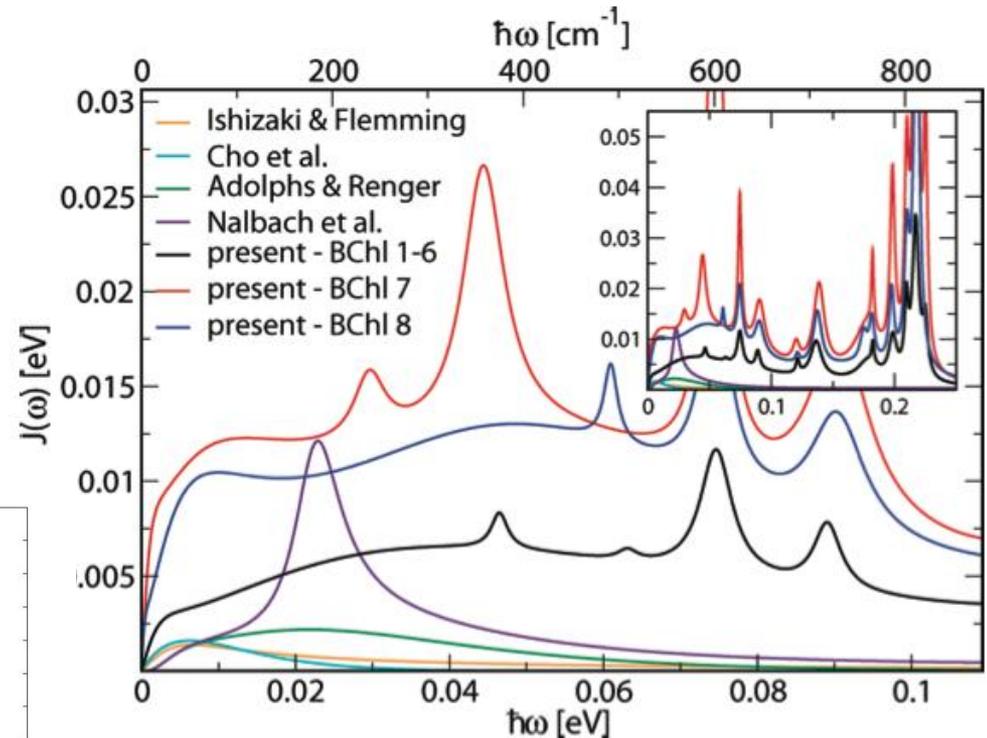
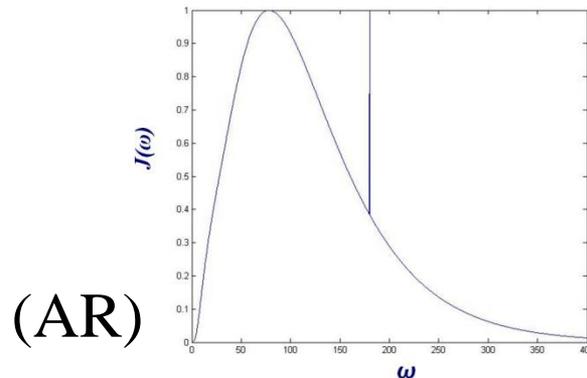
Semiclassical

# Towards a realistic description of long lasting coherence

## Standard open system theory inappropriate for pigment-protein complexes

-Slow fluctuations, long memory  
**Non-Markovian dynamics**

-Strong couplings and spatial correlation  
-Structured spectral function  
**non-perturbative**



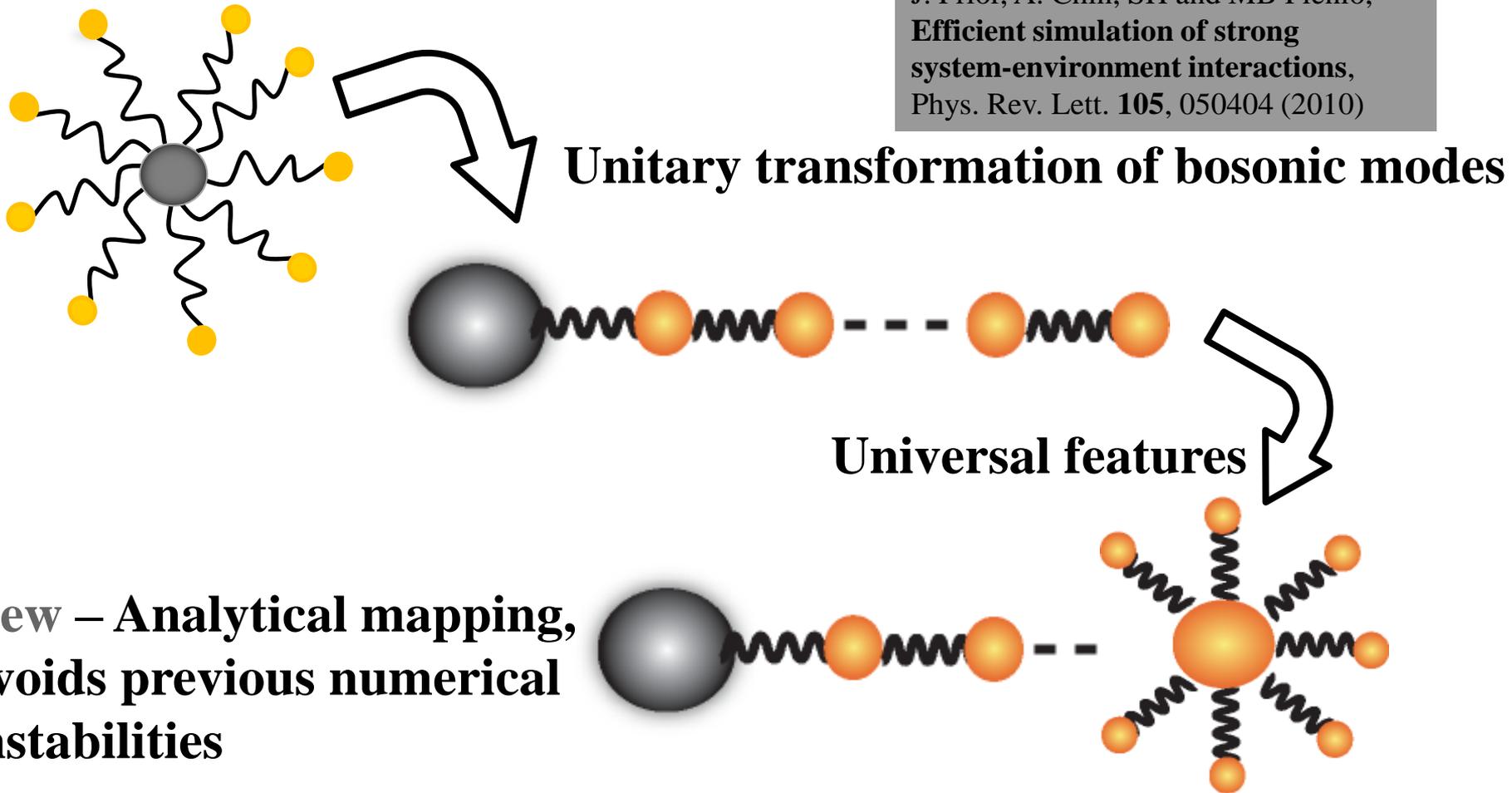
Olbrich et al. J. Phys. Chem Lett. 2. 2011

**Significantly correlated system-bath dynamics**  
-Many-body physics important

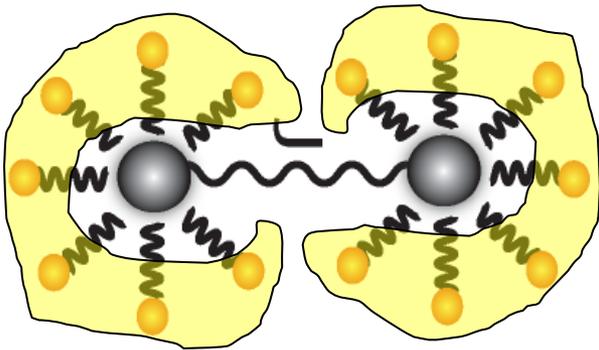
# Efficient exact simulation of many body systems (TEDOPA)

**Main result: Powerful numerical tool with associated physical insight**

J. Prior, A. Chin, SH and MB Plenio,  
**Efficient simulation of strong  
system-environment interactions,**  
Phys. Rev. Lett. **105**, 050404 (2010)

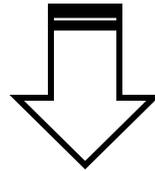


# Non-perturbative description of system-environment interaction for structured environments

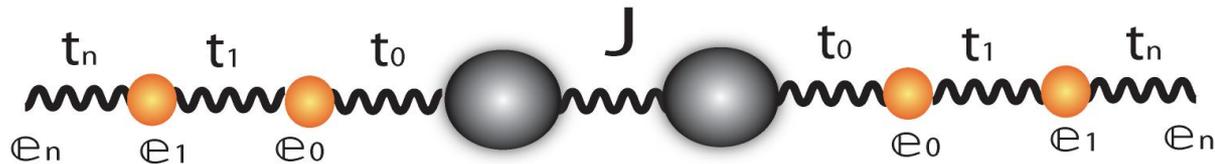
$$H_{res} = \int dx g(x) a_x^\dagger a_x$$


$$V = \int dx h(x) \hat{A}(a_x^\dagger + a_x)$$

$$b_n^\dagger = \int dx U_n(x) a_x^\dagger$$



Exact, thanks to theory of orthonormal polynomials

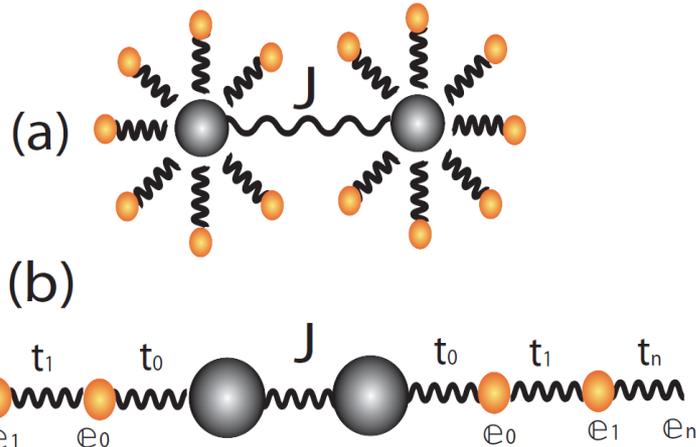


t-DMRG yields dynamics for general spectral densities

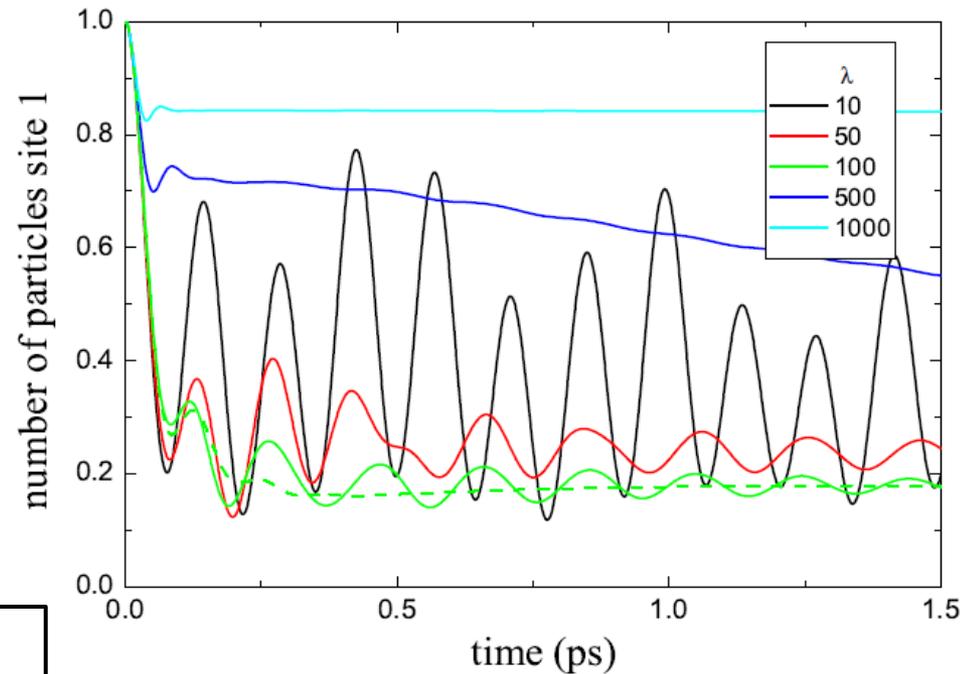
$$c_0 \hat{A}(b_0 + b_0^\dagger) + \sum_{n=0}^{\infty} \omega_n b_n^\dagger b_n + t_n b_{n+1}^\dagger b_n + t_n b_n^\dagger b_{n+1}$$

J. Math. Phys. **51**, 092109 (2010)

# Efficient simulation of strong system-environment interactions: dimer system (building block of complex bio-molecular aggregates)

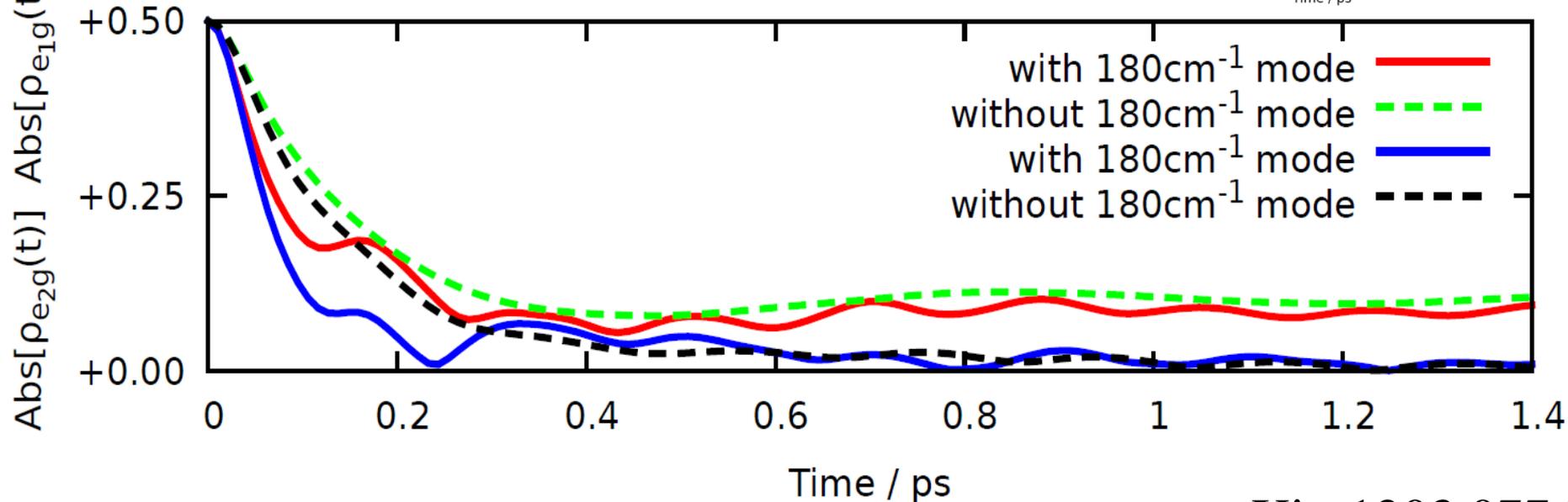
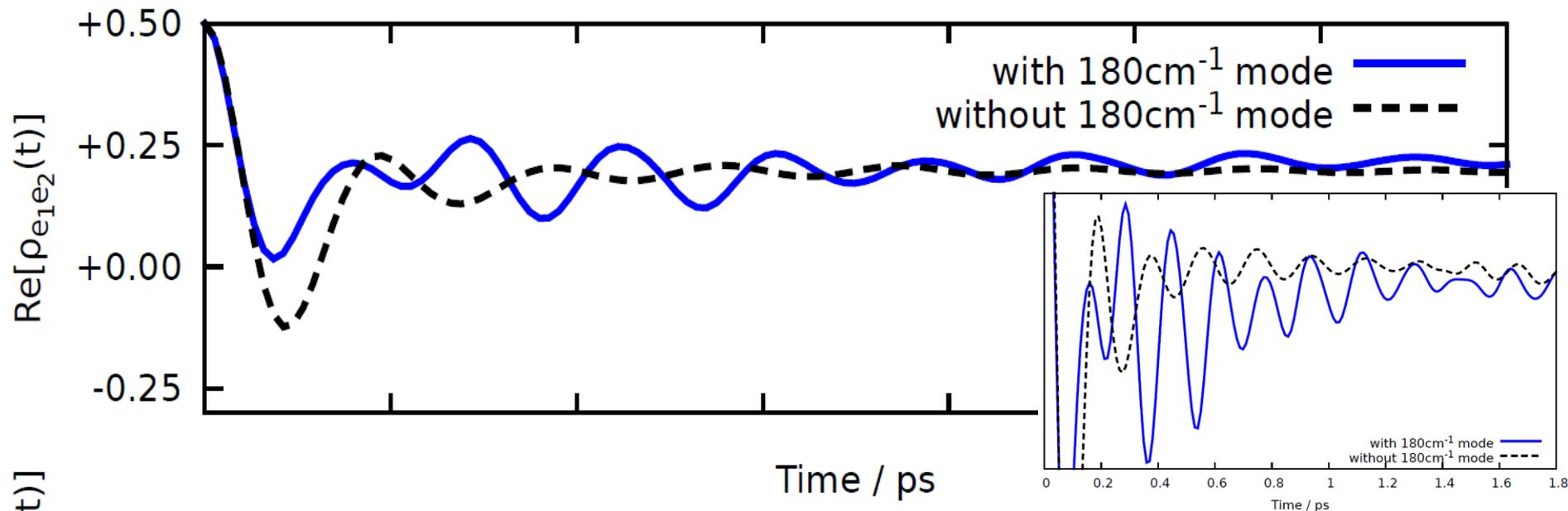


$$J(\omega) = \frac{2\pi\lambda [1000\omega^5 e^{-\left(\frac{\omega}{\omega_1}\right)^{\frac{1}{2}}} + 4.3\omega^5 e^{-\left(\frac{\omega}{\omega_2}\right)^{\frac{1}{2}}}]}{9!(1000\omega_1^5 + 4.3\omega_2^5)} + 4\pi S_H \omega_H^2 \delta(\omega - \omega_H),$$

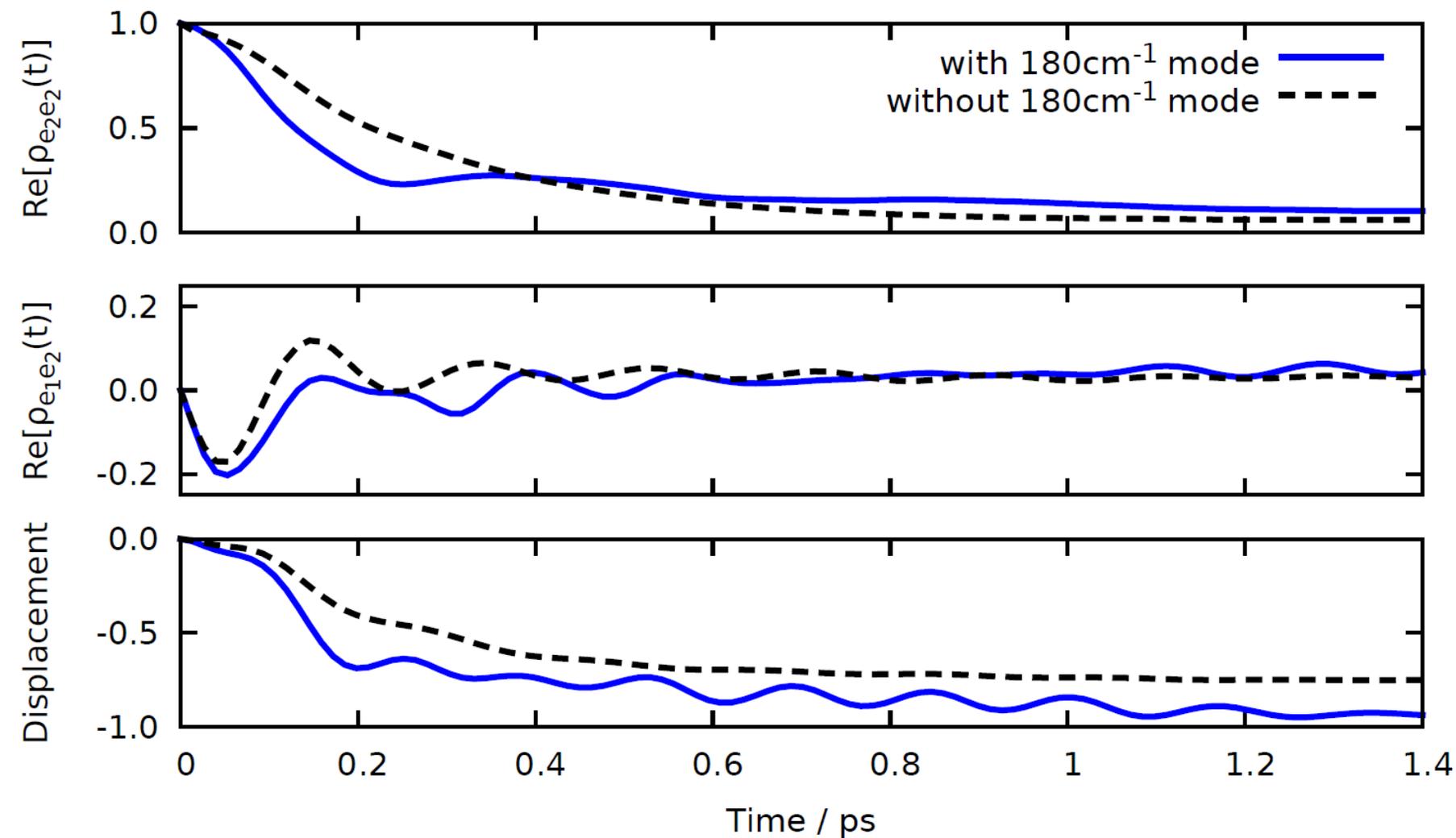


Can obtain good fits to  
beating patterns as reported by Engel's group  
(Hayes et al. NJP. 2. (2010);  
Hayes et al. Faraday Disc. 150. (2011))

# Exact Simulations (T=77 K)



# Exact Simulations



arXiv:1203.0776

**Other results involving localized modes:**

Mancal & Pullerits, Olaya-Castro & Scholes, Kreisbeck & Kramer (**arXiv 2012**)

# Conclusions

**- The basis for efficient and robust transport in the FMO complex is provided by an intricate interplay of environmental protein interactions and coherent pigment-pigment couplings**

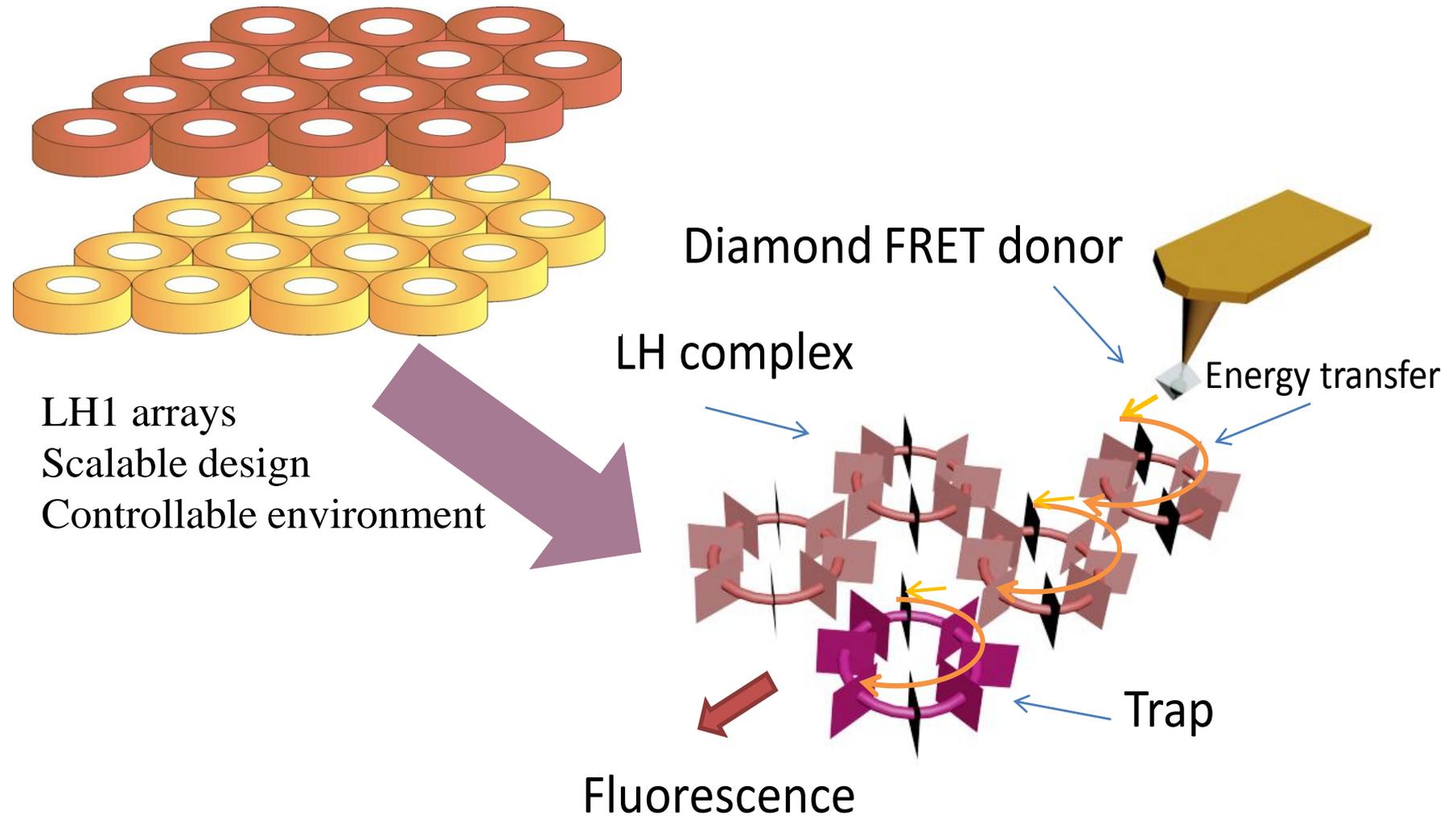
**-Structure in the protein spectral density leads to strong coherent pigment-protein interactions – boundary of *system* and *environment* is obscured. Many-body treatment of open quantum dynamics needed**

**- Non-equilibrium system-bath (coherent) effects may be behind the enhancement of inter-exciton coherences**

# To do ....

- Simulate larger networks at finite temperature and with full structure
- Evaluation of spectral density from excitonically coupled systems (test bed WSCP)
- Compute 2D spectra from non-perturbative dynamics
- Extract clear dynamical features via different experimental techniques (quantum control)

# What about the *big* question?



Collaboration with Robin Gosh (Stuttgart) and Fedor Jelezko (Ulm)



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**corner**





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ulm university universität  
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Alexander von Humboldt  
Stiftung / Foundation

## Main references from our group

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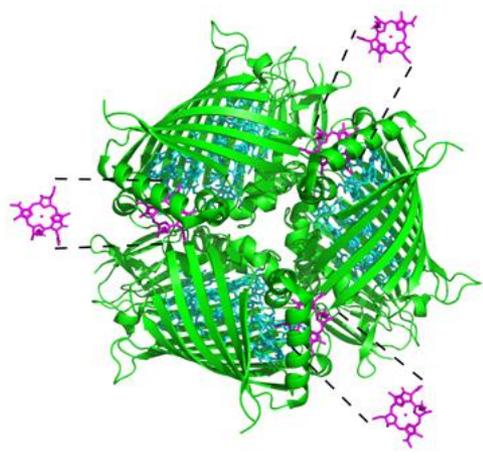
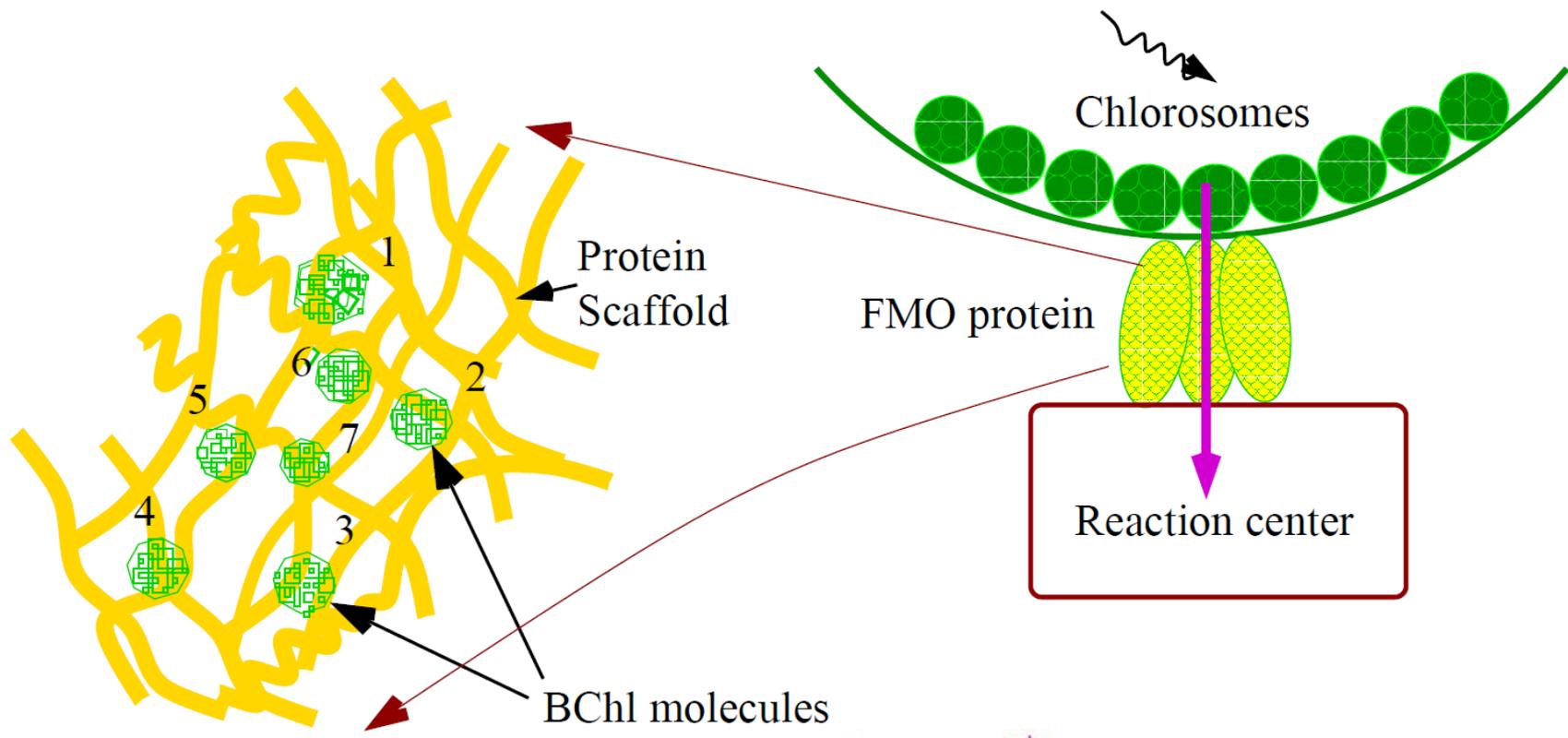
Procedia Chemistry 3, 248 (2011)

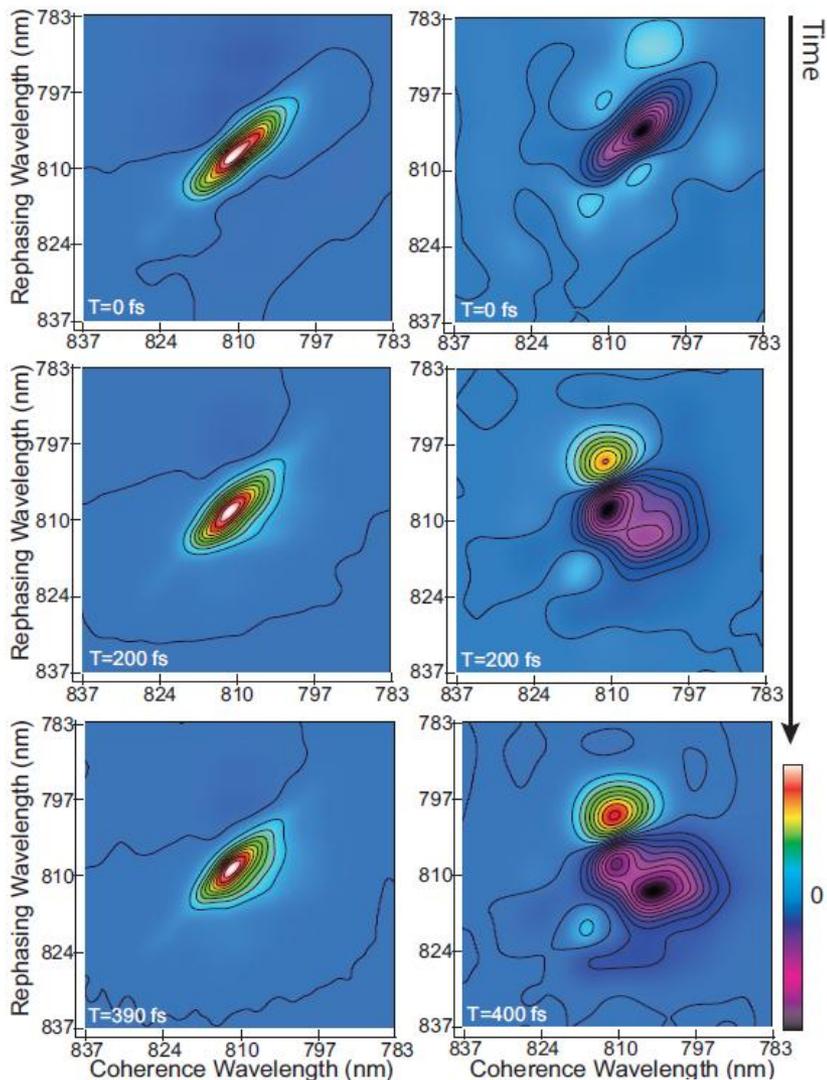
Phys. Rev. A 85, 042331 (2012)

J. Chem. Phys. 136, 155102 (2012)

arXiv:1203.5072

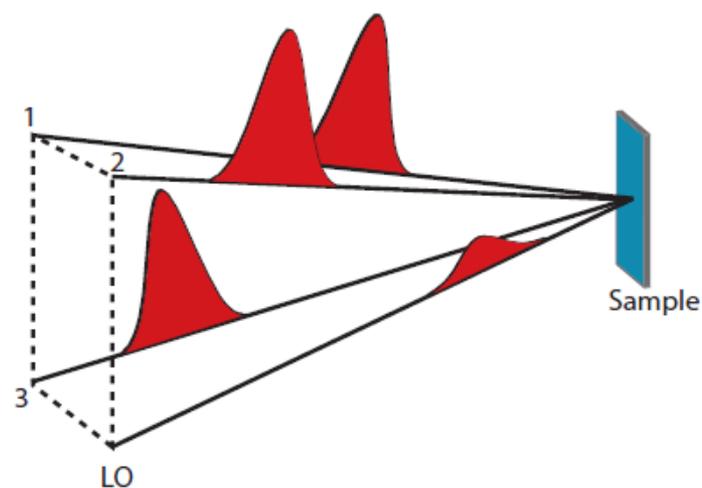
arXiv:1203.0776





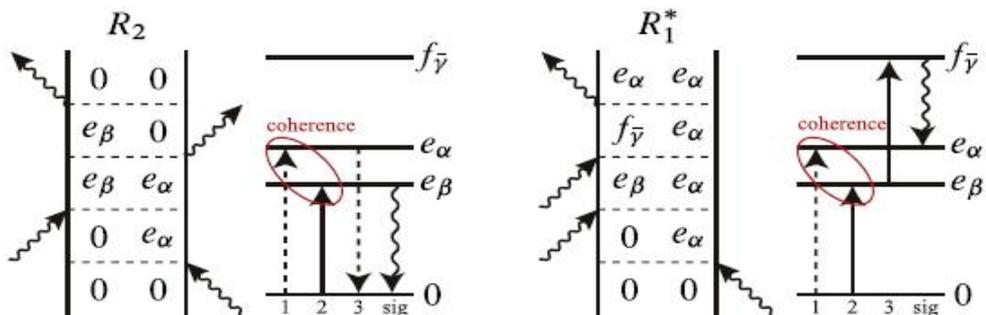
## Cross-peak-specific two-dimensional electronic spectroscopy

Elizabeth L. Read<sup>†‡</sup>, Gregory S. Engel<sup>†‡</sup>, Tessa R. Calhoun<sup>†‡</sup>, Tomáš Mančal<sup>†‡§</sup>, Tae Kyu Ahn<sup>†‡</sup>, Robert E. Blankenship<sup>¶</sup>, and Graham R. Fleming<sup>†¶||</sup>



**Fig. 1.** The evolution of the 2D electronic spectrum of *P. phaeum* FMO is shown. In the conventional 2D spectra (*Left*), the emergence of cross-peaks below the main diagonal is evidenced by the contour lines bowing away from the diagonal by 200 fs. In the cross-peak-specific spectra (*Right*), the emergence of the negative features below the main diagonal corresponds to the cross-peaks largely obscured in the conventional 2D spectra.

### Rephasing



### Nonrephasing

